

# Anomalies in B decays

University of Geneva,  
April 25, 2018

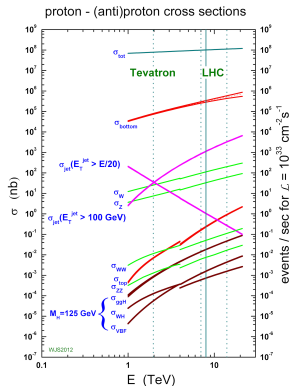
Michel De Cian, EPFL

# B physics

mass GeV/c <sup>2</sup>	$\sim 2.4$	$\sim 1.275$	$\sim 172.44$	0	$\sim 125.09$
charge	2/3	2/3	2/3	0	0
spin	1/2	1/2	1/2	1	0
	<b>u</b> up	<b>c</b> charm	<b>t</b> top	<b>g</b> gluon	<b>H</b> Higgs
<b>QUARKS</b>					
	$\sim 4.8$	$\sim 95$	$\sim 4.18$	0	
	-1/3	-1/3	-1/3	0	
	1/2	1/2	1/2	1	
	<b>d</b> down	<b>s</b> strange	<b>b</b> bottom	<b><math>\gamma</math></b> photon	
<b>GAUGE BOSONS</b>					
	$\sim 0.511$	$\sim 105.67$	$\sim 1.7768$	$\sim 81.19$	
	1	1	1	0	
	1/2	1/2	1/2	1	
	<b>e</b> electron	<b><math>\mu</math></b> muon	<b><math>\tau</math></b> tau	<b>Z</b> Z boson	
<b>LEPTONS</b>					
	$\sim 2.2$	$\sim 1.7$	$\sim 15.5$	$\sim 80.39$	
	0	0	0	1	
	1/2	1/2	1/2	1	
	<b><math>\nu_e</math></b> electron neutrino	<b><math>\nu_\mu</math></b> muon neutrino	<b><math>\nu_\tau</math></b> tau neutrino	<b>W</b> W boson	

$$B^0 = |d\bar{b}\rangle, B^+ = |u\bar{b}\rangle, B_s^0 = |s\bar{b}\rangle, \Lambda_b^0 = |udb\rangle, B_c^+ = |c\bar{b}\rangle$$

- B physics is the study of bound states containing one  $b$  quark and their decays / dynamics.
- They decay in a multitude of final states, allowing the study of a wide range of physics.
- They are copiously produced at the LHC.



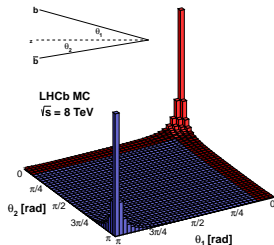
## LHC physics (toy version)



- Take a bunch of protons
- Smash them together and create a mess.
- Spend some millions to build a device to understand it.

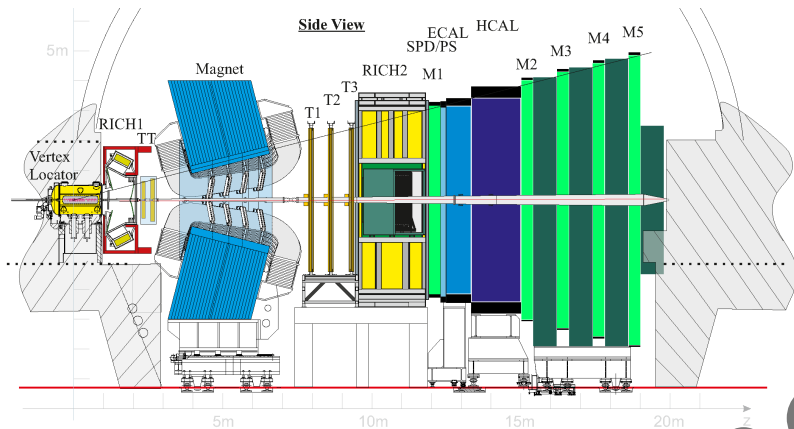


# B physics at the LHC

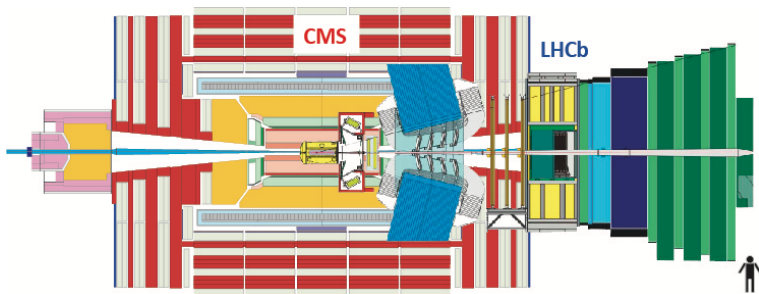


- b hadrons are moderately heavy ( $\sim 5 \text{ GeV}$ ,  $> \Lambda_{QCD}$ ) and are mainly produced in the forward or backward direction at the LHC  
→ build a forward detector.
- B hadrons have "soft" decay products and travel  $\sim 1 \text{ cm}$  before decaying  
→ build a detector with low- $p_T$  capability and good momentum / vertex resolution.
- B hadrons have a large variety of decay channels with different particle species in the final states  
→ need a particle ID.

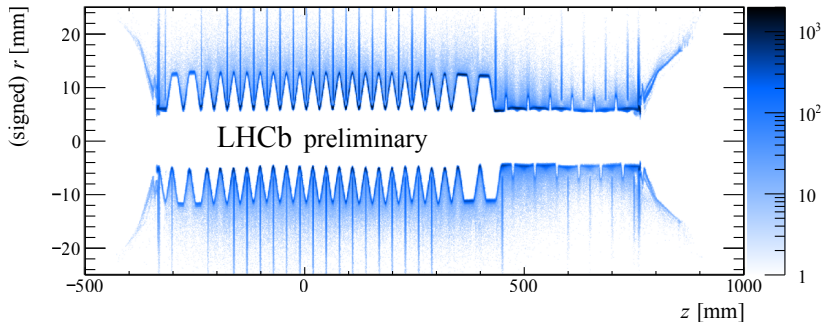
# LHCb



# LHCb: Comparison with CMS

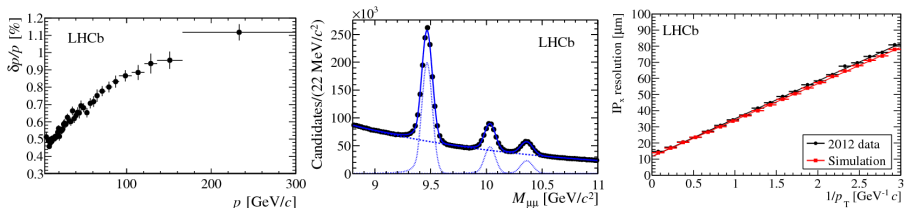


## LHCb: Vertex Locator

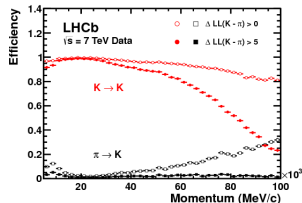


- Velo sensors 8mm from the beam position.
- Allows for very good Impact Parameter and Primary / Secondary Vertex resolution.

# LHCb: Performance numbers

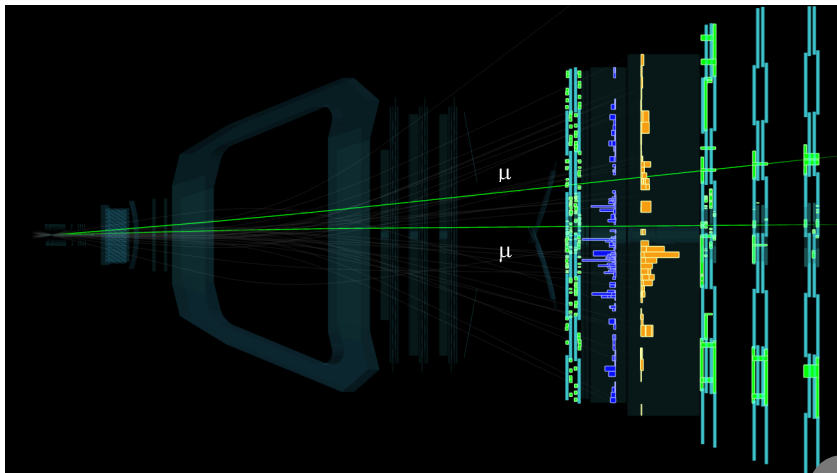


- Excellent momentum / mass resolution:
  - $\frac{\delta p}{p} = 0.5\% (10 \text{ GeV}/c) - 1.0\% (200 \text{ GeV}/c)$
  - $\sigma_m(B_s^0 \rightarrow \mu^+ \mu^-) \approx 20 \text{ MeV}/c^2$
- Impact parameter resolution:
  - $15 + 29/p_T [ \text{GeV}/c ] \mu\text{m}$
- High particle identification efficiency.
  - $\varepsilon_\mu \approx 97\%$  with 1-3%  $\pi \rightarrow \mu$  misidentification
  - $\varepsilon_K \approx 95\%$  with  $\approx 5\%$   $\pi \rightarrow K$  misidentification



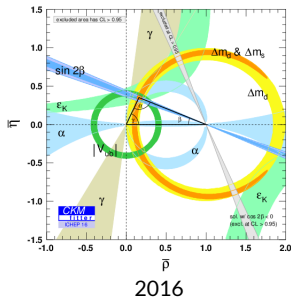
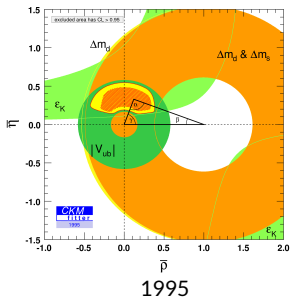


# LHCb: Event display



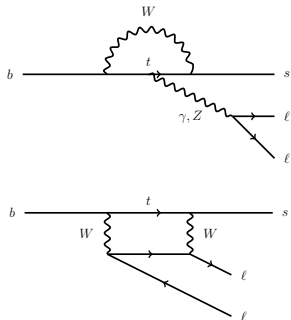
# Rare B Decays

## B decays as a laboratory



- B hadrons are a perfect laboratory to perform measurements of many fundamental physics quantities.
- Decays governed by (electro)weak interaction, but hadronic state itself by strong interaction.
- B physics is by definition flavour physics and strongly linked to the CKM matrix.

# Rare B decays

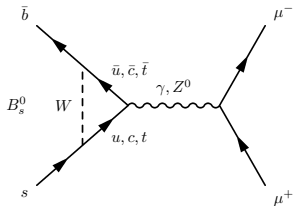


		neutral current		
charged current	Quarks	2.4 MeV $\frac{2}{3}$ $\frac{1}{2}$ <b>u</b> up	1.27 GeV $\frac{2}{3}$ $\frac{1}{2}$ <b>c</b> charm	171.2 GeV $\frac{2}{3}$ $\frac{1}{2}$ <b>t</b> top
		4.8 MeV $-\frac{1}{3}$ $\frac{1}{2}$ <b>d</b> down	104 MeV $-\frac{1}{3}$ $\frac{1}{2}$ <b>s</b> strange	4.2 GeV $-\frac{1}{3}$ $\frac{1}{2}$ <b>b</b> bottom
		<2.2 eV 0 $\frac{1}{2}$ <b><math>\nu_e</math></b> electron neutrino	<0.17 MeV 0 $\frac{1}{2}$ <b><math>\nu_\mu</math></b> muon neutrino	<15.5 MeV 0 $\frac{1}{2}$ <b><math>\nu_\tau</math></b> tau neutrino
		0.511 MeV -1 $\frac{1}{2}$ <b>e</b> electron	105.7 MeV -1 $\frac{1}{2}$ <b><math>\mu</math></b> muon	1.777 GeV -1 $\frac{1}{2}$ <b><math>\tau</math></b> tau
		Leptons		

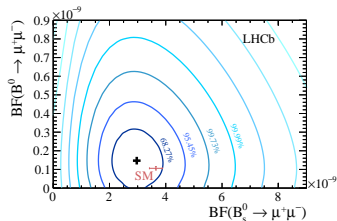
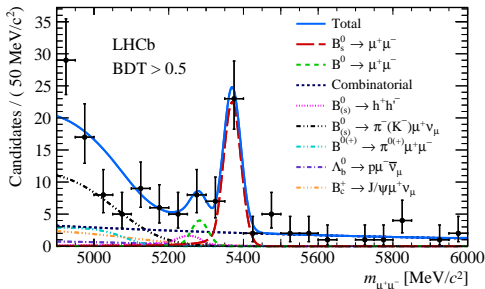
- Will (mostly) take about flavour-changing neutral current decays of b quarks today.
- And only consider electroweak interactions (no gluonic penguins).
- Decays are strongly suppressed, but heavy new particles (beyond the SM) can appear in the loop and alter the final state distributions.
- "Rare B decays":  $\mathcal{B} \sim 10^{-6}$

*B*

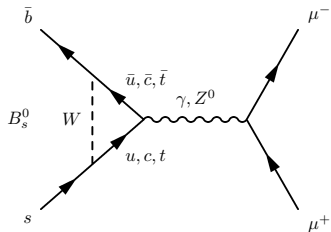
$$B_s^0 \rightarrow \mu^+ \mu^-$$



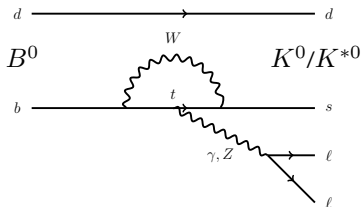
- Strongly helicity suppressed.
- Only leptons in the final state, very clean theoretically.
- $\mathcal{B}(B_s^0 \rightarrow \mu^+ \mu^-) = (3.0 \pm 0.6_{-0.2}^{+0.3}) \cdot 10^{-9}$
- $\mathcal{B}(B^0 \rightarrow \mu^+ \mu^-) < 3.4 \cdot 10^{-10}$
- All compatible with standard model predictions.



$b \rightarrow sl^+l^-$



$$m^2(\gamma, Z) = q^2 = m^2(B_s^0)$$



$$m^2(\gamma, Z) = q^2 > 4m_\ell^2$$

*i.e.* physics depends on  $q^2$

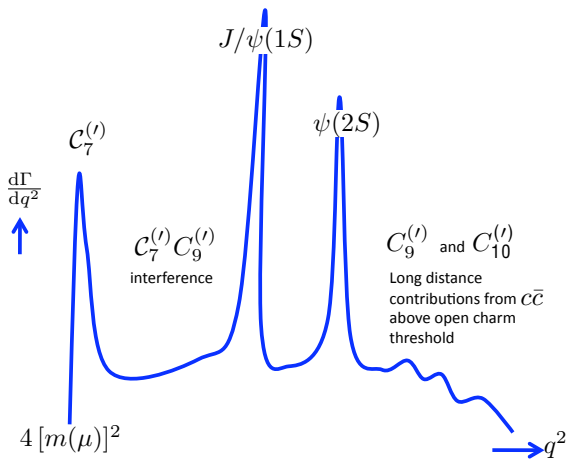


$$\frac{d\mathcal{B}}{dq^2}$$

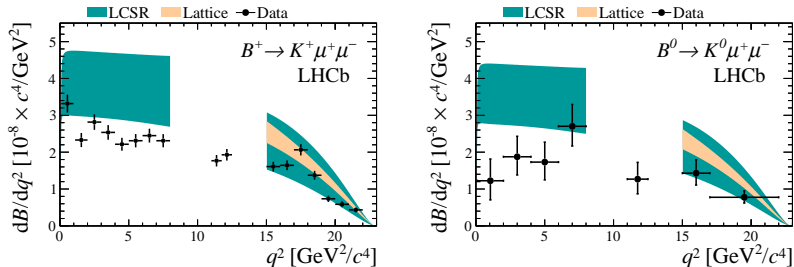
All results with  $3 \text{ fb}^{-1}$  (Run I)



# $q^2$ spectrum for $b \rightarrow s \ell^+ \ell^-$

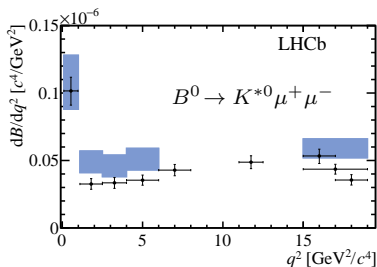
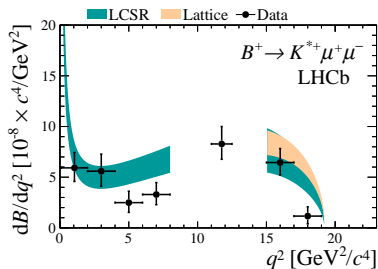


## $b \rightarrow s l^+ l^-$ differential branching fractions



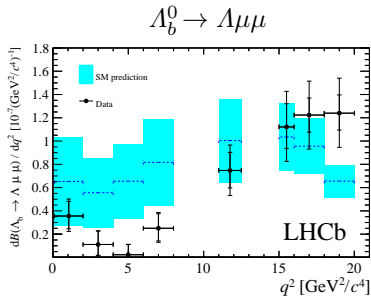
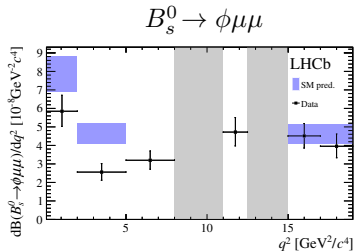
- Measure branching fractions as a function of  $q^2$ .
  - Normalize  $B \rightarrow K \mu^+ \mu^-$  to  $B \rightarrow K J/\psi$
- For  $B^0 \rightarrow K^0 \mu^+ \mu^-$  use decay  $K_s^0 \rightarrow \pi^+ \pi^-$
- Measured values significantly below prediction for low  $q^2$ .

## $b \rightarrow s l^+ l^-$ differential branching fractions



- Measure branching fractions as a function of  $q^2$ 
  - Normalize  $B \rightarrow K^* \mu^+ \mu^-$  to  $B \rightarrow K^* J/\psi$
- For  $B^+ \rightarrow K^{*+} \mu^+ \mu^-$  use decay  $K^{*+} \rightarrow K_S^0 \pi^+$
- Measured values below prediction for low  $q^2$ .

## $b \rightarrow sl^+l^-$ differential branching fractions

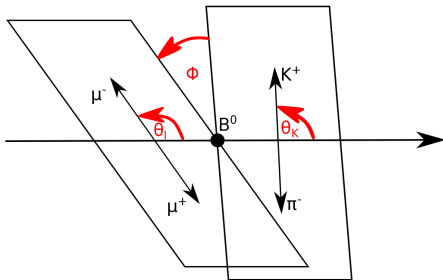


- Measure branching fractions as a function of  $q^2$ 
  - Normalize  $B_s^0 \rightarrow \phi\mu^+\mu^-$  to  $B_s^0 \rightarrow \phi J/\psi$
  - Normalize  $\Lambda_b^0 \rightarrow \Lambda\mu^+\mu^-$  to  $\Lambda_b^0 \rightarrow \Lambda J/\psi$
- Basically all differential branching fractions are lower than their prediction for low values of  $q^2$ .

$$\frac{d\mathcal{B}}{dq^2 d\Omega}$$

All results with  $3 \text{ fb}^{-1}$  (Run I)

## Angular analysis of $B^0 \rightarrow K^{*0} \mu^+ \mu^-$



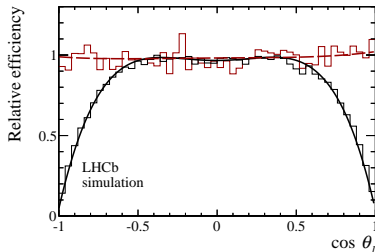
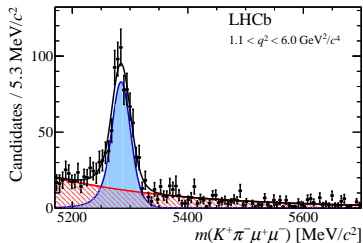
- Differential decay rate of  $P \rightarrow VV$  decays depends on 3 decay angles and an observable, depending on  $q^2$ .
- $$\frac{1}{\Gamma} \frac{d\Gamma}{d \cos \theta_\ell d \cos \theta_K d\phi} = \frac{9}{32\pi} \sum_i J_i(q^2) f(\cos \theta_\ell, \cos \theta_K, \phi)$$
- Best studied case in LHCb:  $B^0 \rightarrow K^{*0} \mu^+ \mu^-$

## Angular analysis of $B^0 \rightarrow K^{*0} \mu^+ \mu^-$

$$\frac{d^4(\Gamma + \bar{\Gamma})}{d \cos \theta_\ell d \cos \theta_K d\phi dq^2} = \frac{9}{32\pi} \left[ \frac{3}{4}(1 - F_L) \sin^2 \theta_K + F_L \cos^2 \theta_K + \right. \\ \left. \frac{1}{4}(1 - F_L) \sin^2 \theta_K \cos 2\theta_\ell - F_L \cos^2 \theta_K \cos 2\theta_\ell + \right. \\ S_3 \sin^2 \theta_K \sin^2 \theta_\ell \cos 2\phi + S_4 \sin 2\theta_K \sin 2\theta_\ell \cos \phi + \\ S_5 \sin 2\theta_K \sin \theta_\ell \cos \phi + S_6 \sin^2 \theta_K \cos \theta_\ell + \\ S_7 \sin 2\theta_K \sin \theta_\ell \sin \phi + \\ \left. S_8 \sin 2\theta_K \sin 2\theta_\ell \sin \phi + S_9 \sin^2 \theta_K \sin^2 \theta_\ell \sin 2\phi \right]$$

- Call the coefficient in front of the angular expressions "observable".
- Angular terms are (almost all) orthogonal.
- $S_6 = \frac{4}{3} A_{FB} = \frac{4}{3} \frac{\# \cos \theta_\ell > 0 - \# \cos \theta_\ell < 0}{\# \cos \theta_\ell > 0 + \# \cos \theta_\ell < 0}$ : Forward-backward asymmetry of the leptons.
- $F_L$ : Longitudinal polarization of the  $K^{*0}$

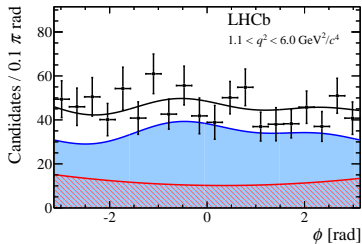
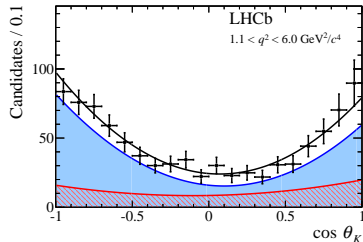
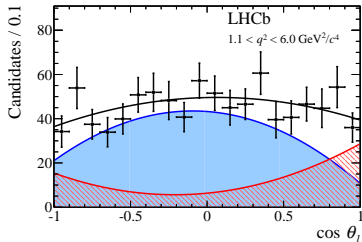
# Angular analysis of $B^0 \rightarrow K^{*0} \mu^+ \mu^-$



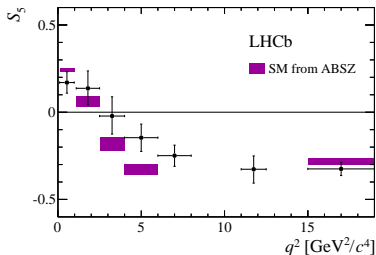
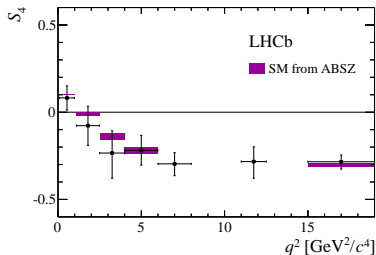
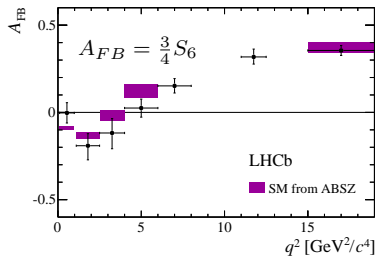
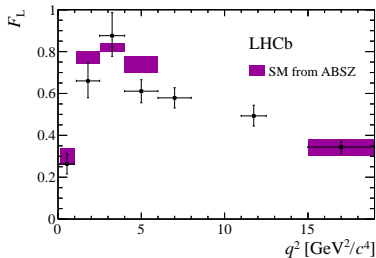
- Use a BDT to select the events, in total  $\approx 2400$  signal candidates.
- Use simulated sample of phase-space generated events to determine the effect of the acceptance and selection.



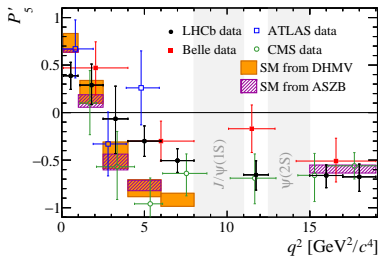
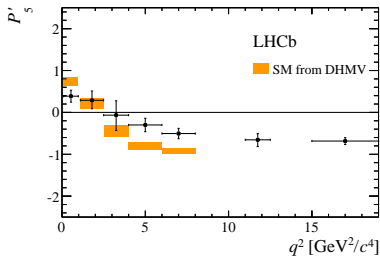
# Angular analysis of $B^0 \rightarrow K^{*0} \mu^+ \mu^-$



# Angular analysis of $B^0 \rightarrow K^{*0} \mu^+ \mu^-$

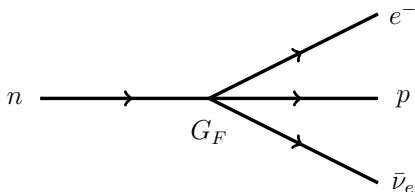


# Angular analysis of $B^0 \rightarrow K^{*0} \mu^+ \mu^-$



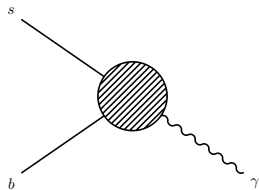
- $P'_5 = \frac{S_5}{\sqrt{1-F_L}}$
- The  $P'_i$  observables are less prone to hadronic form-factor uncertainties than the  $S_i$  ones (when using so-called "soft form-factors").
- Measurement also by Belle, CMS, ATLAS (but with less statistical power)
- Global significance is about  $3.4\sigma$  from the SM (LHCb measurement alone).

## Wilson coefficients (I)

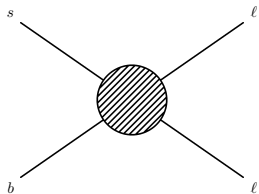


- Need a framework to describe all these different types of processes, with as little assumptions as possible.
- Fermi solved this problem already 85 years ago for the  $\beta$  decay by introducing a point interaction.
- $G_F$  is a coupling constant that gives the strength of the interaction, as long as  $E \ll m_W$ .

## Wilson coefficients (II)



$\mathcal{O}_7$

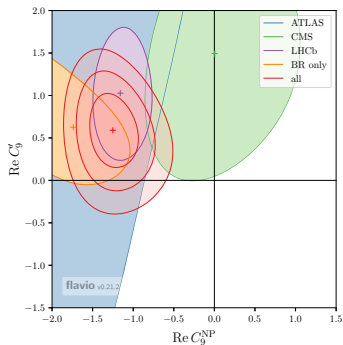
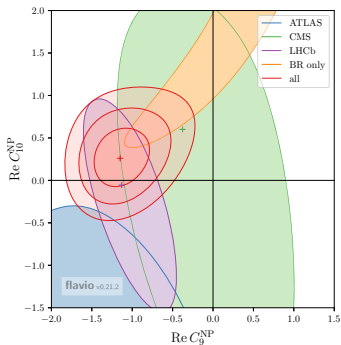


$\mathcal{O}_{9V}, \mathcal{O}_{10A}$

$$\mathcal{H}_{eff} = -4 \frac{G_F}{\sqrt{2}} V_{tb} V_{ts}^* \sum_i \mathcal{C}_i \mathcal{O}_i$$

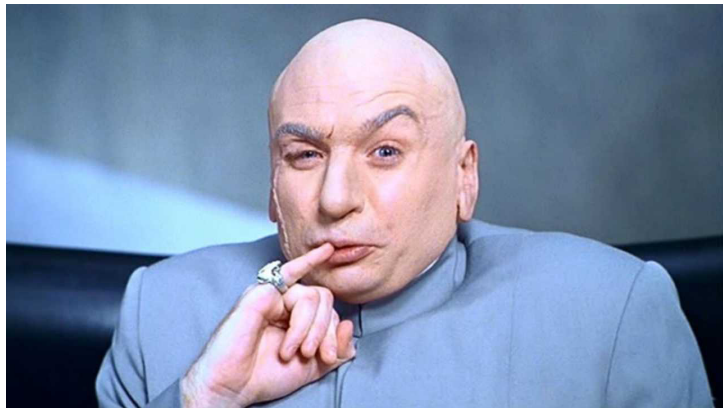
- Do the same thing for the 3 possible interactions in  $b \rightarrow s \ell^+ \ell^-$  processes.
- $G_F$  is Fermi constant,  $V_{tb}$ ,  $V_{ts}^*$  CKM elements.
- $\mathcal{C}_i$  are called Wilson coefficients, they are (complex) numbers.
- Derive  $\mathcal{C}_i$  from all measurements and combine them in global fits.

## Global fits (part I)

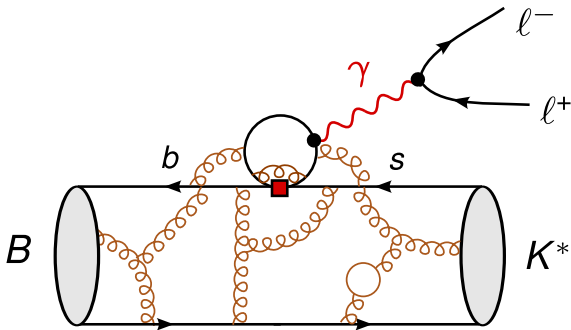


- Only consider  $C_9$  and  $C_{10}$ .
- 0.0 / 0.0 is the Standard Model.
- Does not include lepton universality measurements (see later).
- $> 3\sigma$  away from Standard Model. Is it new physics?

## The villain



## The villain



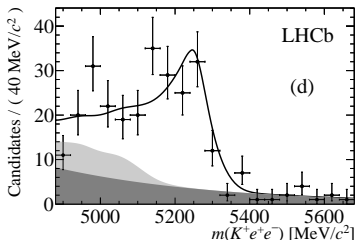
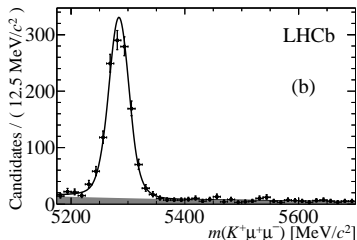
- This could mimic a new physics effect in  $C_9$ , and is not included in the uncertainties of the hadronic form-factors.
- One could measure effect of charm-loops by a precise analysis of the  $\mu\mu$  invariant mass ( $\rightarrow$  backup).



# LFU

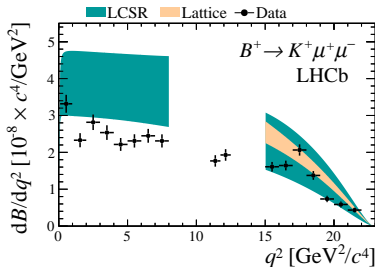
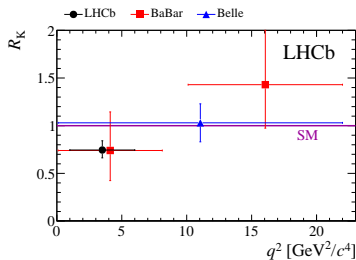
All results with  $3 \text{ fb}^{-1}$  (Run I)

# Lepton Flavour universality in loop decays



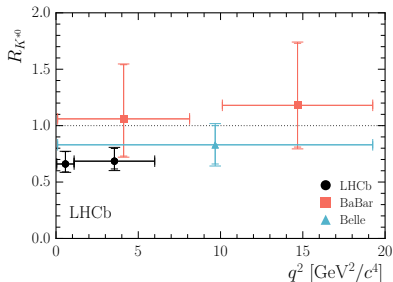
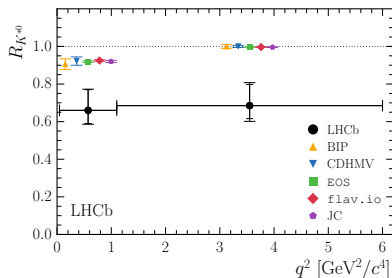
- Measure  $R_K = \frac{\mathcal{B}(B^+ \rightarrow K^+ \mu^+ \mu^-)}{\mathcal{B}(B^+ \rightarrow K^+ e^+ e^-)}$  in  $q^2 \in [1, 6] \text{ GeV}^2/c^4$
- Use  $B^+ \rightarrow J/\psi K^+, J/\psi \rightarrow \ell\ell$  as normalization and control channel.
- Electrons are more challenging than muons, due to lower reconstruction efficiency and energy loss due to bremsstrahlung.
- Hadronic uncertainties cancel in the ratio.
- $2.6\sigma$  deviation from the SM,  
 $\mathcal{B}(B^+ \rightarrow K^+ e^+ e^-)$  compatible with SM predictions.

# Lepton Flavour universality in loop decays



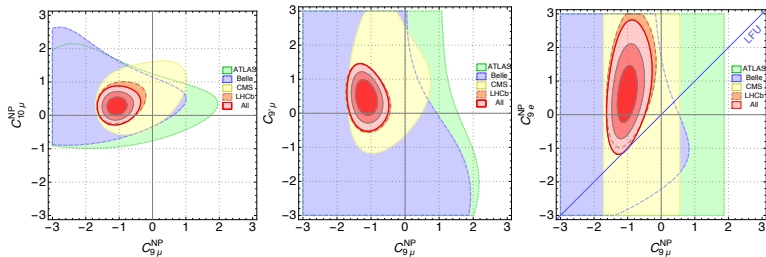
- $2.6\sigma$  from the standard model prediction in  $1 \text{ GeV}^2/c^4 < q^2 < 6 \text{ GeV}^2/c^4$
- $\mathcal{B}(B^+ \rightarrow K^+ e^+ e^-)$  alone is measured to be compatible with the SM.
- Hm...

# Lepton Flavour universality in loop decays



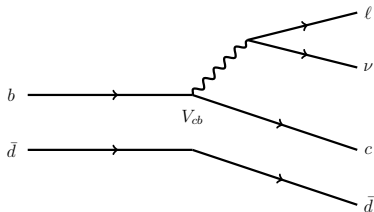
- Consider  $R_{K^*} = \frac{\mathcal{B}(B^0 \rightarrow K^{*0} \mu^+ \mu^-)}{\mathcal{B}(B^0 \rightarrow K^{*0} e^+ e^-)}$
- Similar strategy as for  $R_K$ . Use  $B^0 \rightarrow J/\psi K^*$ ,  $J/\psi \rightarrow \ell\ell$  as normalization and control channel.
- Compatible at 2.2 and 2.4 $\sigma$  with SM prediction for low and intermediate  $q^2$  region.

## Global fits (part II)



- Can introduce a different Wilson coefficient  $C_9$  for muons and electrons and redo global fits.
- Compatible with only deviations in the muon channels and not in the electron channels.

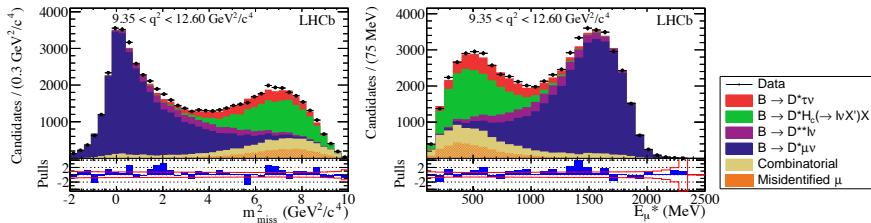
## LFU in $\bar{B}^0 \rightarrow D^{*+} \ell \nu$



- Measure lepton flavour universality in charged-current (tree) decays.
- Measure  $R_{D^*} = \frac{\mathcal{B}(\bar{B}^0 \rightarrow D^{*+} \tau^- \nu)}{\mathcal{B}(\bar{B}^0 \rightarrow D^{*+} \mu^- \nu)}$

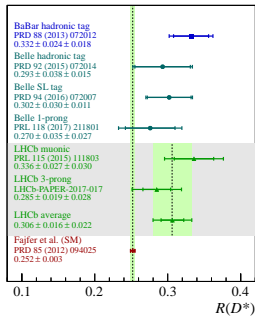
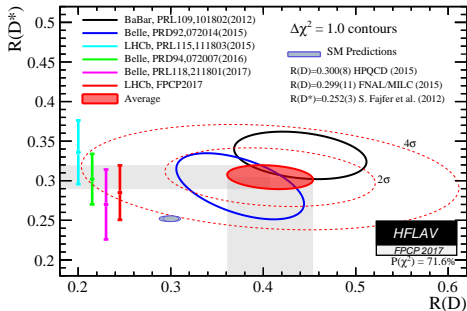


# LFU in $\bar{B}^0 \rightarrow D^{*+} \ell \nu$ , muonic mode



- Measure lepton flavour universality in charged-current (tree) decays.
- Measure  $R_{D^{*}} = \frac{\mathcal{B}(\bar{B}^0 \rightarrow D^{*+} \tau^{-} \nu)}{\mathcal{B}(\bar{B}^0 \rightarrow D^{*+} \mu^{-} \nu)}$
- Use  $\tau^{-} \rightarrow \mu^{-} \nu \nu$ , i.e.  $\tau$  and  $\mu$  modes have the same final state.
  - Distinguish with kinematical distributions
- $R_{D^{*}, \text{exp}, \mu} = 0.336 \pm 0.027(\text{stat}) \pm 0.030(\text{syst})$
- $R_{D^{*}, \text{SM}} = 0.252 \pm 0.003$
- $\approx 2\sigma$  from the SM prediction.

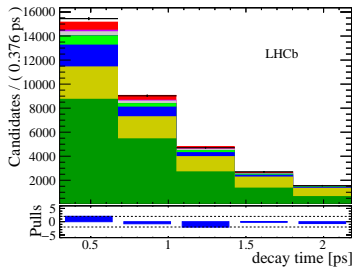
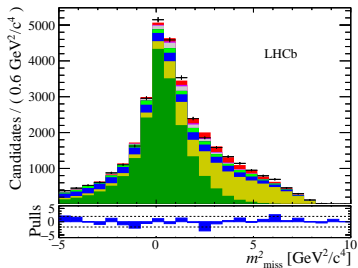
# LFU in $\bar{B}^0 \rightarrow D^{*+} \ell \nu$ , 3-prong mode



- Use  $\tau^- \rightarrow \pi^- \pi^+ \pi^- \nu$
- Use  $\bar{B}^0 \rightarrow D^{*+} 3\pi$  as normalisation channel, and known ratio  $\mathcal{B}(\bar{B}^0 \rightarrow D^{*+} 3\pi) / \mathcal{B}(\bar{B}^0 \rightarrow D^{*+} \mu \nu)$  to calculate  $R_{D^*}$
- $R_{D^*, \text{exp}, 3\pi} = 0.286 \pm 0.019(\text{stat}) \pm 0.025(\text{sys}) \pm 0.021(\text{BR})$



# LFU in $B_c^+ \rightarrow J/\psi \ell \nu$



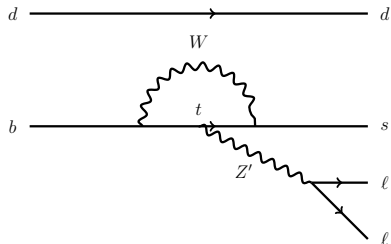
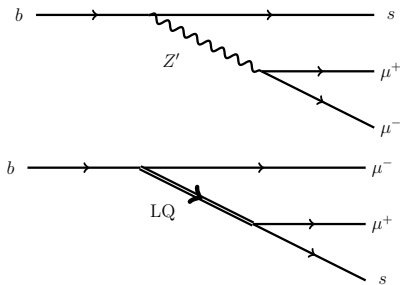
- Measure  $R_{J/\psi} = \frac{B(B_c^+ \rightarrow J/\psi \tau^+ \nu)}{B(B_c^+ \rightarrow J/\psi \mu^+ \nu)}$
- Including first measurement of  $B_c^+ \rightarrow J/\psi \tau^+ \nu$  ( $\tau^- \rightarrow \mu^- \nu \nu$ )
- $R_{J/\psi, \text{theo}} = 0.25 - 0.28$
- $R_{J/\psi, \text{LHCb}} = 0.71 \pm 0.17(\text{stat}) \pm 0.18(\text{syst})$  (compatible within  $2\sigma$ )
- Systematic uncertainties dominated by limited size of simulation and knowledge of  $B_c^+ \rightarrow J/\psi$  form factors.

# Models

- There are a plethora of models that try to explain all anomalies simultaneously.
- Have to explain: Difference in branching fractions, angular distributions, lepton flavour on tree and loop level.
- Possible masses of BSM particles:

tree level, unsuppressed $\sim 30 \text{ TeV}$	loop level, unsuppressed $\sim 2.5 \text{ TeV}$
tree level, MFV $\sim 6 \text{ TeV}$	loop level, MFV $\sim 0.5 \text{ TeV}$

## $Z'$ and Leptoquarks



- Can introduce a  $Z'$  that causes a flavour-changing neutral current on tree level or loop level.
- Leptoquarks can simultaneously explain  $R_{D^{*0}}$  (tree-level leptoquarks),  $R_K$  (loop-level leptoquarks) and muon  $g - 2$  (PRL116, 141802 (2016))

## Optimist's point of view



- Many measurements show a deviation, and when combined, it is significant.
- The pattern is somewhat consistent, as shown by global fits.
- The effects are observed by several (independent) measurements and experiments.
- No large uncertainty in the theoretical prediction has been discovered.



## Pessimist's point of view

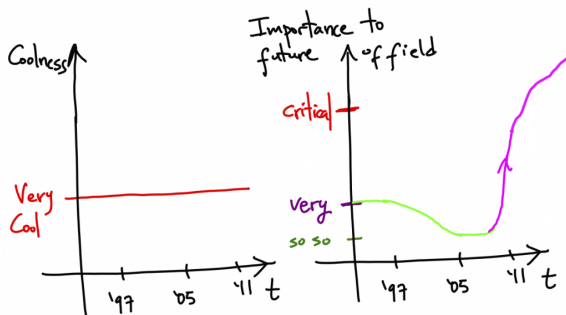


- "The effect is of a magnitude that remains close to the limit of detectability, or many measurements are necessary because of the very low statistical significance of the results."

*I. Langmuir on pathological science.*



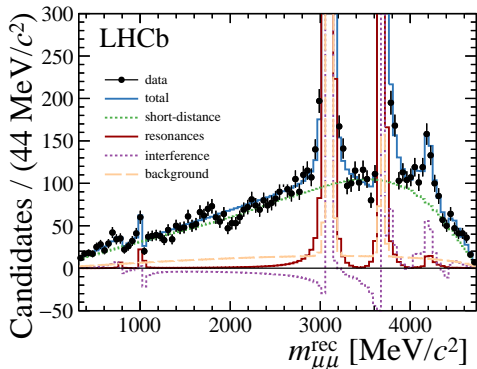
## Conclusion



- B decays are an exciting field of research.
- Several intriguing deviations from the SM have shown up in flavour-changing neutral and charged currents.
- The combination could hint to a deviation from the Standard Model.
- Nature of these anomalies will hopefully soon be resolved.

**Backup**

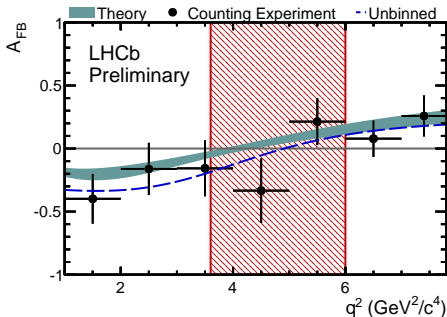
## Short- and long-distance effects in $B \rightarrow K \mu^+ \mu^-$



- Need to perform same analysis for  $B^0 \rightarrow K^{*0} \mu^+ \mu^-$  to understand effect of charm-loops.

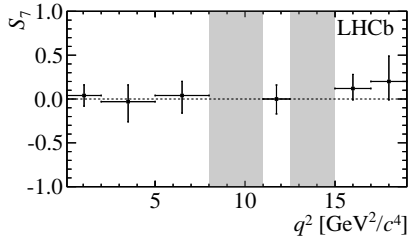
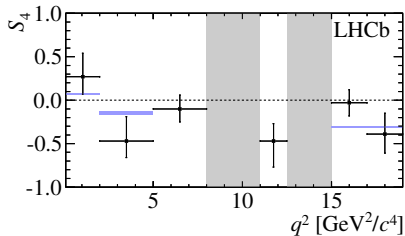
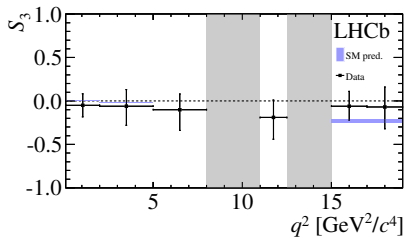
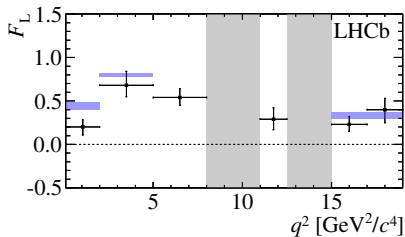


# Zero-crossing point in $B^0 \rightarrow K^{*0} \mu^+ \mu^-$

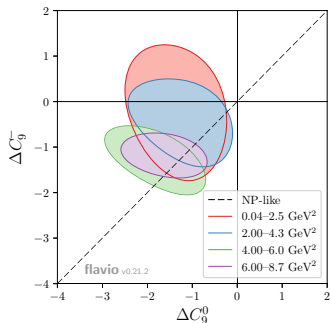
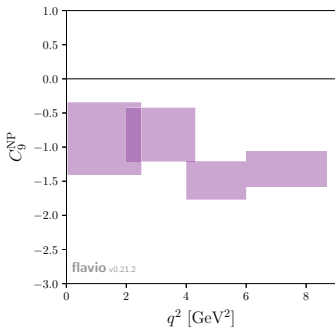


- Zero-crossing point of  $A_{FB}$  theoretically clean.
- For example:  $S_6 = \frac{4}{3} A_{FB} = \frac{4}{3} \frac{\# \cos \theta_\ell > 0 - \# \cos \theta_\ell < 0}{\# \cos \theta_\ell > 0 + \# \cos \theta_\ell < 0}$ .  
Forward-backward asymmetry of the leptons.

# Angular analysis of $B_s^0 \rightarrow \phi \mu^+ \mu^-$

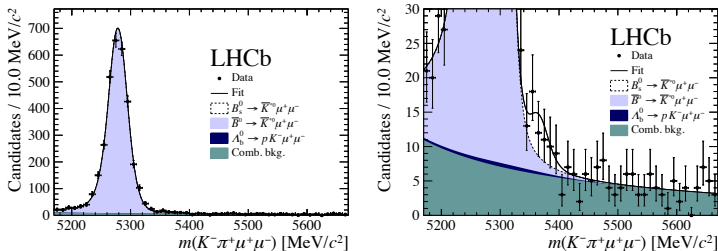


## $q^2$ dependence of $C_9$



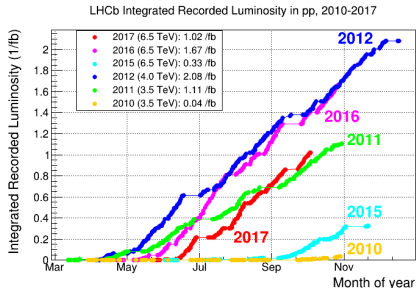
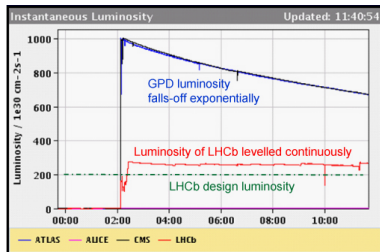
- Check shift of  $C_9$  as a function of  $q^2$ . Should be constant (assuming the new physics is heavy enough).
- That's a hint, but not a confirmation for non-hadronic BSM effects.

# Analyses of $b \rightarrow d\ell\ell$ transitions



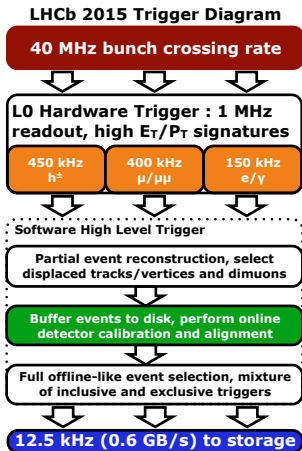
- First evidence of  $B_s^0 \rightarrow K^*\mu^+\mu^-$  ( $3.4\sigma$ )
- Using  $3\text{ fb}^{-1}$  of Run I and  $1.6\text{ fb}^{-1}$  of Run II.
- $\mathcal{B}(B_s^0 \rightarrow K^*\mu^+\mu^-) = (3.0 \pm 1.0(\text{stat}) \pm 0.2(\text{syst}) \pm 0.3(\text{ext})) \cdot 10^{-8}$
- With the upgrade of LHCb from 2021, differential decay rates can be measured.

# Luminosity levelling



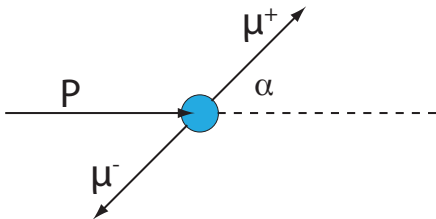
- LHCb does not run at the maximum instantaneous luminosity, as the occupancy in the forward region would be too high.
  - $\mu \approx 1.1$  for 25ns running.
- Luminosity for LHCb is leveled such that it is constant within a fill.
- Achieved by displacing the beams.

## Trigger (Run II)



- Have the same reconstruction (charged and neutral particles) in the software trigger and offline.
- Perform a alignment & calibration after first stage of software trigger, *i.e.* automatically.

## A toy angular analysis

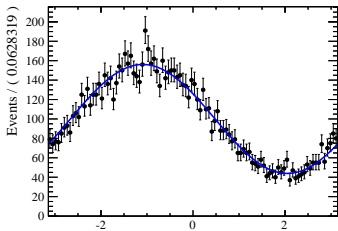


- A particle decays into two particles, with angle  $\alpha$ .
- Suppose we can formulate the angular distribution as:

$$\frac{d\Gamma}{d\alpha} = \frac{1}{2\pi} [A \cos \alpha + B \sin \alpha + C] \quad \alpha \in [-\pi, \pi]$$

- The angular terms are given by kinematics / spin only.
- Remember:  $\frac{d\sigma}{d\Omega}(e^+e^- \rightarrow \mu^+\mu^-) = \frac{\alpha^2}{4s} (1 + \cos^2 \theta)$

## A toy angular analysis



$$\frac{d\Gamma}{d\alpha} = \frac{1}{2\pi} [A \cos \alpha + B \sin \alpha + C] \quad \alpha \in [-\pi, \pi]$$

- The coefficients contain the physics-information we are interested in.
- Do: Run an experiment, collect data, select your decay, plot number of events as a function of  $\alpha$ .
- Fit the angular distribution in collision data with the pdf and extract the coefficients.