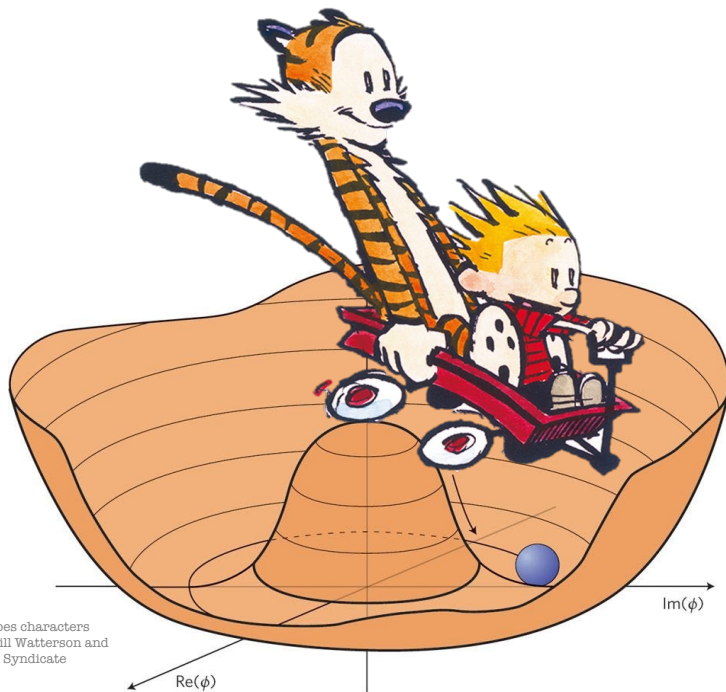


Living with a *pre-teen* Higgs boson

(12 years with the Higgs boson, and the next 12 years)

A tale in 12 chapters



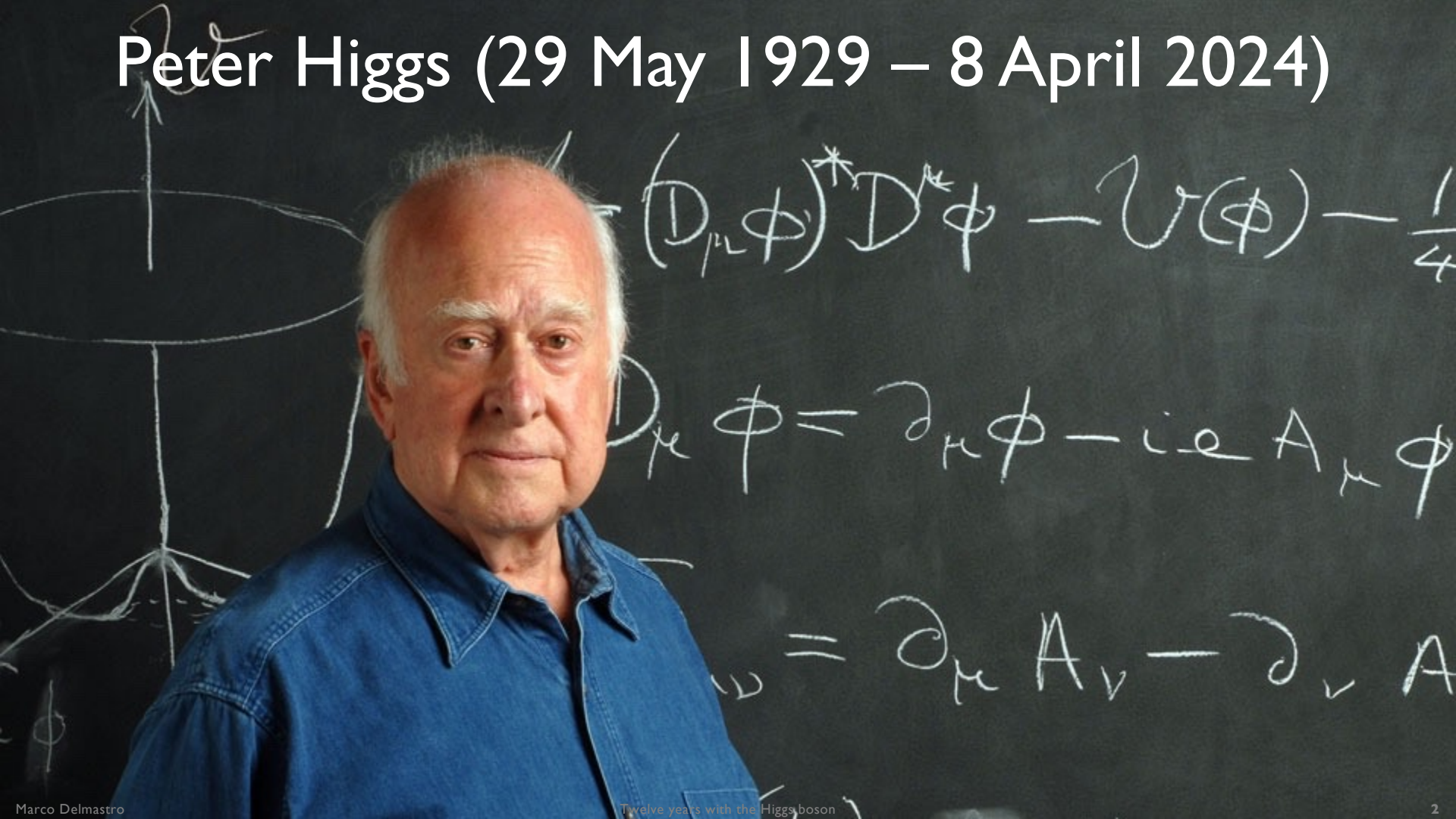
Marco Delmastro



Laboratoire d'Annecy de Physique des Particules

Université de Genève, 8/5/2024

Peter Higgs (29 May 1929 – 8 April 2024)



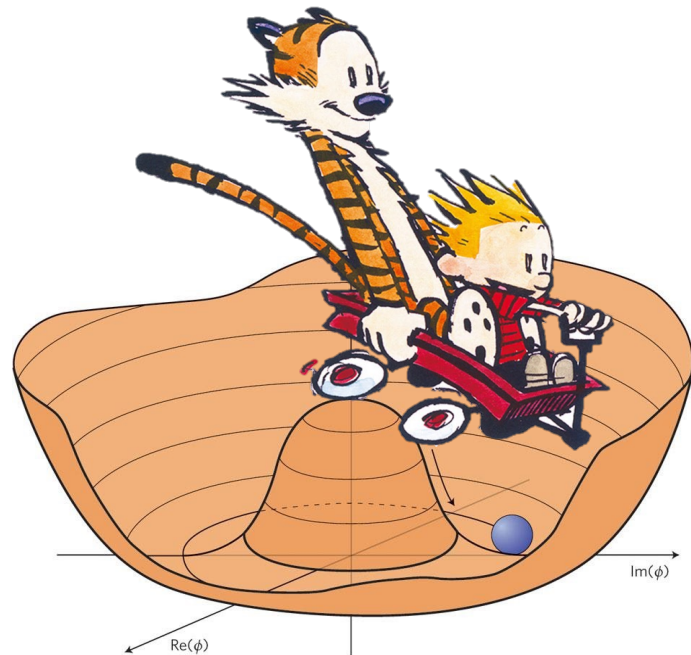
$$(D_\mu \phi)^\dagger D^\mu \phi - \mathcal{V}(\phi) - \frac{1}{4}$$

$$D_\mu \phi = \partial_\mu \phi - ie A_\mu \phi$$

$$F_{\mu\nu} = \partial_\mu A_\nu - \partial_\nu A_\mu$$

1.

Why the Higgs boson?



A fundamental question



How do
elementary
particles
get mass?

This is not a QFT lecture...

All microscopic phenomena seems to be well described by a remarkably simple *theory* (that we continue to call “model” only for historical reasons...)

$$\mathcal{L}_{\text{SM}} = \mathcal{L}_{\text{gauge}}(\psi_i, A_a) + \mathcal{L}_{\text{Higgs}}(H, A_a, \psi_i)$$

massless particles interact by exchanges of forces carried by *massless* gauge bosons. QFT describes *free fields*, interactions generated by *gauge symmetries* (thus *renormalizable*)

mass terms cannot be directly added without breaking theory gauge invariance (i.e. mass is not a fundamental property, rather an emerging phenomenon)

	mass → charge → spin →	≈2.3 MeV/c ² 2/3 1/2 u up	≈1.275 GeV/c ² 2/3 1/2 c charm	≈173.07 GeV/c ² 2/3 1/2 t top	0 0 1 g gluon	≈126 GeV/c ² 0 0 H Higgs boson
QUARKS		≈4.8 MeV/c ² -1/3 1/2 d down	≈95 MeV/c ² -1/3 1/2 s strange	≈4.18 GeV/c ² -1/3 1/2 b bottom	0 0 1 γ photon	
		0.511 MeV/c ² -1 1/2 e electron	105.7 MeV/c ² -1 1/2 μ muon	1.777 GeV/c ² -1 1/2 τ tau	0 0 1 Z Z boson	
LEPTONS		<2.2 eV/c ² 0 1/2 ν_e electron neutrino	<0.17 MeV/c ² 0 1/2 ν_μ muon neutrino	<15.5 MeV/c ² 0 1/2 ν_τ tau neutrino	±1 1 W W boson	GAUGE BOSONS

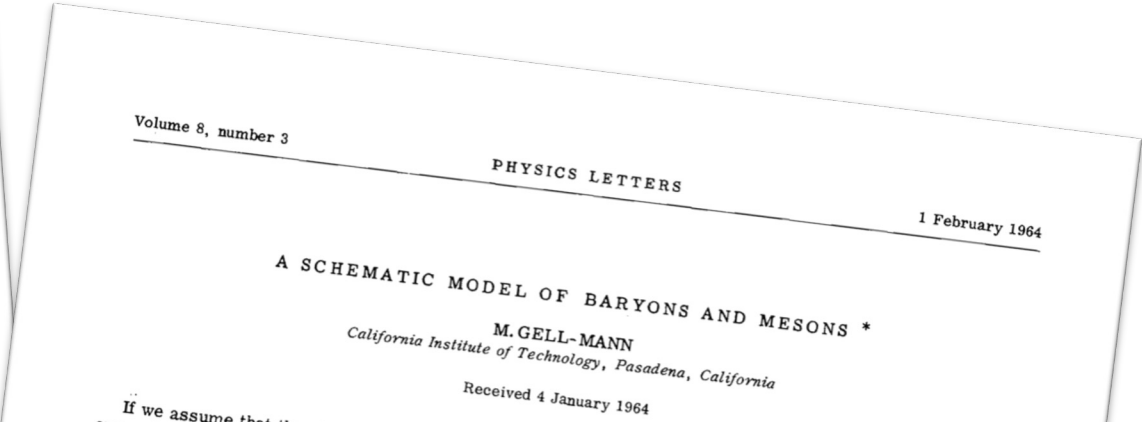
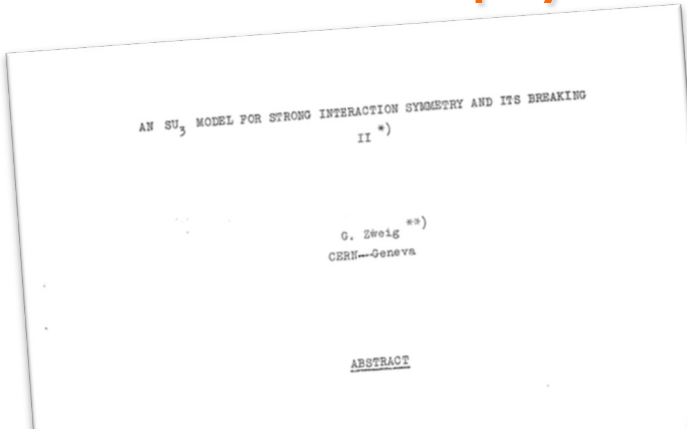
Theory would not make sense without a *Higgs field* (or something similar) responsible for *spontaneous symmetry breaking* and *mass generation*...

2.

An idea from 1964



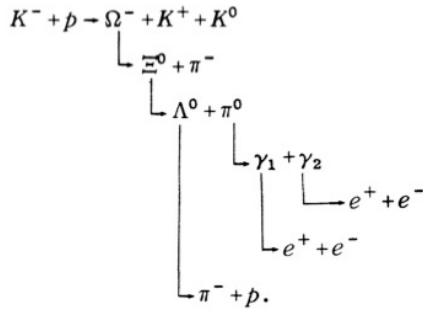
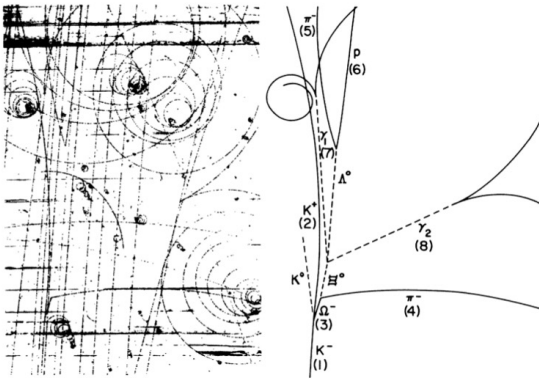
What's hot in physics in 1964?



VOLUME 12, NUMBER 8 PHYSICAL REVIEW LETTERS 24 FEBRUARY 1964

OBSERVATION OF A HYPERON WITH STRANGENESS MINUS THREE*

Both mesons are three fundamental particles into an isospin doublet (number 1/3 and is from Eightfold Way) is also interactions. The breaking is universal, being due to extensive space-time interactions. Predictions for both meson and baryon interactions are presented concerning the rates for the interaction theory to predict rates for experimental search.



For all known baryons and mesons, an example of such a triplet has spin 1/2 and baryons d^-, s^-, u^0 and b^0 mesons.

The same scheme can be used for integral values for the baryon number. The basic triplet has the following baryon number 1/3: $u^2/3, d^{-1/3},$ and $s^{-2/3}$ of the members of the baryon octet. The combinations of the members of the octet are now being made out. The lowest energy combination observed, while the quark model (q) similarly gives

*) Version I is...
**) This work was supported by the Office of Naval Research and the National Science Foundation.

8419/TH.412
21 February 1964

FIG. 2. Photograph and line diagram of event showing decay of Ω^- .

EVIDENCE FOR THE 2π DECAY OF THE K_2^0 MESON*†

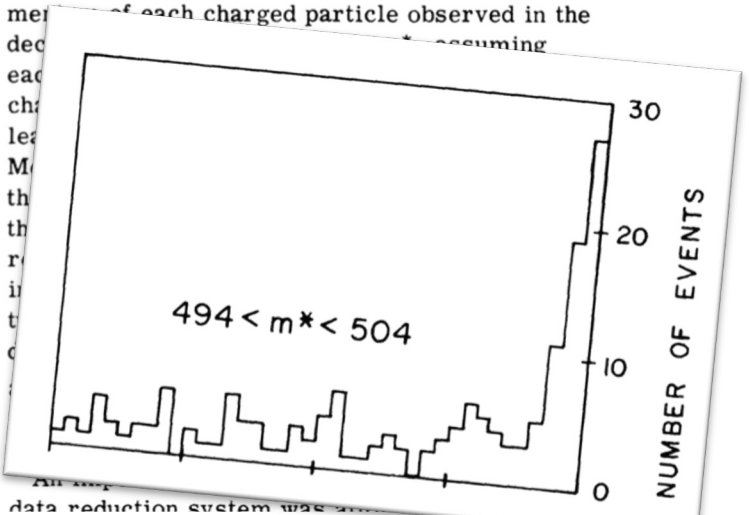
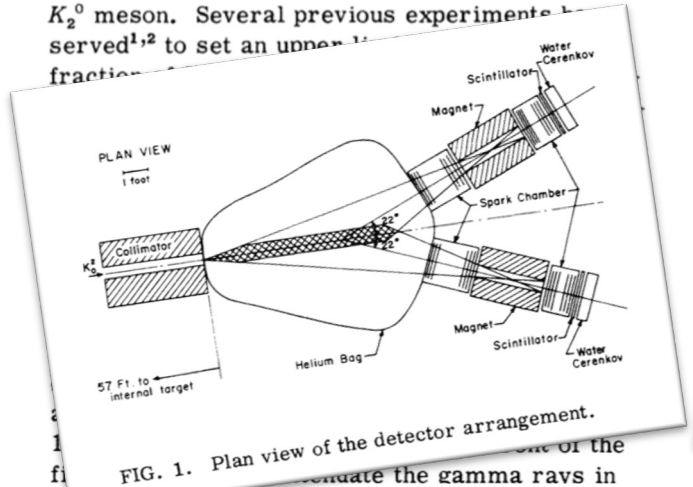
J. H. Christenson, J. W. Cronin,‡ V. L. Fitch,‡ and R. Turlay§

Princeton University, Princeton, New Jersey

(Received 10 July 1964)

This Letter reports the results of experimental studies designed to search for the 2π decay of the K_2^0 meson. Several previous experiments have served^{1,2} to set an upper limit on the branching fraction.

The analysis program computed the vector momenta of each charged particle observed in the decay, assuming



The Brout-Englert-Higgs mechanism is also proposed...

... and barely noticed!

BROKEN SYMMETRY AND THE MASS OF GAUGE VECTOR MESONS*
F. Englert and R. Brout
Faculté des Sciences, Université Libre de Bruxelles, Bruxelles, Belgium
(Received 26 June 1964)

14 citations in the first 5 years

BROKEN SYMMETRIES, MASSLESS PARTICLES AND GAUGE FIELDS
P. W. HIGGS
Tait Institute of Mathematical Physics, University of Edinburgh, Scotland
Received 27 July 1964

18 citations in the first 5 years

VOLUME 13, NUMBER 16 PHYSICAL REVIEW LETTERS 19 OCTOBER 1964

BROKEN SYMMETRIES AND THE MASSES OF GAUGE BOSONS
Peter W. Higgs
Tait Institute of Mathematical Physics, University of Edinburgh, Edinburgh, Scotland
(Received 31 August 1964)

11 citations in the first 5 years

GLOBAL CONSERVATION LAWS AND MASSLESS PARTICLES*
G. S. Guralnik,† C. R. Hagen,‡ and T. W. B. Kibble
Department of Physics, Imperial College, London, England
(Received 12 October 1964)

23 citations in the first 5 years

It would take 10 more years to become “mainstream”...

PHYSICAL REVIEW

VOLUME 155, NUMBER 5

25 MARCH 1967

1967

Symmetry Breaking in Non-Abelian Gauge Theories*

T. W. B. KIBBLE

Department of Physics, Imperial College, London, England

(Received 24 October 1966)

Photon remains massless

1967

Electroweak Symmetry Breaking
incorporated into
Standard Model

A MODEL OF LEPTONS*

Steven Weinberg†

Laboratory for Nuclear Science and Physics Department,
Massachusetts Institute of Technology, Cambridge, Massachusetts

(Received 17 October 1967)

A. Salam, “Weak And Electromagnetic Interactions,” Originally printed in “N. Svartholm: Elementary Particle Theory, Proceedings of The Nobel Symposium Held 1968 at Lerum, Sweden” 367-377 (1968).

RENORMALIZABLE LAGRANGIANS FOR MASSIVE YANG-MILLS FIELDS

G. 't HOOFT

Institute for Theoretical Physics, University of Utrecht

Received 13 July 1971

REGULARIZATION AND RENORMALIZATION OF GAUGE FIELDS

G. 't HOOFT and M. VELTMAN

*Institute for Theoretical Physics *, University of Utrecht*

Received 21 February 1972

1971-72

Standard Model is renormalizable

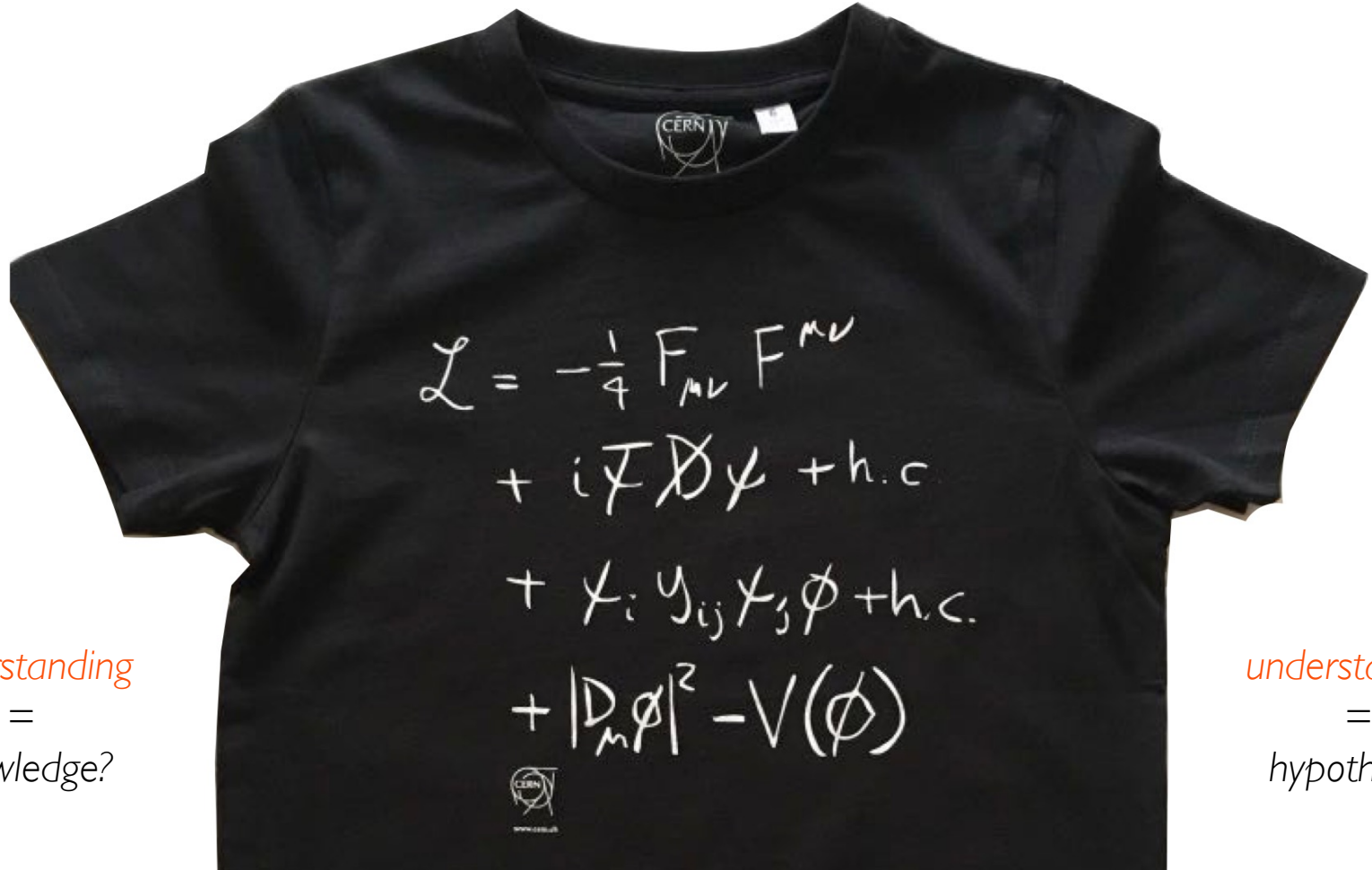
3.



A closer look to the SM Lagrangian



“This equation neatly sums up our current **understanding** of fundamental particles and forces”



understanding
=
knowledge?

understanding
=
hypothesis?

Gauge sector


This part is very well known (one could even say since the discovery of Maxwell equations in 1880), elegant and very well experimentally verified!



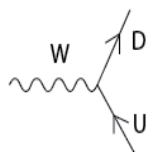
X is any fermion in the Standard Model.



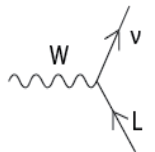
X is electrically charged.



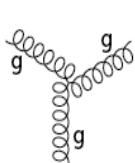
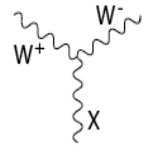
X is any quark.



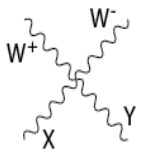
U is a up-type quark;
D is a down-type quark.



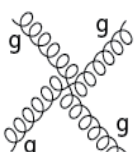
L is a lepton and nu is the corresponding neutrino.

X is a photon or Z-boson.



X and Y are any two electroweak bosons such that charge is conserved.



$$\mathcal{L} = -\frac{1}{4} F_{\mu\nu} F^{\mu\nu} + i\bar{\psi} \not{D} \psi$$

$$+ \chi_i y_{ij} \chi_j \phi + h.c.$$

$$+ |D_\mu \phi|^2 - V(\phi)$$

This equation neatly sums up our current understanding of fundamental particles and forces.

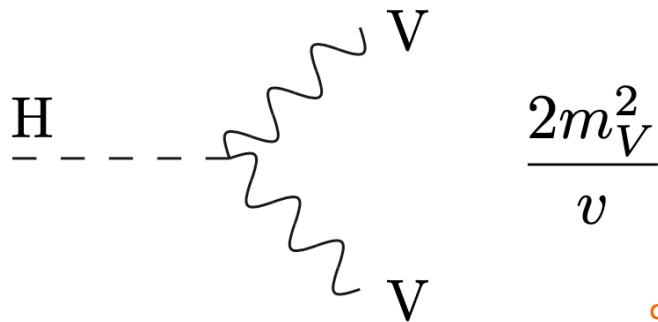
understanding = knowledge

Electro-weak symmetry breaking

Still a gauge interaction (nothing new, one might be tempted to say...), only this time **between bosons and a scalar field!**

Nothing like this directly observed before 2012 (but tested before by precision electroweak interactions)

Experimentally observed in 2012 with Higgs boson discovery...



...this term could not exist without a v.e.v.

$$\begin{aligned} \mathcal{L} = & -\frac{1}{4} F_{\mu\nu} F^{\mu\nu} \\ & + i\bar{\psi} \not{D} \psi \\ & + \chi_i Y_{ij} \chi_j \phi + \text{h.c.} \\ & + |D_\mu \phi|^2 - V(\phi) \end{aligned}$$

This equation neatly sums up our current understanding of fundamental particles and forces.

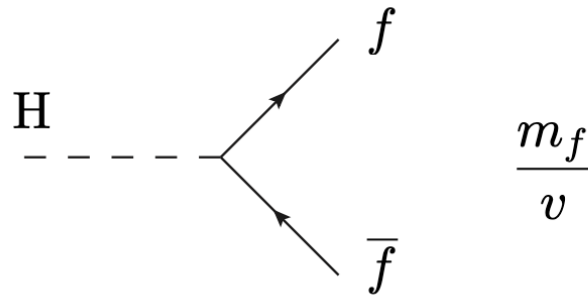
Fermion masses via Yukawa interaction

Before 2012 this term was almost a conjecture... and even after the discovery in 2012, it remains an “anomalous” term (it’s *not a gauge interaction*, and no interaction with this structure has *ever been observed in nature*)

Is the Higgs boson responsible for the EW symmetry breaking also responsible for the masses of fermions?

Is the Higgs boson responsible for the masses of all fermions?

understanding = hypothesis



$$\begin{aligned} \mathcal{L} = & -\frac{1}{4} F_{\mu\nu} F^{\mu\nu} \\ & + i\bar{\psi} \not{D} \psi \\ & + \chi_i Y_{ij} \chi_j \phi + \text{h.c.} \\ & + |D_\mu \phi|^2 - V(\phi) \end{aligned}$$

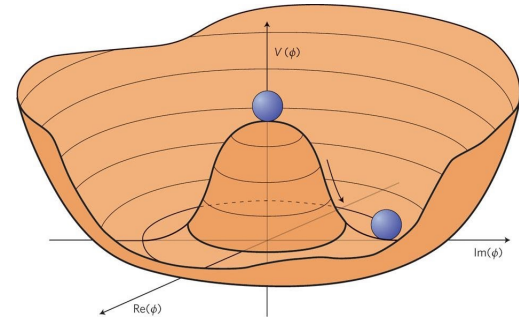
This equation neatly sums up our current understanding of fundamental particles and forces.

$$\mathcal{L} = -\frac{1}{4} F_{\mu\nu} F^{\mu\nu} + i\bar{\psi}\not{D}\psi + \chi_i Y_{ij} \chi_j \phi + h.c. + |D_\mu \phi|^2 - V(\phi)$$

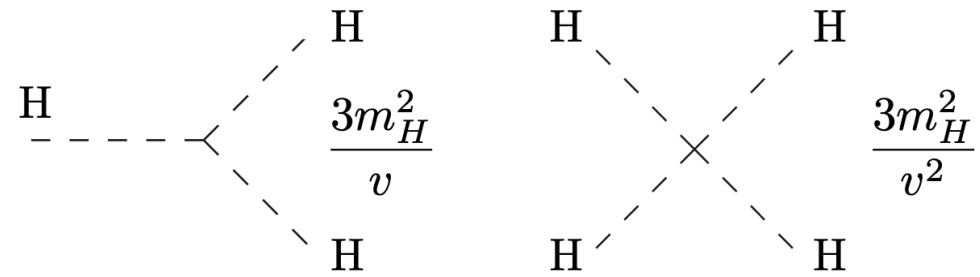
Higgs potential

$$V(\Phi) = -\mu^2 \Phi^\dagger \Phi + \lambda (\Phi^\dagger \Phi)^2$$

It defines all characteristics of Higgs field and defines Higgs self-interaction properties. Nothing like this has **ever been observed in nature**



Is the shape of the Higgs potential that postulated by the Standard Model?



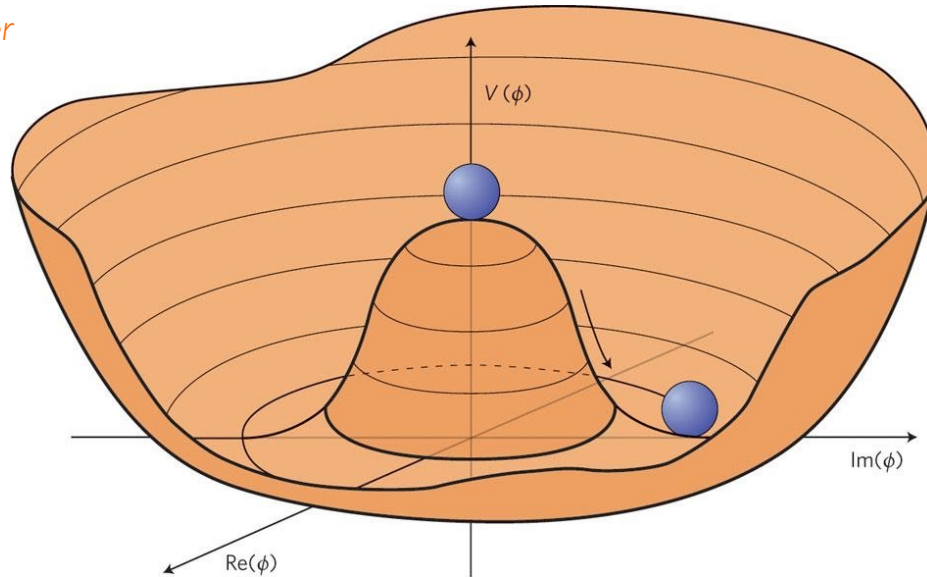
This equation neatly sums up our current understanding of fundamental particles and forces.

A couple of things about the Higgs potential...

$$V(\Phi) = -\mu^2 \Phi^\dagger \Phi + \lambda (\Phi^\dagger \Phi)^2$$

The potential is *symmetric under rotations in Φ space*

It has a minimum which is *not at $\langle \Phi \rangle = 0$* known as the *vacuum expectation value (vev) v* that defines the *electroweak scale*



$$\mu^2 < 0$$

$$\lambda > 0$$

The mass of the Higgs boson is not predicted by the theory!

$$v = \frac{\mu}{\sqrt{\lambda}} = \frac{m_W}{g} = 246 \text{ GeV}$$

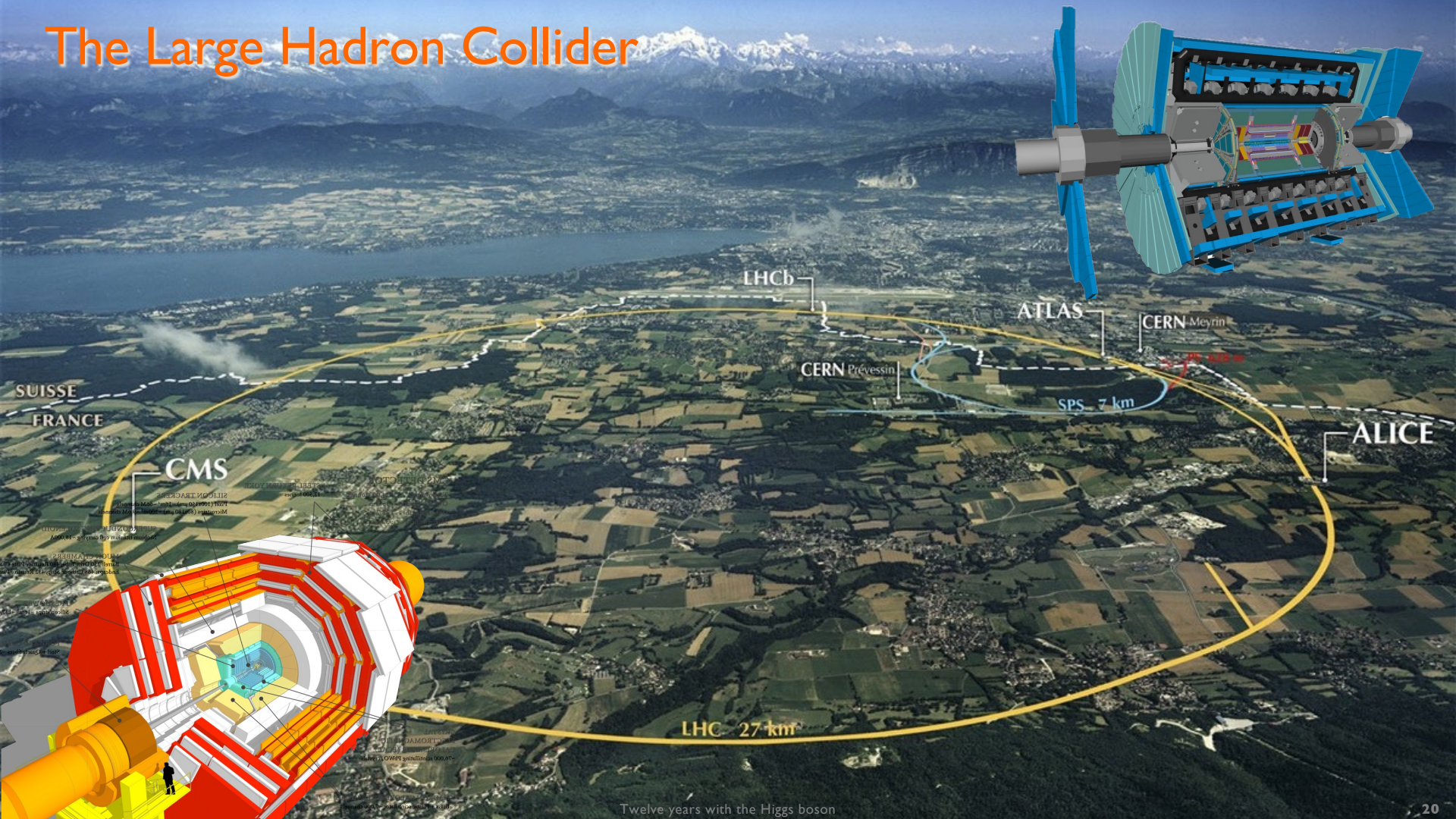
$$m_H = \sqrt{2}\mu = \sqrt{2\lambda}v$$

4.

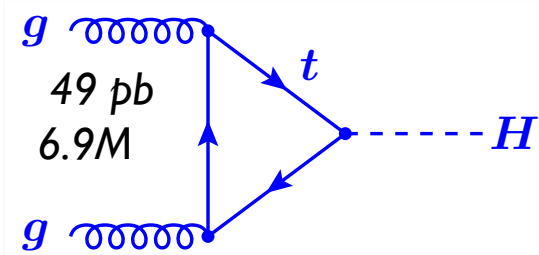


A textbook discovery

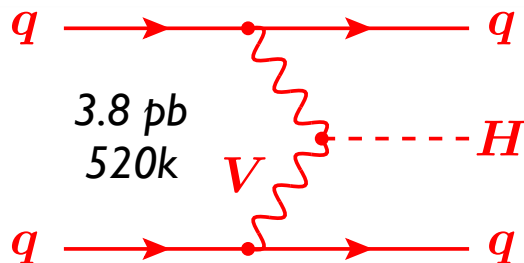
The Large Hadron Collider



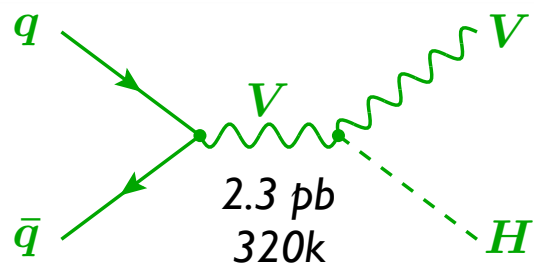
Higgs boson production at the LHC



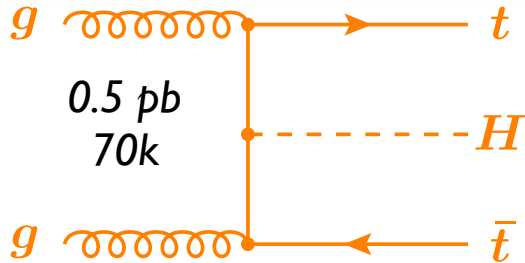
gluon-gluon fusion:
main production mode at
LHC



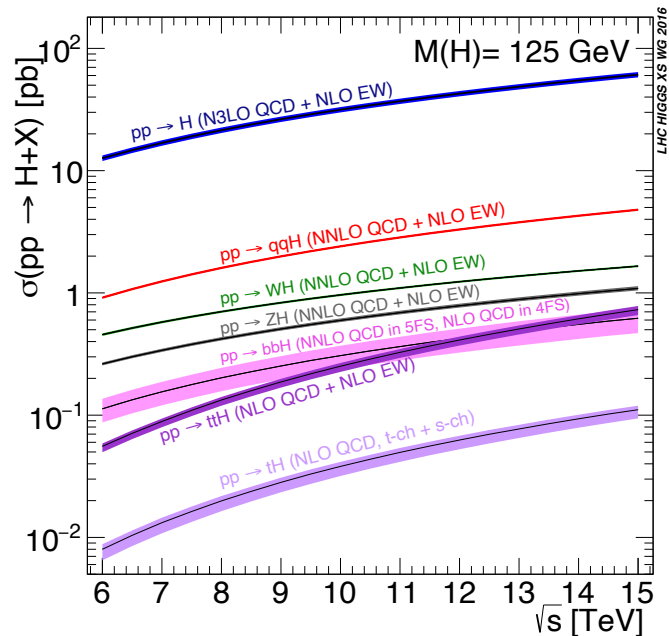
Vector Boson Fusion
2 well-separated forward
jets



VH
tag W and Z
boson decays



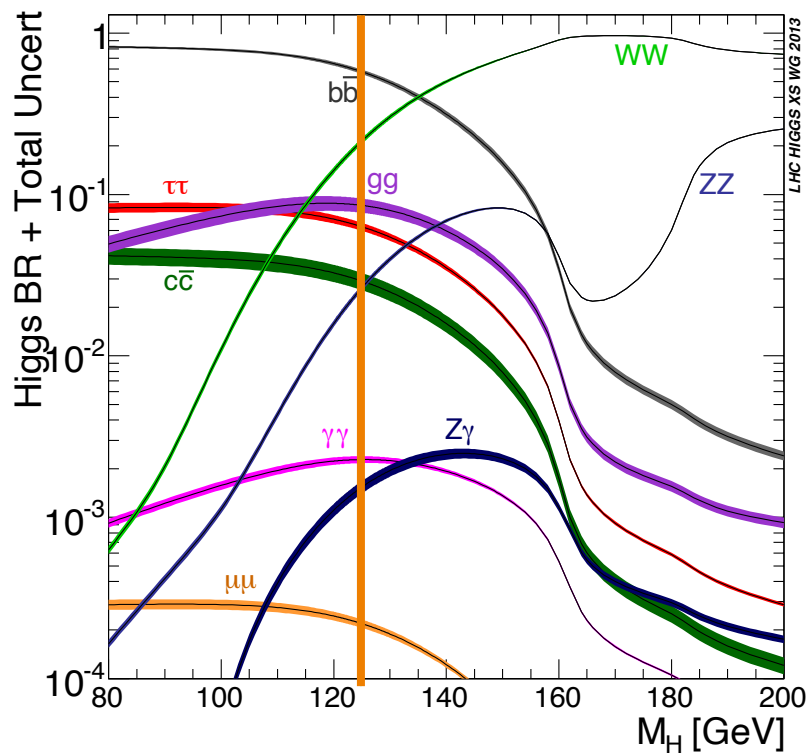
ttH
tag 2 top quarks



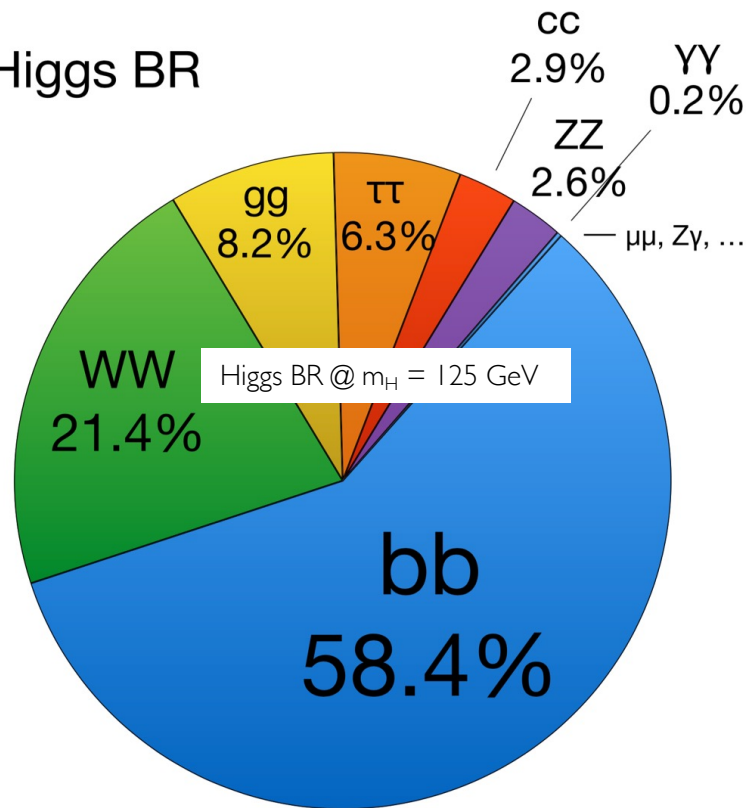
$\sigma[\text{pb}] @ 13 \text{ TeV}$

Higgs produced in 140 fb^{-1} in one experiment

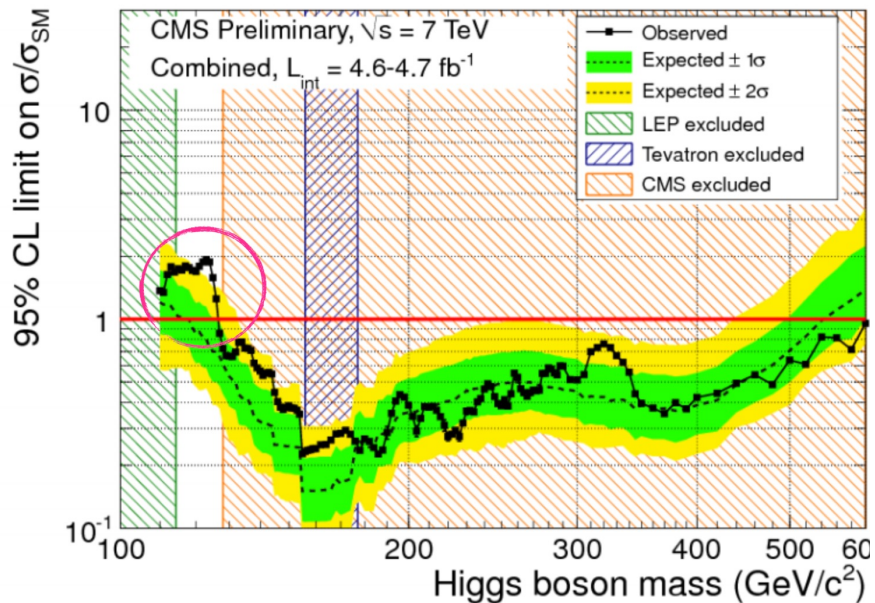
The Higgs boson is an unstable particle



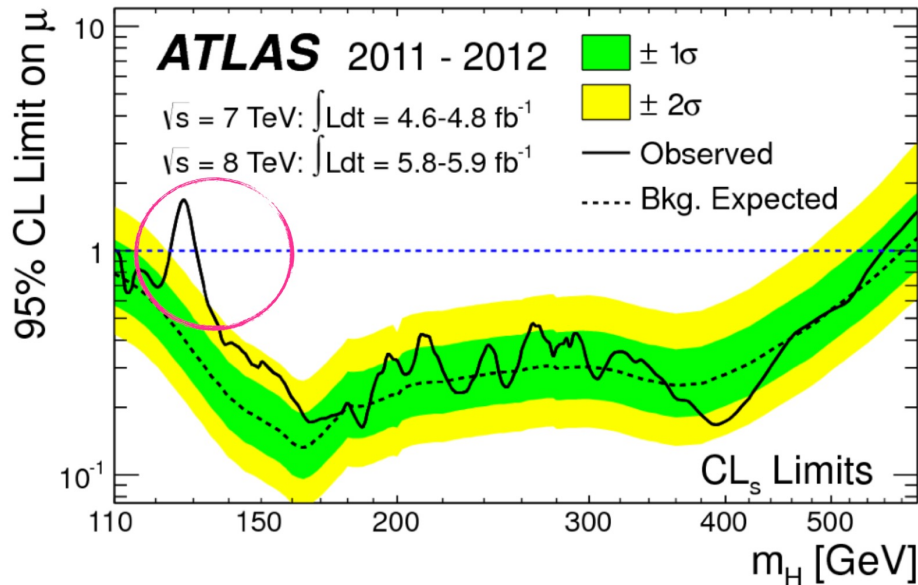
Higgs BR



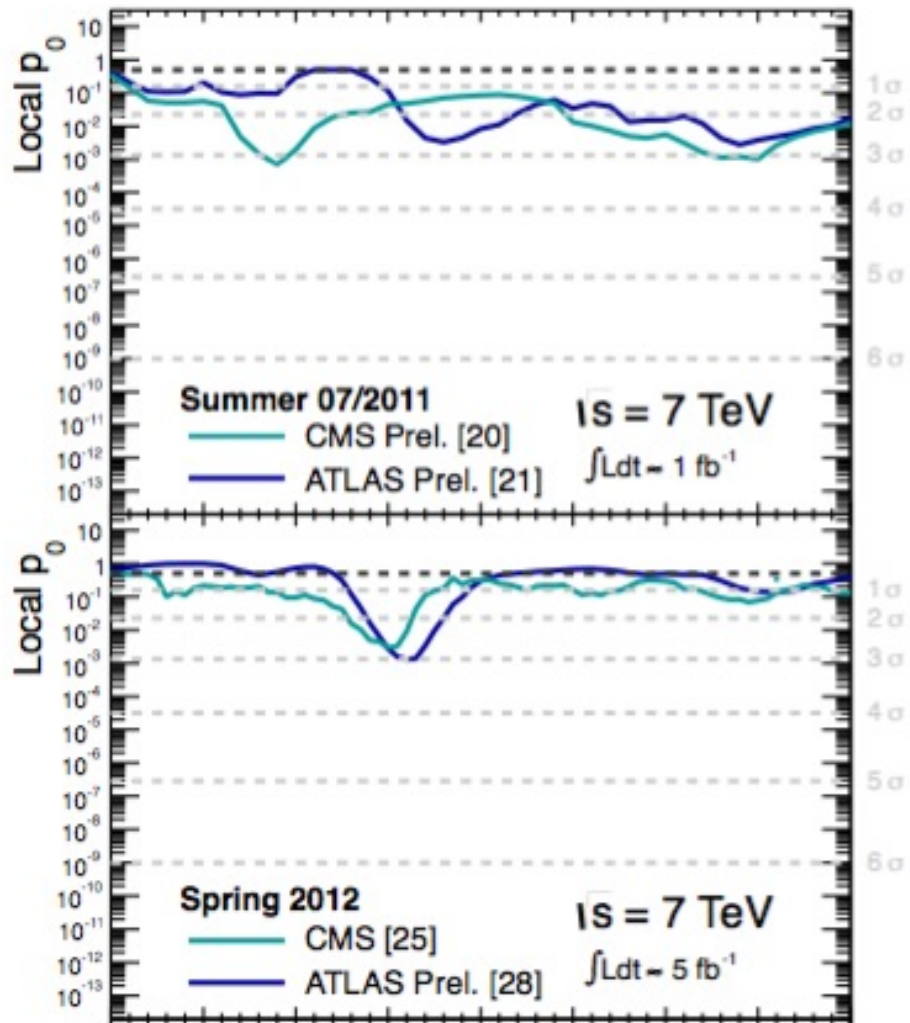
Almost exactly 12 years ago...

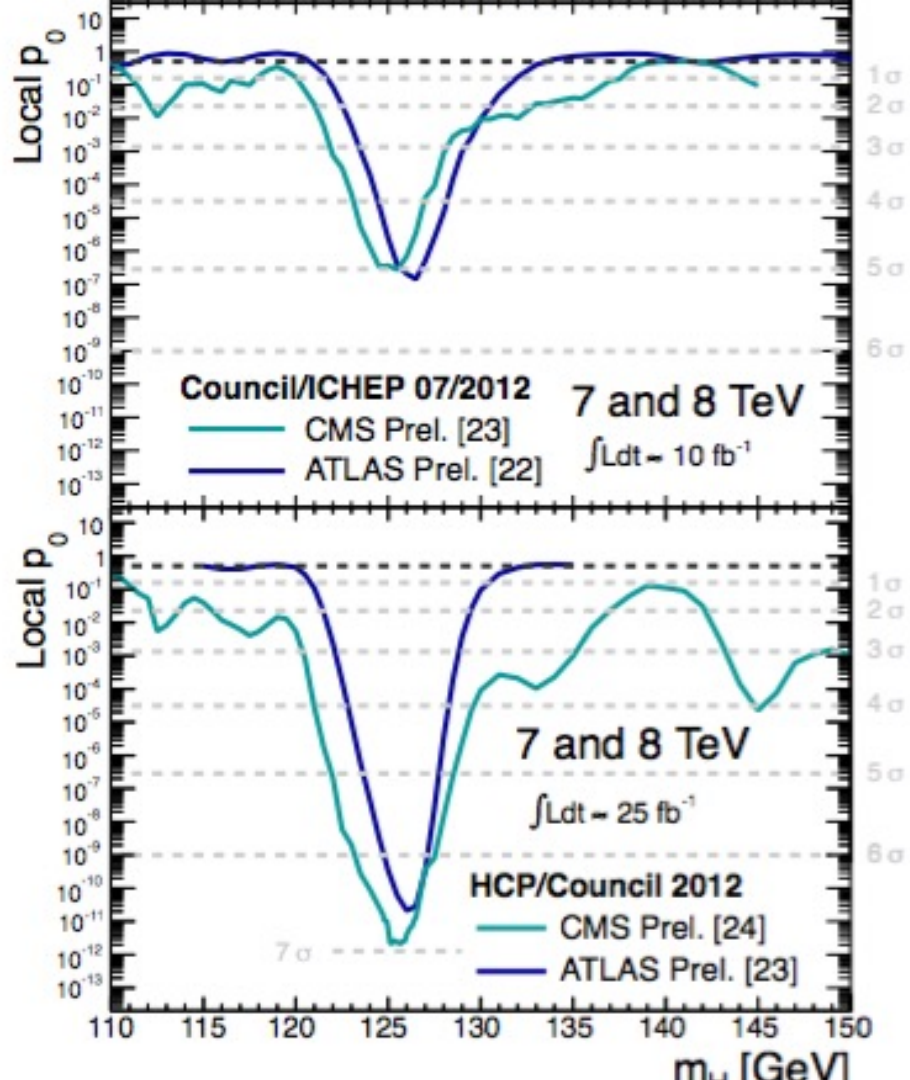


End of 2011...



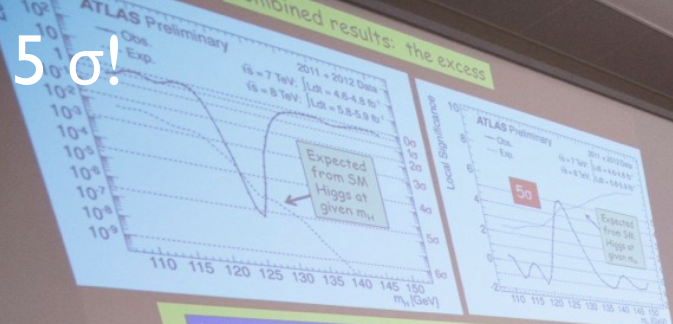
... June 2012...





$5\sigma!$

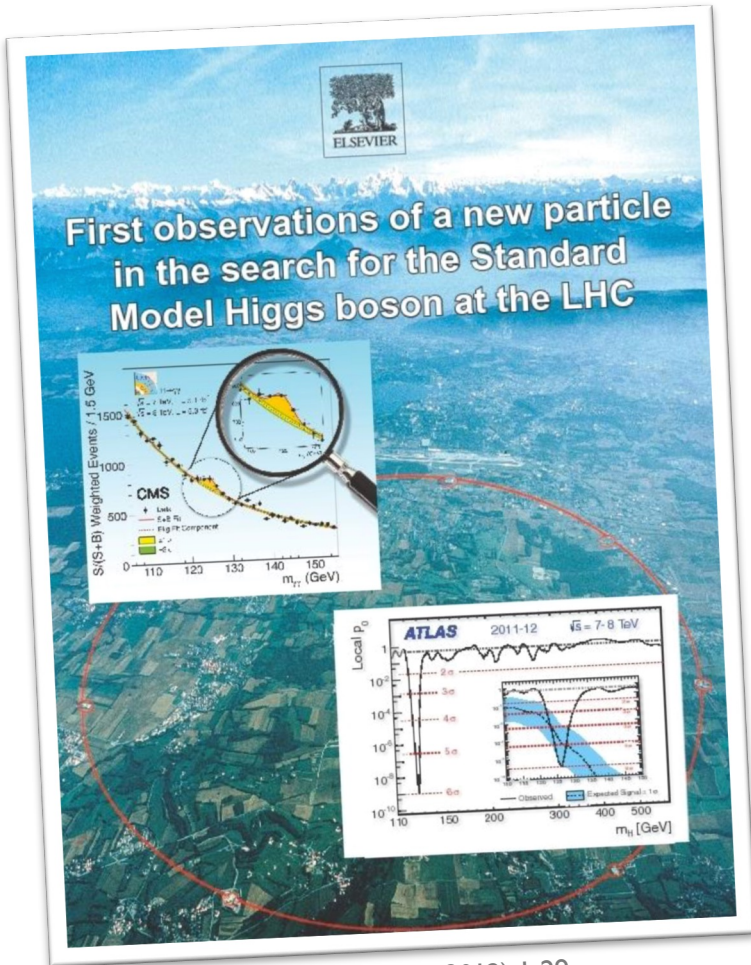
Combined results: the excess



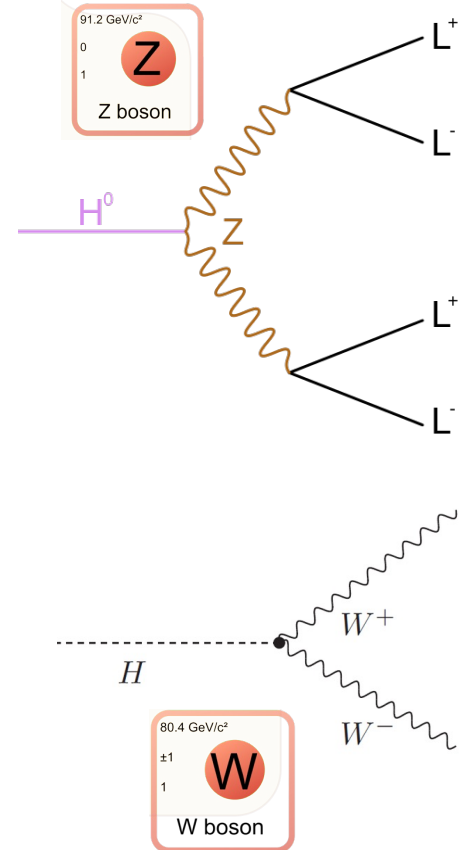
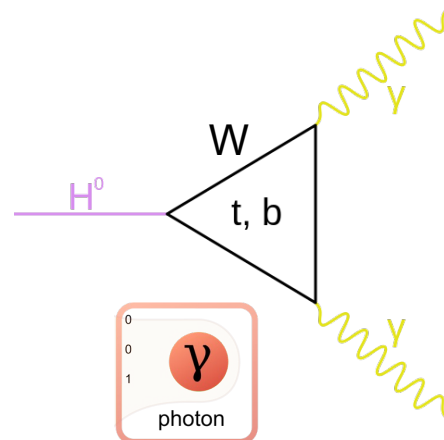
Maximum excess observed at	$m_h = 126.5 \text{ GeV}$
Local significance (including energy-scale systematics)	5.0σ
Probability of background up-fluctuation	3×10^{-7}
Expected from SM Higgs $m_h = 126.5$	4.6σ
Global significance: $4.1\text{-}4.3 \sigma$ (for LEE over 110-600 or 110-180 GeV)	

Fabiola Gianotti revealing the combined significance of the ATLAS Higgs searches on on July 4 2012

Higgs boson discovery channels



ATLAS: PLB 716 (2012) 1-29
 CMS: PLB 716 (2012) 30



5.

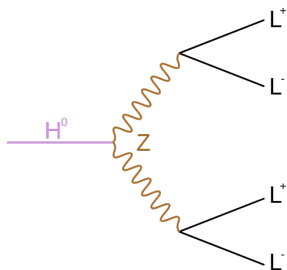
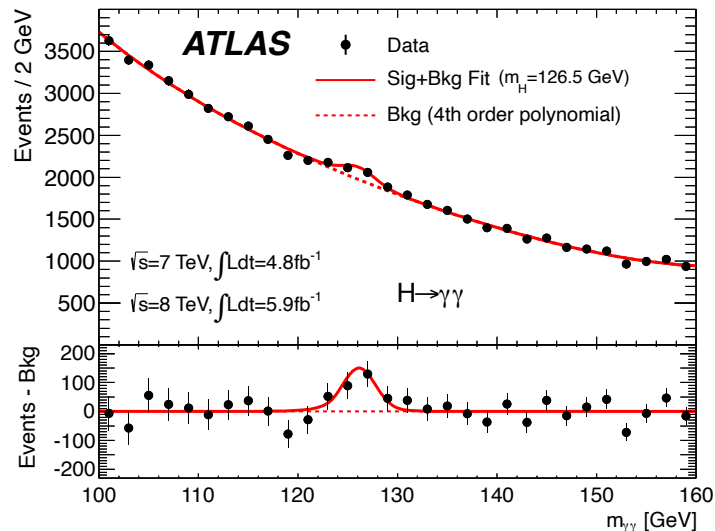
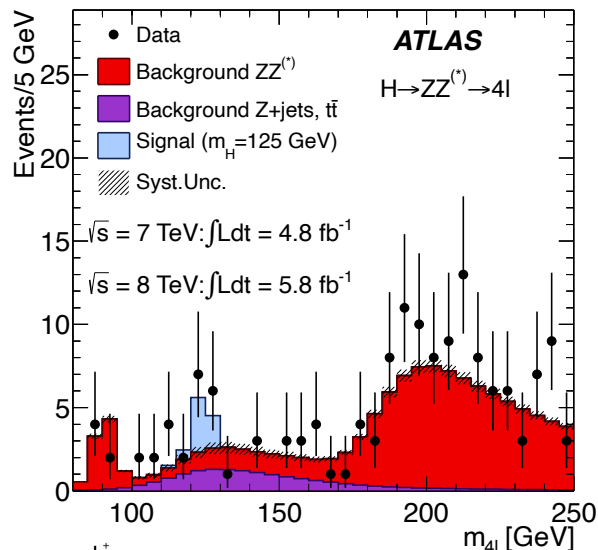
12 years of Higgs measurements



A long way since the discovery...

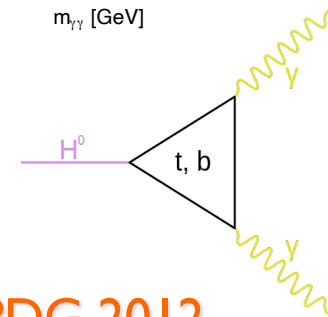
2012

Phys Lett B 716 (1) 1-29 2012



H^0 DECAY MODES	Fraction (Γ_i/Γ)	p (MeV/c)
$W W^*$	seen	—
$Z Z^*$	seen	—
$\gamma\gamma$	seen	—
$b\bar{b}$	possibly seen	—
$\tau^+ \tau^-$	possibly seen	—

Mass Limits for the Standard Model Higgs
 Mass $m > 122$ and none 127–600 GeV, CL = 95%

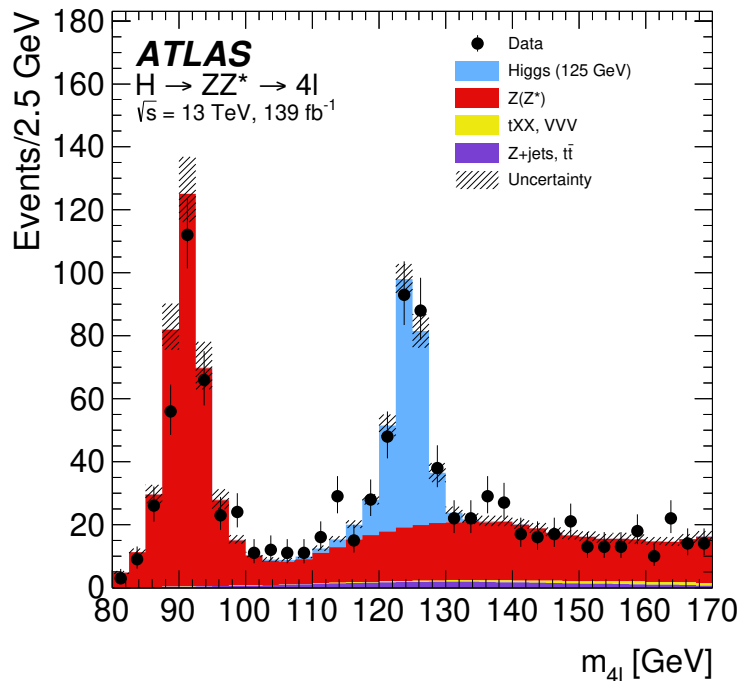


PDG 2012

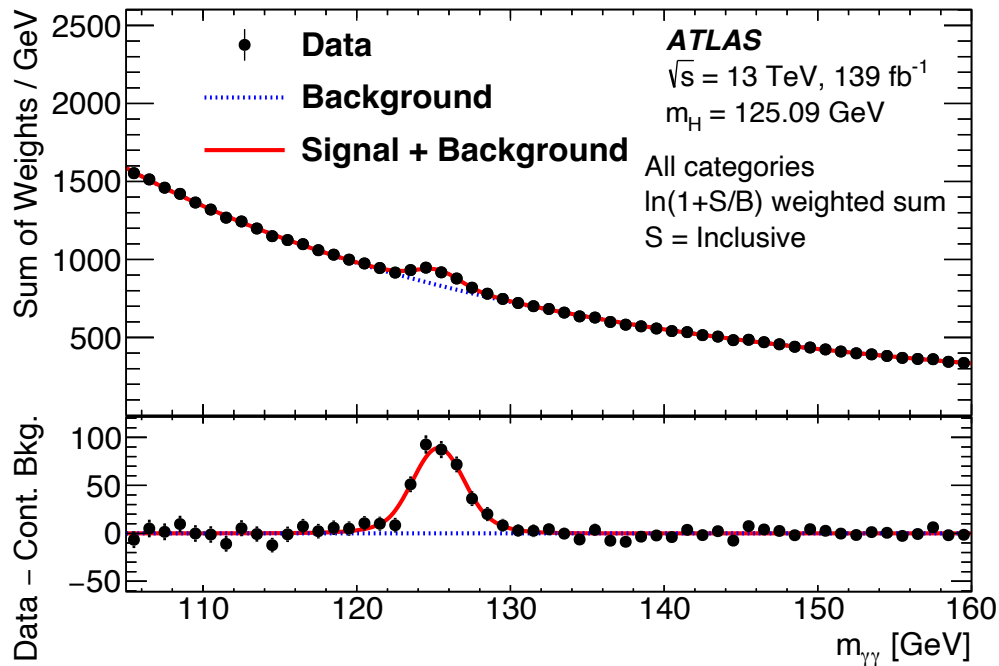
J. Beringer *et al.* (Particle Data Group), PR **D86**, 010001 (2012)

A long way since the discovery...

YESTERDAY



[Eur. Phys. J. C 80 \(2020\) 942](#)



[JHEP 07 \(2023\) 088](#)

A long way since the discovery...

PDG 2022

H^0

$J = 0$

Mass $m = 125.25 \pm 0.17$ GeV ($S = 1.5$)
 Full width $\Gamma = 3.2_{-2.2}^{+2.8}$ MeV (assumes equal
 on-shell and off-shell effective couplings)

H^0 Signal Strengths in Different Channels

Combined Final States = 1.13 ± 0.06

$W W^* = 1.19 \pm 0.12$

$Z Z^* = 1.01 \pm 0.07$

$\gamma\gamma = 1.10 \pm 0.07$

$c\bar{c}$ Final State = 37 ± 20

$b\bar{b} = 0.98 \pm 0.12$

$\mu^+\mu^- = 1.19 \pm 0.34$

$\tau^+\tau^- = 1.15_{-0.15}^{+0.16}$

$Z\gamma < 3.6$, CL = 95%

$\gamma^*\gamma$ Final State = 1.5 ± 0.5

$t\bar{t}H^0$ Production = 1.10 ± 0.18

tH^0 production = 6 ± 4

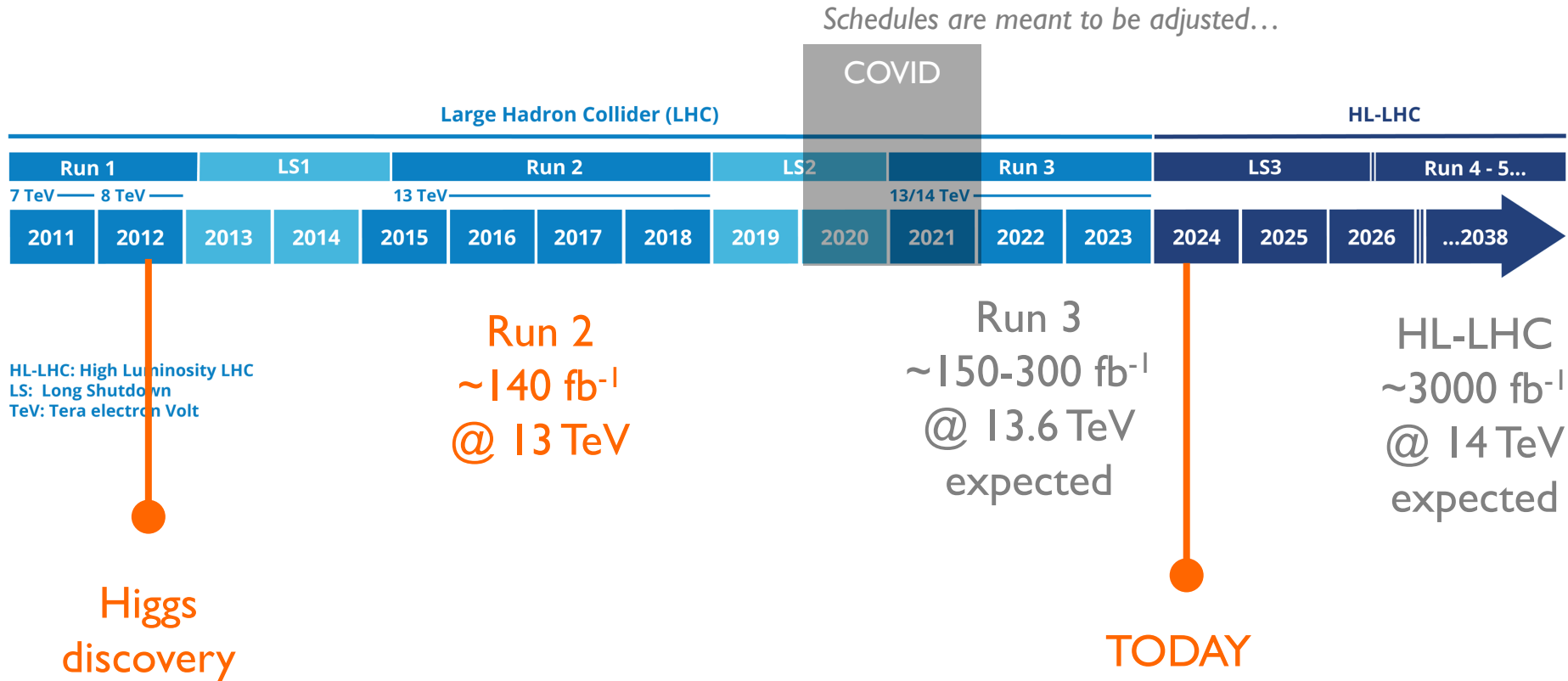
H^0 Production Cross Section in pp Collisions at $\sqrt{s} = 13$ TeV =
 56 ± 4 pb

H^0 DECAY MODES

H^0 DECAY MODES	Fraction (Γ_i/Γ)	Confidence level	ρ (MeV/c)
e^+e^-	$< 3.6 \times 10^{-4}$	95%	62625
$Z\rho(770)$	$< 1.21\%$	95%	29423
$Z\phi(1020)$	$< 3.6 \times 10^{-3}$	95%	29417
$J/\psi\gamma$	$< 3.5 \times 10^{-4}$	95%	62587
$J/\psi J/\psi$	$< 1.8 \times 10^{-3}$	95%	62548
$\psi(2S)\gamma$	$< 2.0 \times 10^{-3}$	95%	62571
$\Upsilon(1S)\gamma$	$< 4.9 \times 10^{-4}$	95%	62268
$\Upsilon(2S)\gamma$	$< 5.9 \times 10^{-4}$	95%	62224
$\Upsilon(3S)\gamma$	$< 5.7 \times 10^{-4}$	95%	62197
$\Upsilon(nS)\Upsilon(mS)$	$< 1.4 \times 10^{-3}$	95%	—
$\rho(770)\gamma$	$< 8.8 \times 10^{-4}$	95%	62623
$\phi(1020)\gamma$	$< 4.8 \times 10^{-4}$	95%	62621
$e\mu$	LF $< 6.1 \times 10^{-5}$	95%	62625
$e\tau$	LF $< 2.2 \times 10^{-3}$	95%	62612
$\mu\tau$	LF $< 1.5 \times 10^{-3}$	95%	62612
invisible	$< 19\%$	95%	—

R.L. Workman *et al* (PDG), Prog. Theor. Exp. Phys. **2022**, 083C01 (2022)

12 years with the Higgs boson, and the next 12+ years...



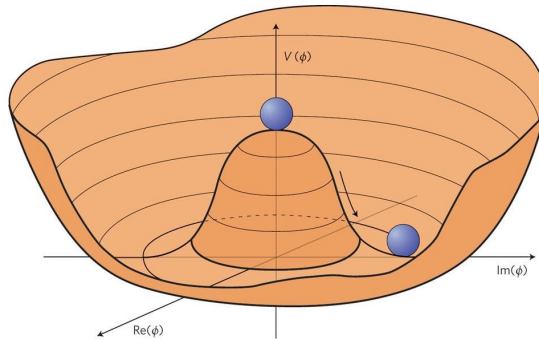
6.

How well
do we know
the “basic”

Higgs properties?



Higgs boson mass and width



$$V(\Phi) = -\mu^2 \Phi^\dagger \Phi + \lambda (\Phi^\dagger \Phi)^2$$

$$\nu = \frac{\mu}{\sqrt{\lambda}}$$

$$V(\Phi) = V_0 + \frac{1}{2} m_H^2 H^2 + \lambda \nu H^3 + \frac{1}{4} \lambda H^4$$

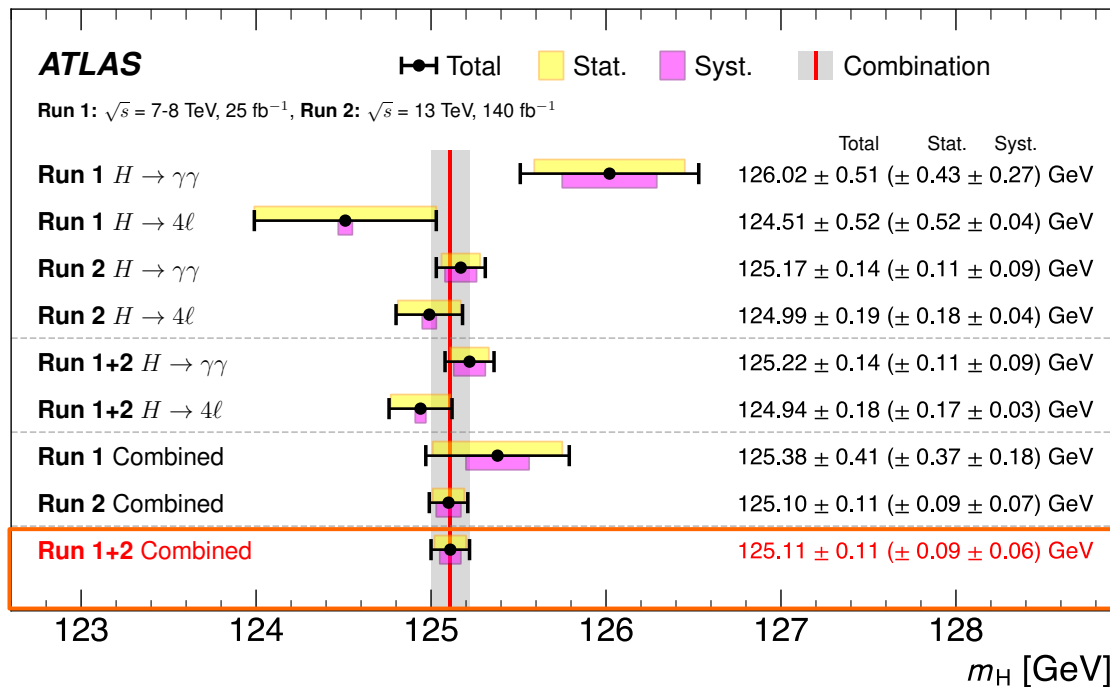
$$m_H = \sqrt{2} \mu = \sqrt{2 \lambda} \nu$$

How well do we know the Higgs mass?

$$m_H = \sqrt{2}\mu$$

- m_H Not predicted by theory
- All Higgs coupling properties depend on m_H
- Excellent measurements by ATLAS and CMS
 - ✓ Precision depends on energy (photons, electrons) and momentum (muons) calibration

Phys. Rev. Lett. 131 (2023) 251802



110 MeV
(0.09%)

Higgs width: a problem difficult to tackle directly!

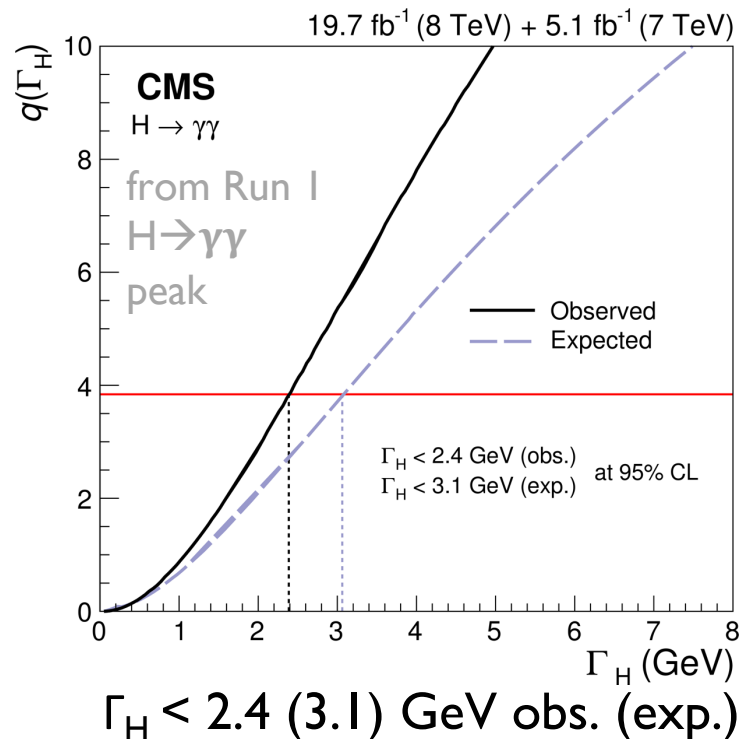
Total Higgs natural width in SM is **small!**

- ✓ Too small to be accessed experimentally at LHC from resonance line-shape in analysis where peak can be reconstructed...

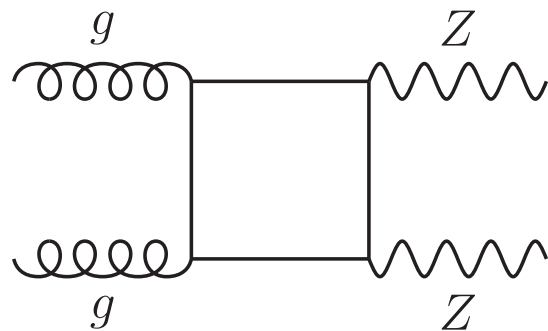
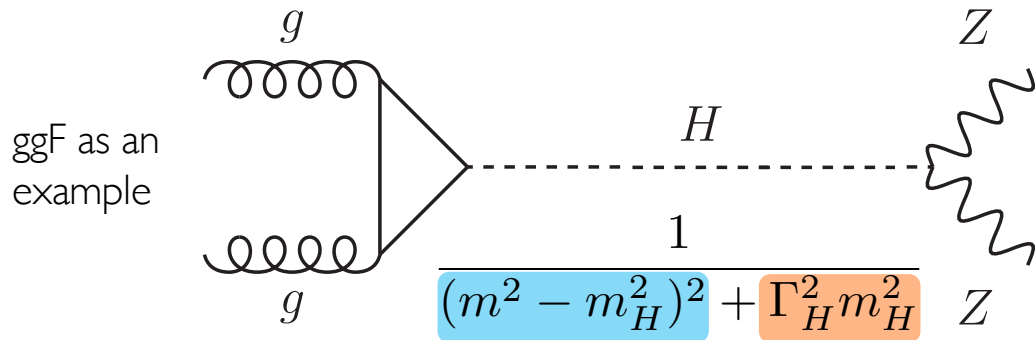
$$\Gamma_H^{\text{SM}} \text{ dominated by } \Gamma(H \rightarrow b\bar{b}) \approx \frac{N_c g_w^2 m_b^2 m_H}{32\pi m_W^2}$$

$$\Gamma_H^{\text{SM}} = 4.07 \text{ MeV}$$

Direct measurement severely limited by detector resolution! One (old) example:



The Higgs boson as propagator: width from off-shell Higgs

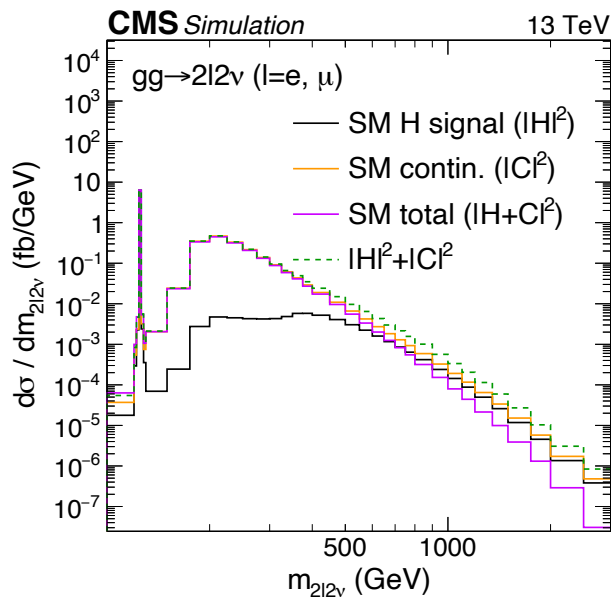


- **Interference** impacts both total cross section and $m(VV)$ line-shape
- Assuming **on-shell and off-shell couplings are equal**:

$$\frac{\mu_{\text{off-shell}}}{\mu_{\text{on-shell}}} = \frac{\Gamma}{\Gamma_{\text{SM}}}$$

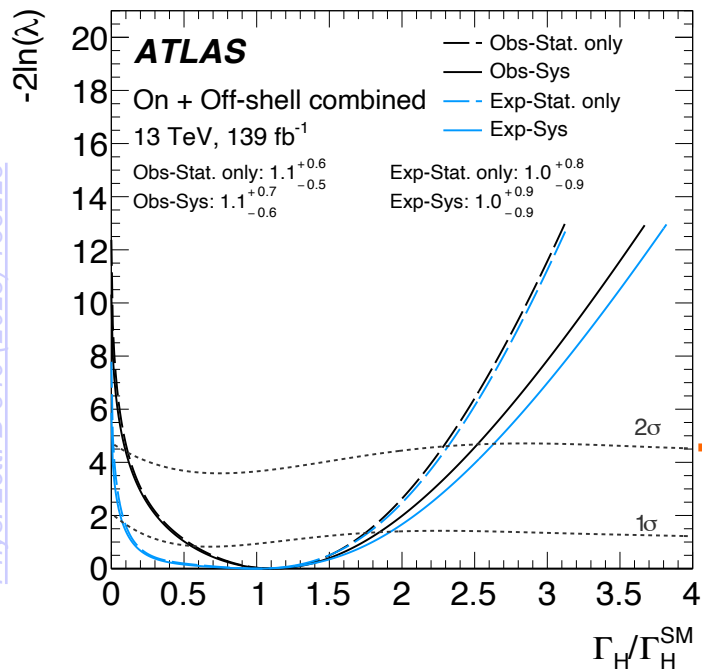
$$\begin{aligned}
 vv &= gg \\
 vv &= WW, ZZ, Z\gamma, \gamma\gamma
 \end{aligned}$$

$$\sigma_{vv \rightarrow H \rightarrow 4\ell}^{\text{on-shell}} \propto \frac{g_{\text{gluon}}^2 g_V^2}{\Gamma_H} \quad \sigma_{vv \rightarrow H \rightarrow 4\ell}^{\text{off-shell}} \propto g_{\text{gluon}}^2 g_V^2$$



Measurements of the Higgs width from off-shell production

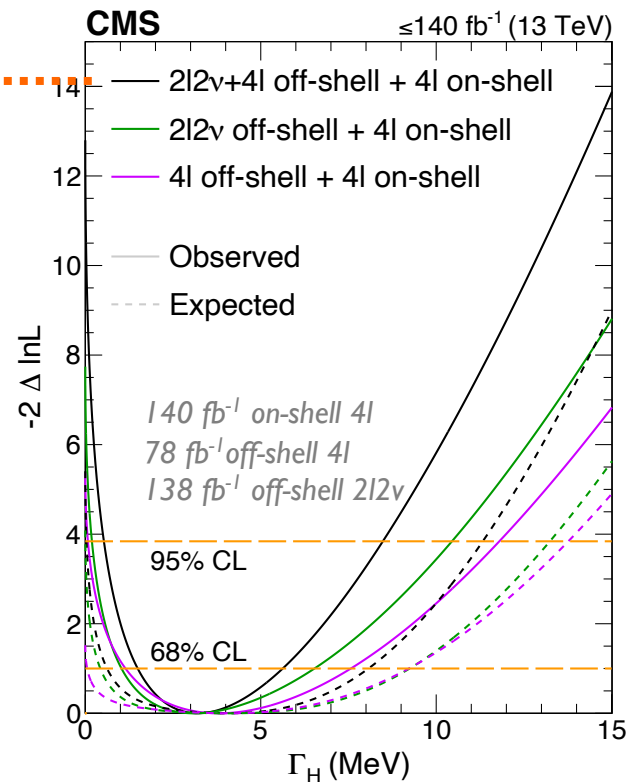
Measurements in **4l** and **2l2v** final states and for **different production modes**



CMS
 $\Gamma_H = 3.2^{+2.5}_{-1.7}$ MeV

~3 σ evidence for off-shell H production

ATLAS
 $\Gamma_H = 4.5^{+3.3}_{-2.5}$ MeV



arXiv:2202.06923

Phys. Lett. B 846 (2023) 138223

How well will we know the Higgs width?

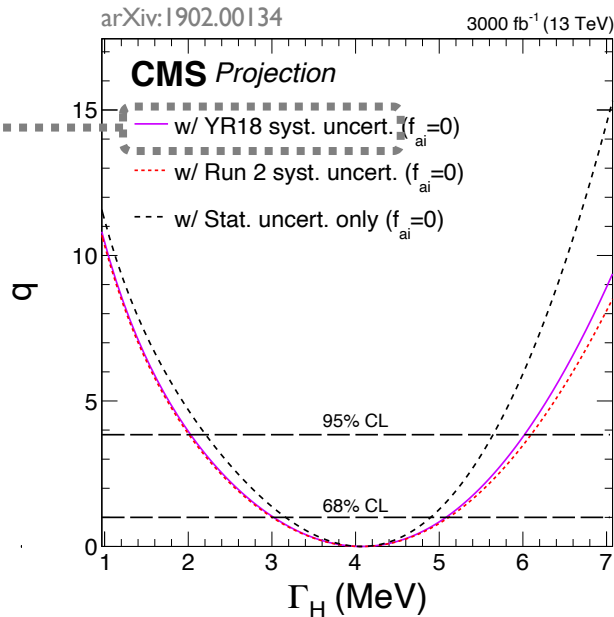
- Combined ATLAS + CMS width sensitivity with 3000 fb^{-1} , based on the off-shell measurement

✓ $\Gamma_H = 4.1 \pm 0.8 \text{ MeV}$

- More conservative assumptions on theoretical uncertainties than Run 2 results
 - i.e. signal and background k-factors

“YR18” systematics: uncertainties reduced wrt Run 2 according to the improvements expected to be reached at the end of HL-LHC e.g. theory uncertainties generally halved. See

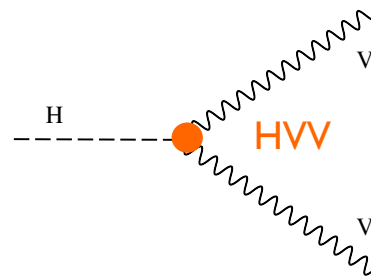
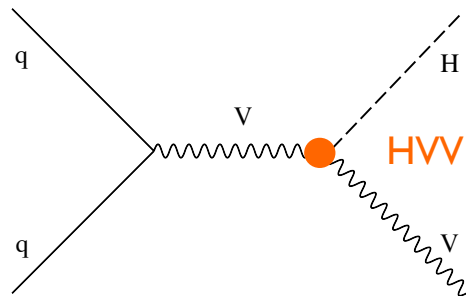
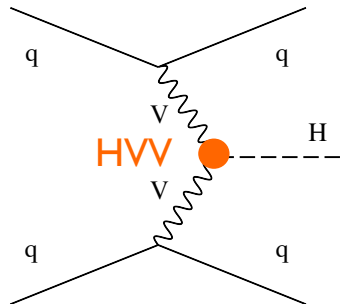
<https://twiki.cern.ch/twiki/bin/view/LHCPhysics/HLHELHCCCommonSystematics>



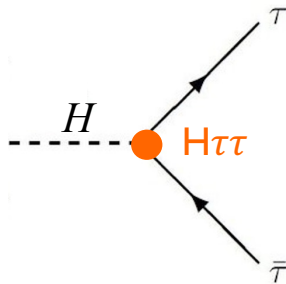
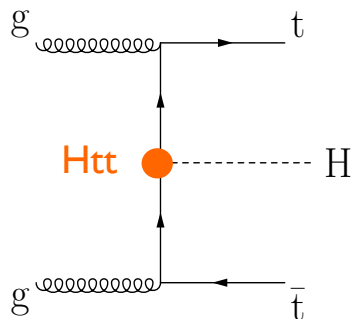
Studies of the Higgs CP properties

Spin is property of the particle, CP of the coupling...

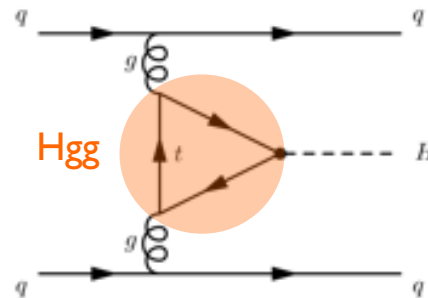
coupling to EW vector bosons



coupling to fermions



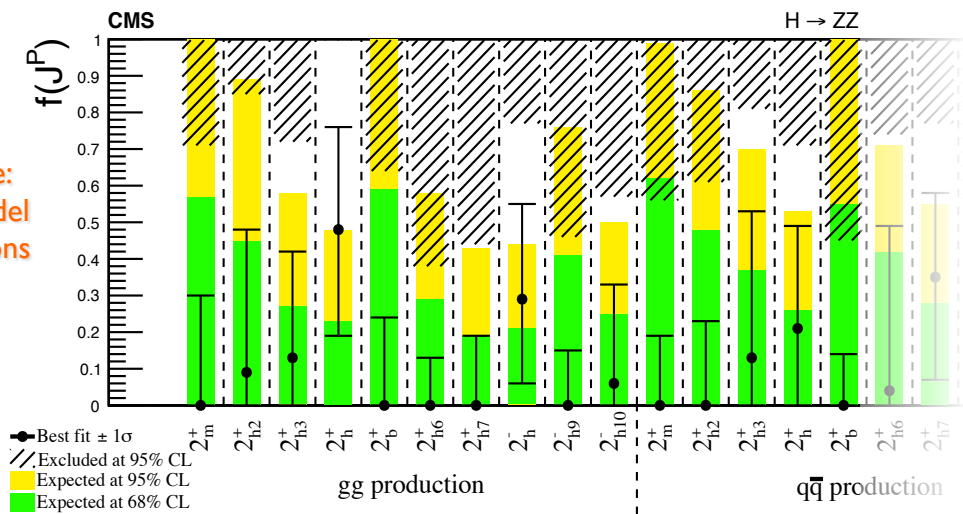
coupling to gluons



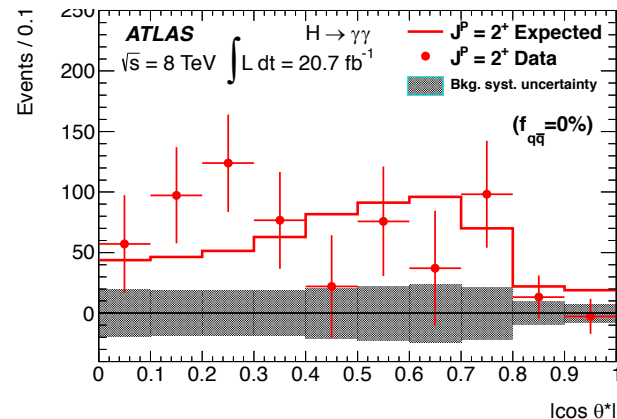
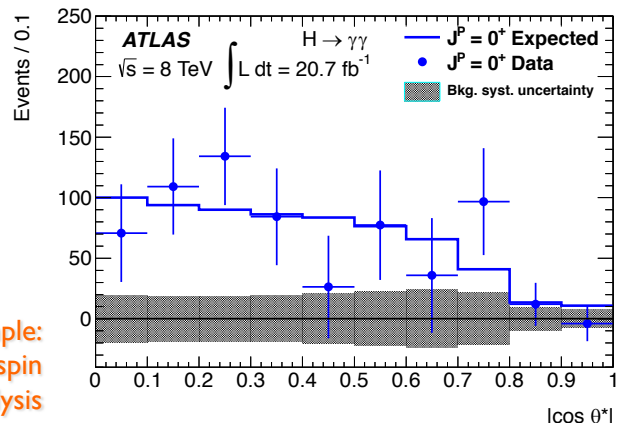
SM Higgs spin and CP properties

- SM Higgs has spin 0 and positive (even) parity ($J^{CP} = 0^{++}$)
- At the end of Run I we knew Higgs had spin 0...
 - ✓ Spin 1 and 2 hypotheses excluded at $> 99.9\%$ CL using $H \rightarrow \gamma\gamma$, $H \rightarrow ZZ^*$ and $H \rightarrow WW^*$

Example:
J=2 model
exclusions
with
 $H \rightarrow ZZ$



Example:
 $H \rightarrow \gamma\gamma$ spin
analysis



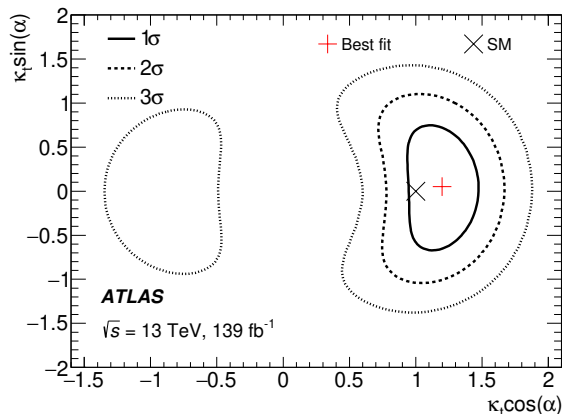
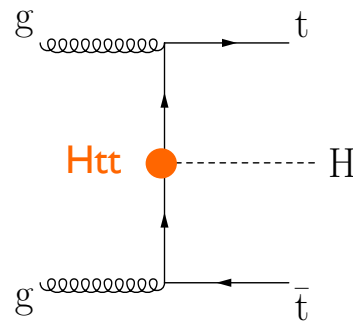
$$\cos \theta^* = \frac{\sinh(\eta_{\gamma_1} - \eta_{\gamma_2})}{\sqrt{1 + (p_T^{\gamma\gamma} / m_{\gamma\gamma})^2}} \cdot \frac{2p_T^{\gamma_1} p_T^{\gamma_2}}{m_{\gamma\gamma}^2}$$

$\gamma\gamma$ polar angle ϑ^* with respect to Z-axis in Collins-Soper frame

CP properties of Higgs-top coupling with $t\bar{t}H$

Effective Lagrangian for Yukawa coupling to top quarks parameterized by CP-Even and CP-odd components

$$\mathcal{L}_{t\bar{t}H} = -\kappa'_t y_t \phi \bar{\psi}_t (\cos \alpha + i\gamma_5 \sin \alpha) \psi_t$$



ATLAS $t\bar{t}H$
 $H \rightarrow \gamma\gamma$

pure CP-odd coupling

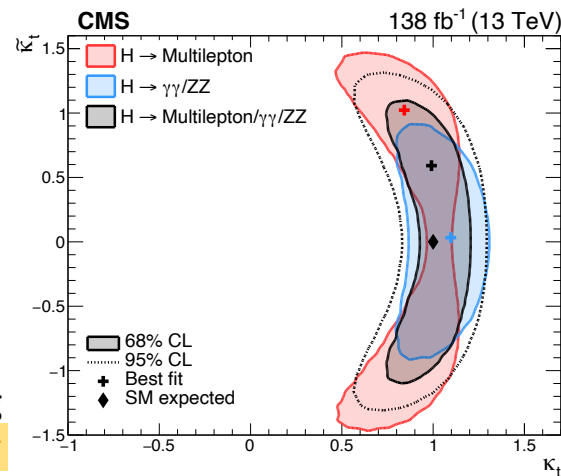
excluded at 3.9σ

Two examples:

CMS $t\bar{t}H$
 $H \rightarrow ML$
 $+ H \rightarrow \gamma\gamma$
 $+ H \rightarrow ZZ$

pure CP-odd coupling

excluded at 3.7σ



$|\alpha| > 43$ deg excluded @ 95% CL.

$$f_{CP}^{Htt} = \frac{\tilde{\kappa}_t^2}{\tilde{\kappa}_t^2 + \kappa_t^2} \quad \mathcal{A}(Hff) = \frac{m_f}{v} \bar{\psi}_f (\kappa_t + i\tilde{\kappa}_t \gamma_5) \psi_f$$

$$|f_{Htt}| = (\sin \alpha)^2 \quad |f_{Htt}| = 0.28 (< 0.55 \text{ at } 1\sigma)$$

7.

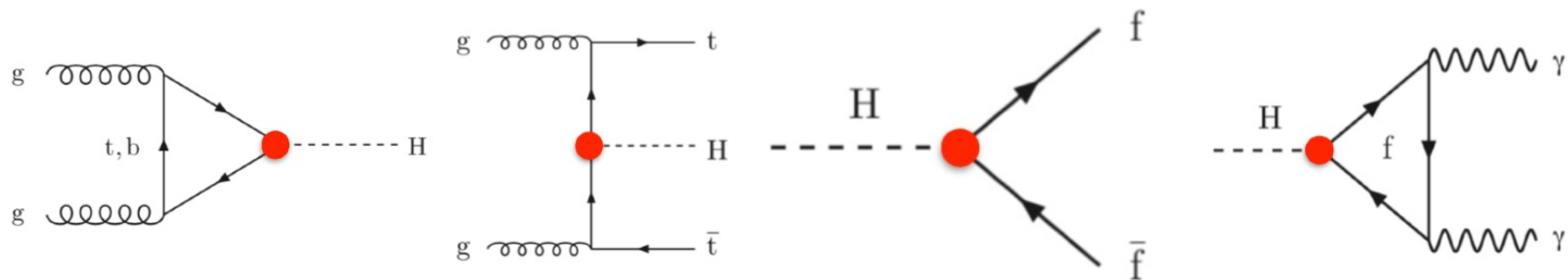
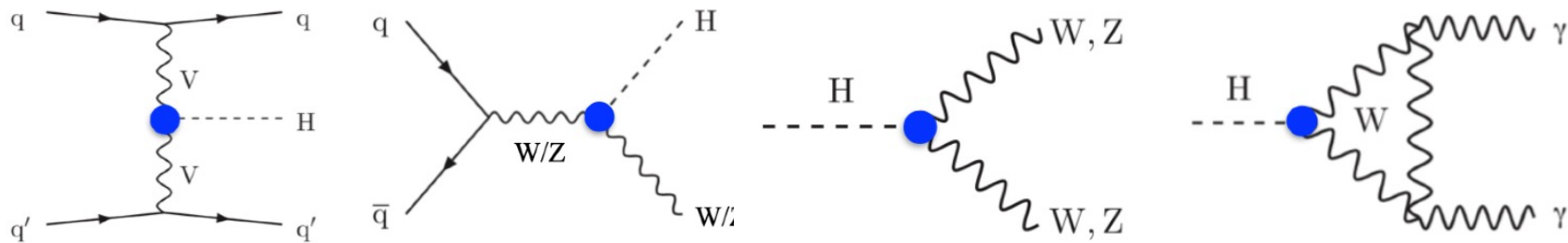


How well
do we know
the Higgs couplings?



Couplings to...

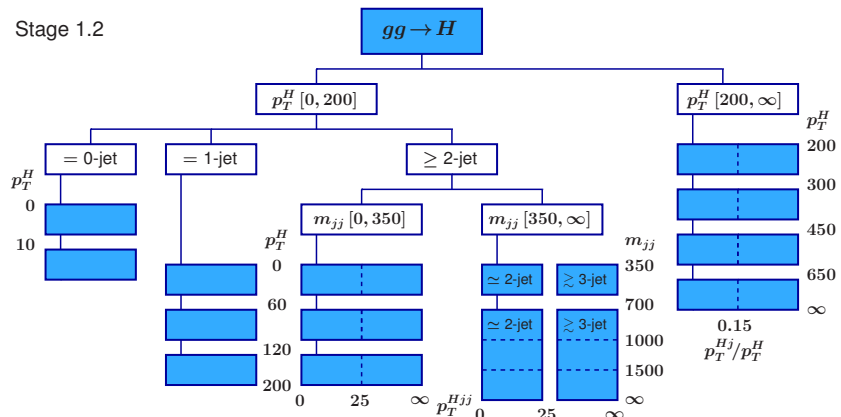
... gauge bosons



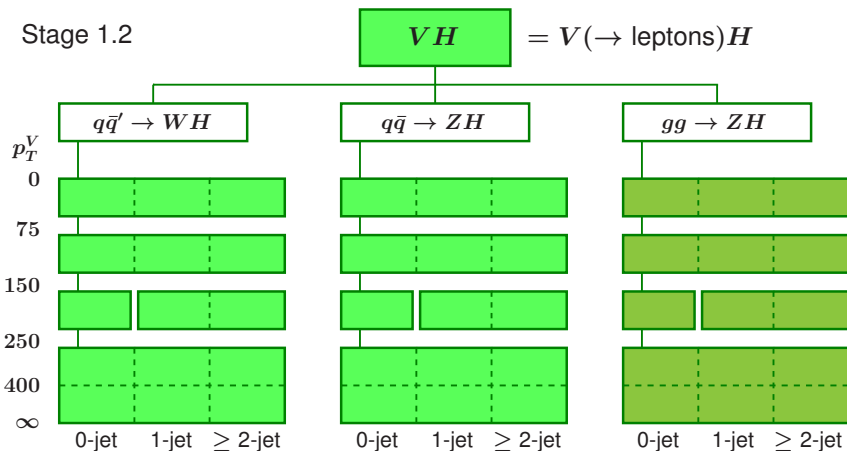
... fermions

STXS: Simplified Template Cross Sections

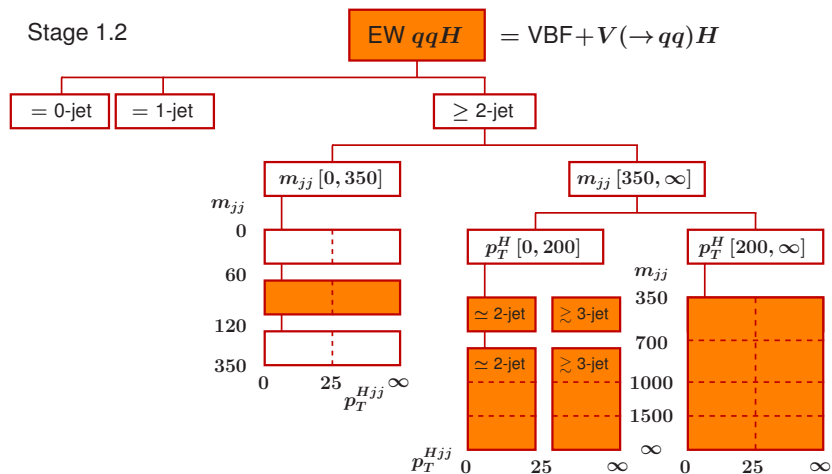
Stage 1.2



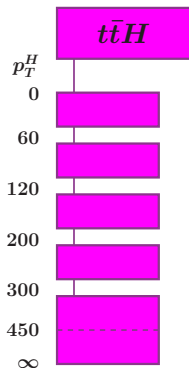
Stage 1.2



Stage 1.2

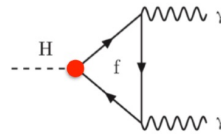


Stage 1.2

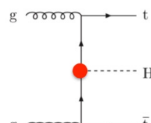
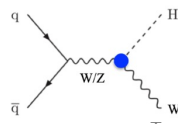
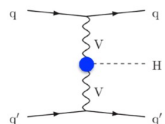
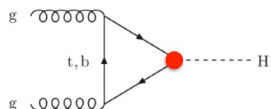
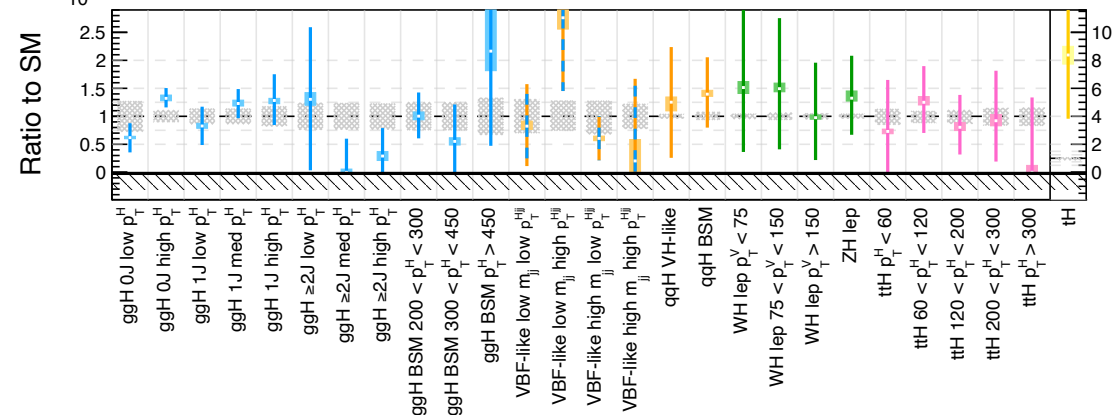
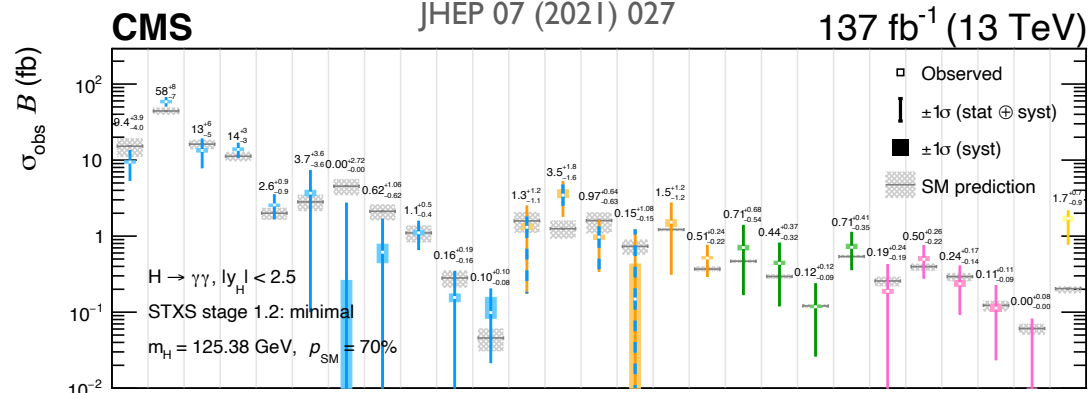


- Cross sections per production mode in different regions of phase space
 - ✓ Reduce dependency on theory
 - ✓ Allow MVA analysis
 - ✓ Exploit sensitivity of each channel specific region of phase space (designed for combination)
 - ✓ Inclusive in decays (for now)
 - ✓ Binning defined to allow better estimate of theoretical uncertainties, and easier merging when insufficient experimental sensitivity
 - ✓ Largely model-independent approach to **test for BSM deviations in kinematic distributions**

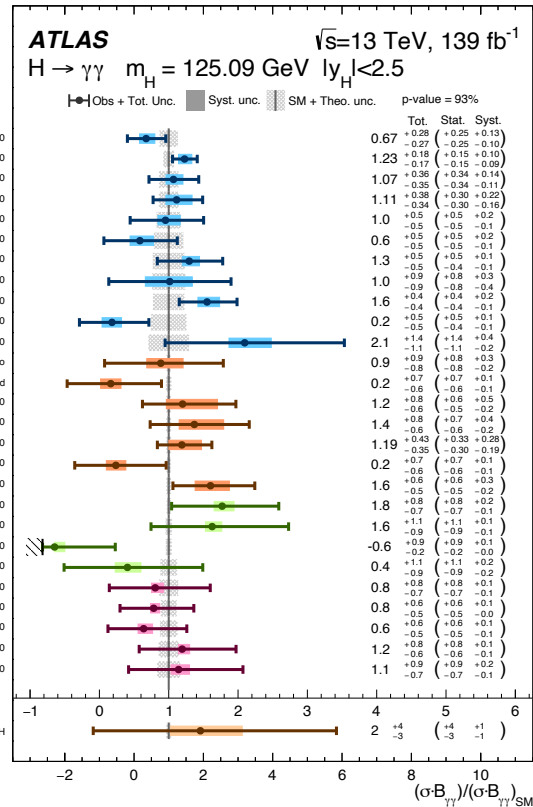
STXS at work: $H \rightarrow \gamma\gamma$



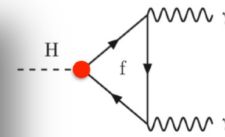
JHEP 07 (2023) 088



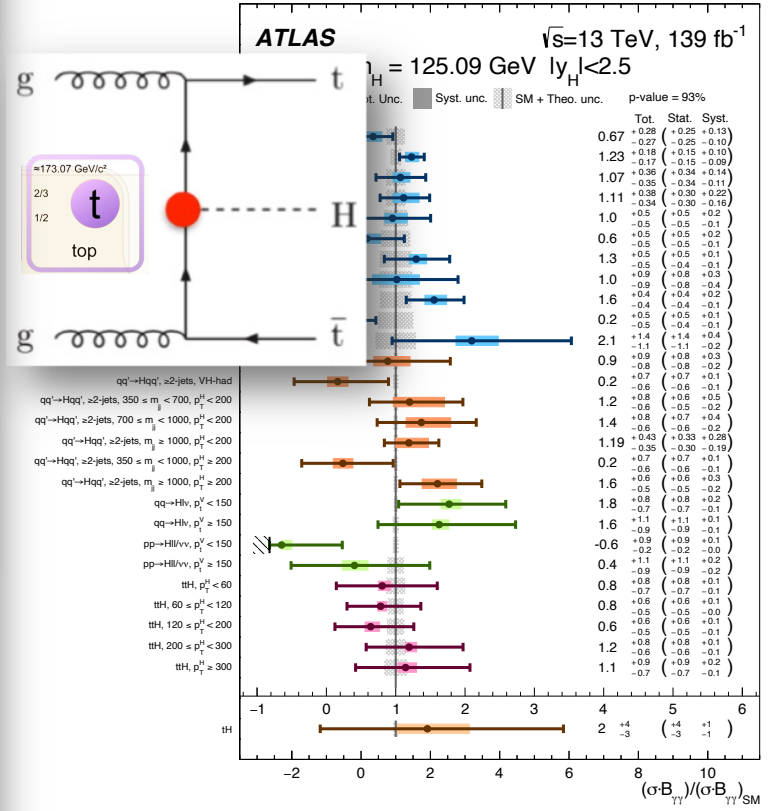
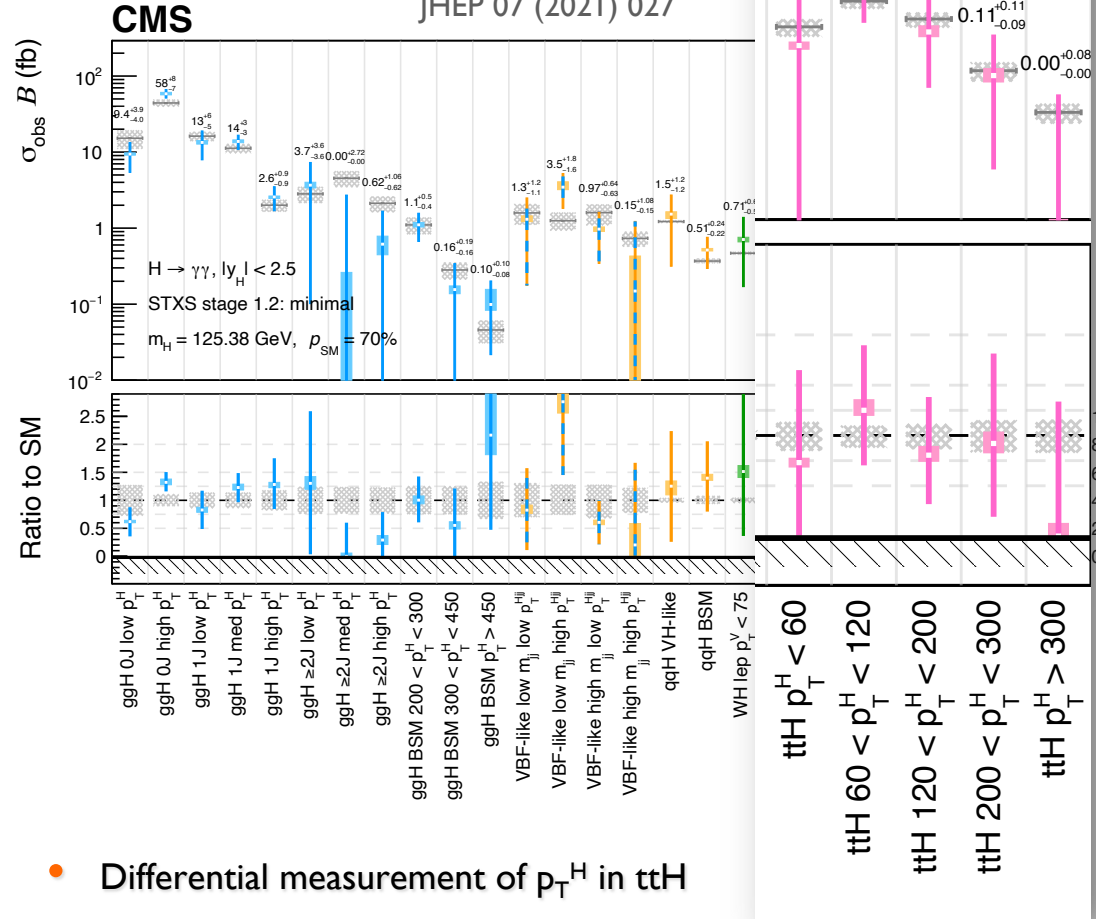
Twelve years with the Higgs boson



STXS at work: $H \rightarrow \gamma\gamma$

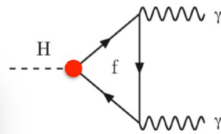


arXiv:2207.00348

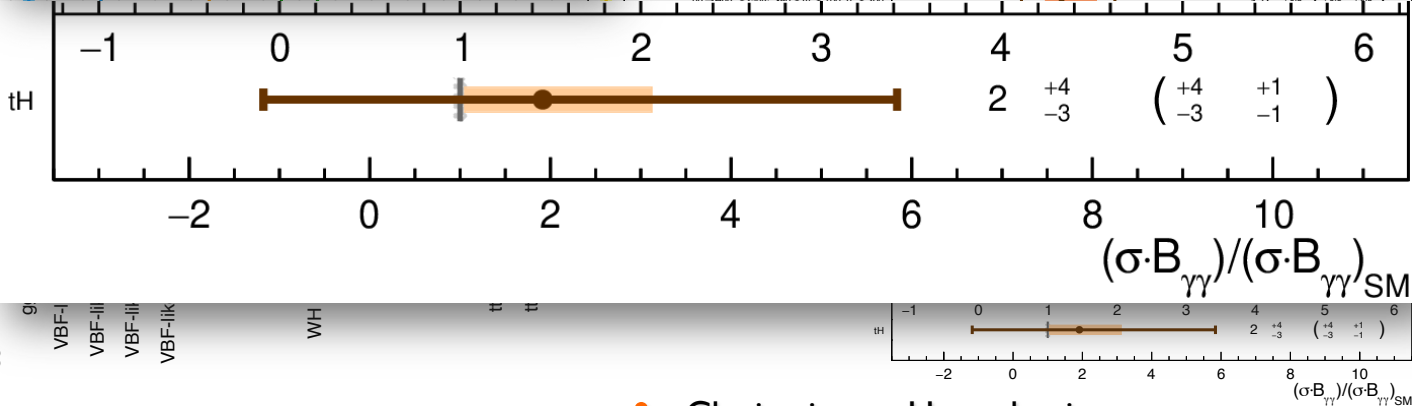
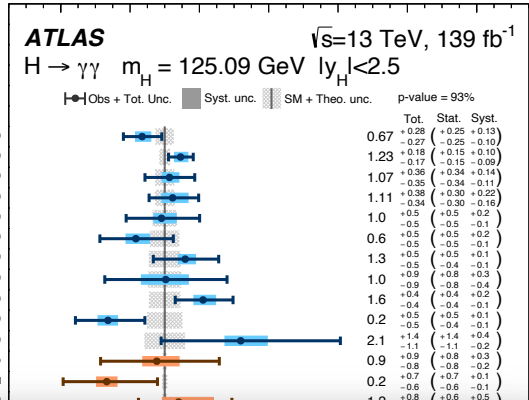
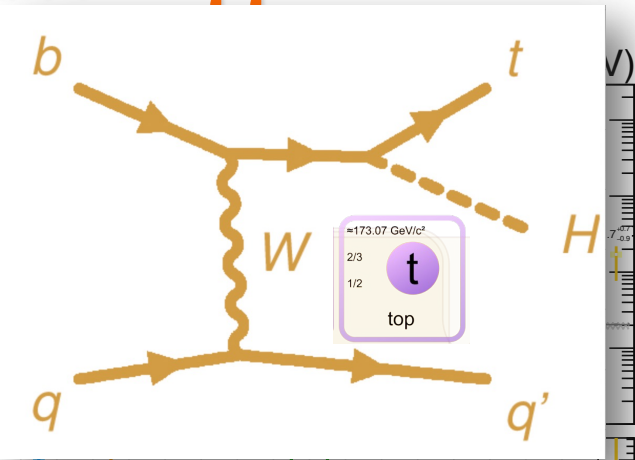
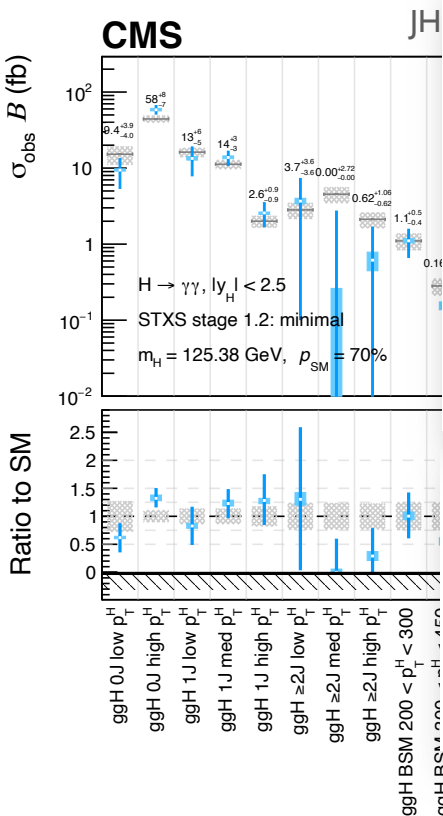


- Differential measurement of p_T^H in ttH

STXS at work: $H \rightarrow \gamma\gamma$



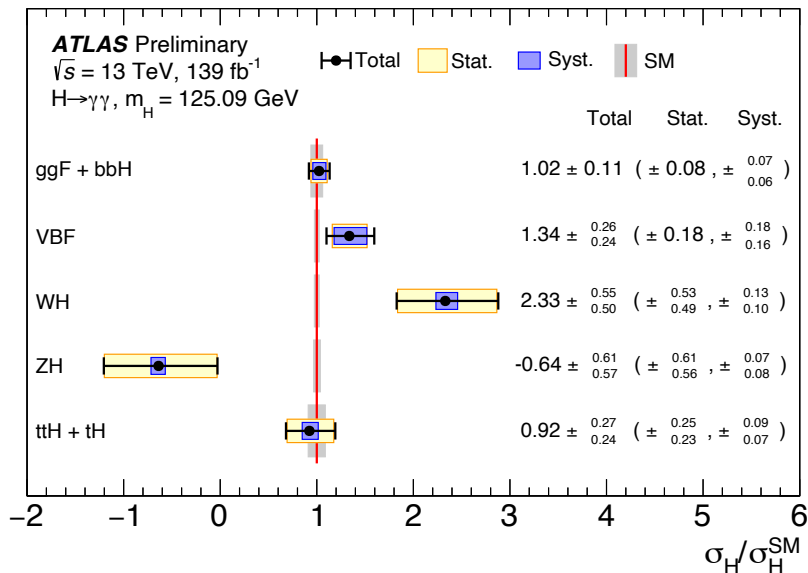
arXiv:2207.00348



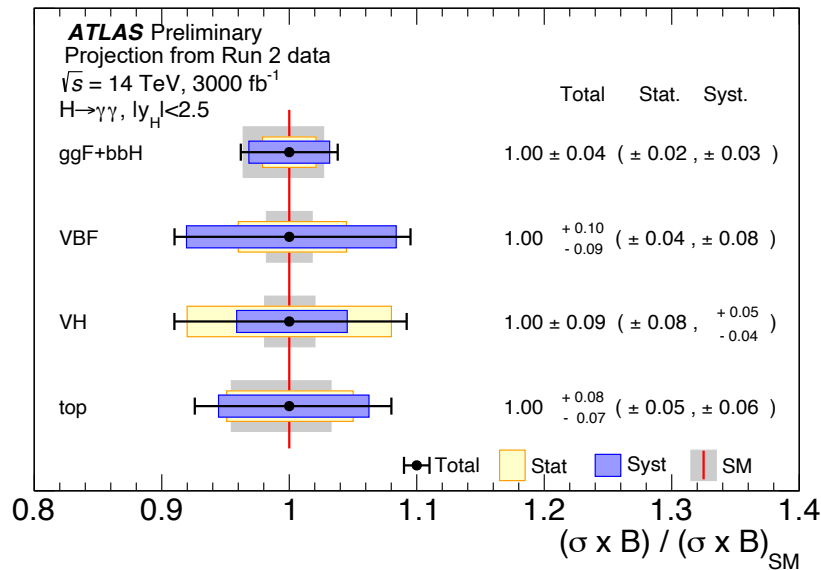
- Closing in on tH production
- ✓ ~10 × σ_{tH}SM excluded at 95% CL

H \rightarrow $\gamma\gamma$ as an example: Run 2 vs HL-LHC

- While STXS measurement remains main current goal, traditional “production-mode” results (a.k.a. “Stage 0” STXS) still measured by more powerful decay channels and in combination
- Already at Run 2 some measurement **limited by systematics: more and more true in the future!**



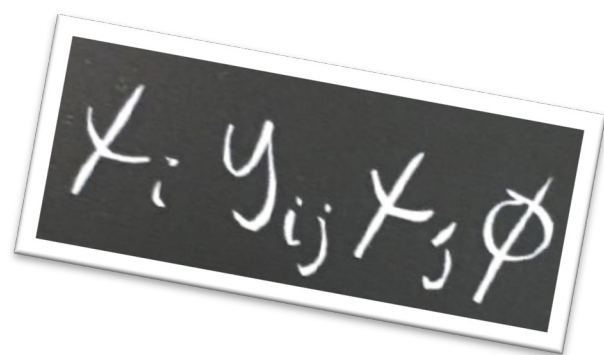
ATLAS-CONF-2020-026



arXiv:1902.00134

8.

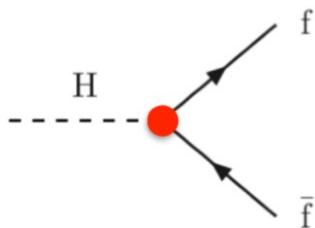
zoom on Higgs coupling to fermions



PDG 2012

$b\bar{b}$
 $\tau^+\tau^-$

possibly seen
possibly seen



PDG 2022

$c\bar{c}$ Final State = 37 ± 20
 $b\bar{b} = 0.98 \pm 0.12$
 $\mu^+\mu^- = 1.19 \pm 0.34$
 $\tau^+\tau^- = 1.15^{+0.16}_{-0.15}$
 $Z\gamma < 3.6, \text{ CL} = 95\%$
 $\gamma^*\gamma$ Final State = 1.5 ± 0.5
 $t\bar{t}H^0$ Production = 1.10 ± 0.18



Why is Yukawa interaction important?

In SM Yukawa interaction between Higgs and fermions gives fermions their mass...

- **Why is elementary mass important? Two ideas...**

- ✓ Mass of quark u and d is responsible for difference of mass between neutron and proton, thus of proton stability, thus of existence of hydrogen atoms...

proton (up+up+down): $2.2 + 2.2 + 4.7 + \dots = 938.3$ MeV

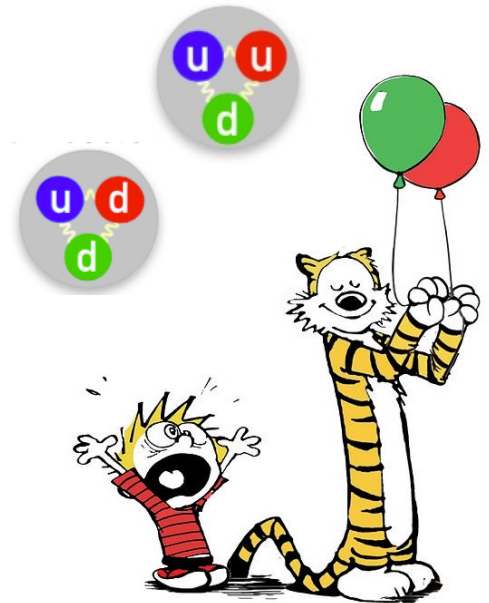
neutron (up+down+down): $2.2 + 4.7 + 4.7 + \dots = 939.6$ MeV

- ✓ Mass of electron determines Bohr radius, thus dimensions of atoms, thus all chemistry

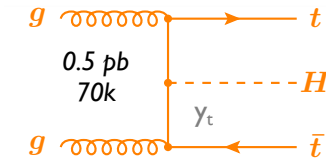
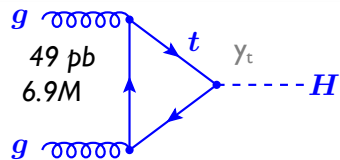
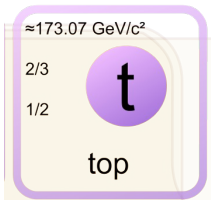
$$a_0 = \frac{4\pi\epsilon_0 \hbar^2}{m_e e^2} = \frac{\hbar}{m_e c \alpha}$$



$$m_i = -\frac{y_i \nu}{\sqrt{2}} \quad \nu = \frac{\mu}{\sqrt{\lambda}}$$

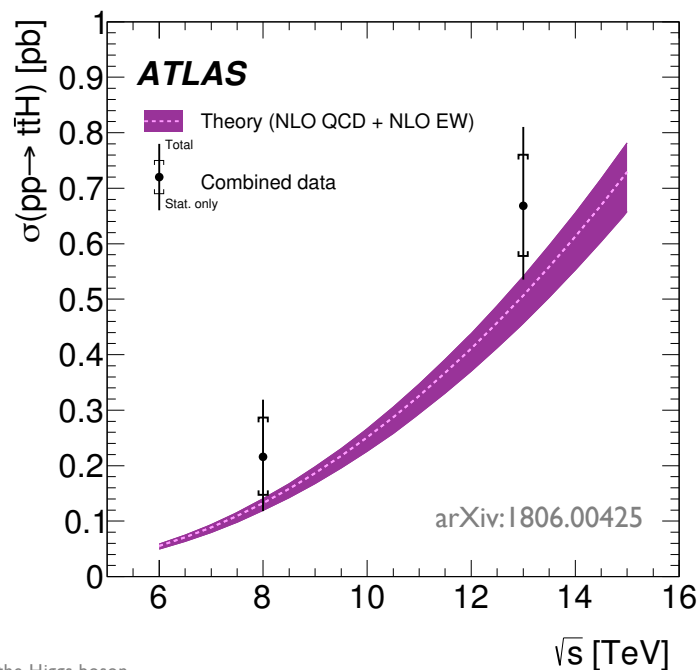
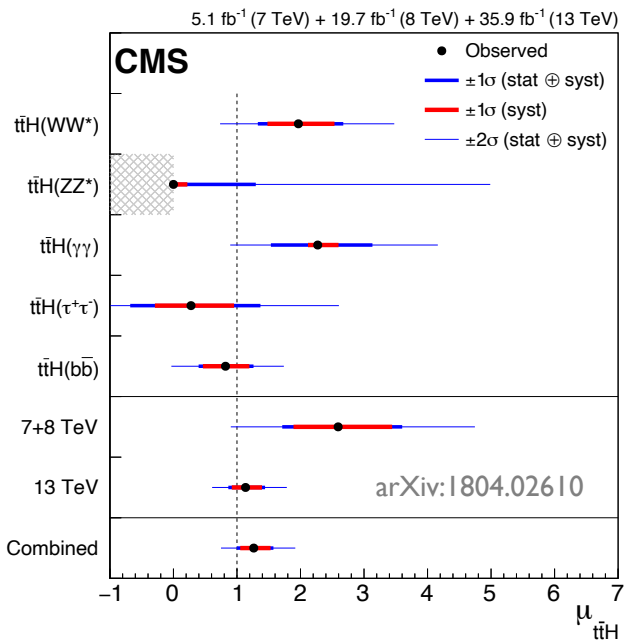


Coupling to top quark: direct observation of $t\bar{t}H$ in 2018!



CMS (Run I + 36 fb⁻¹ Run 2): 5.2 σ (4.2 σ exp)

ATLAS: (Run I + 36-80 fb-I Run 2): 6.3 σ (5.1 σ exp)





 **ATLAS**
EXPERIMENT

Run: 331742

Event: 1873900334

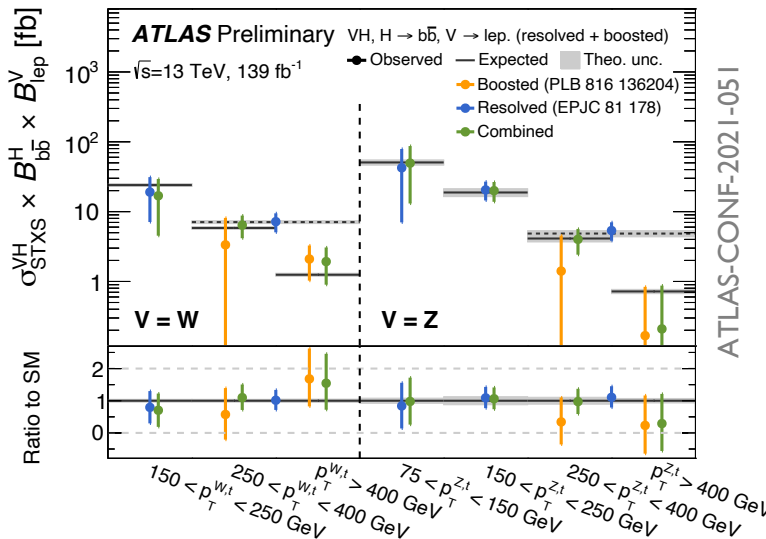
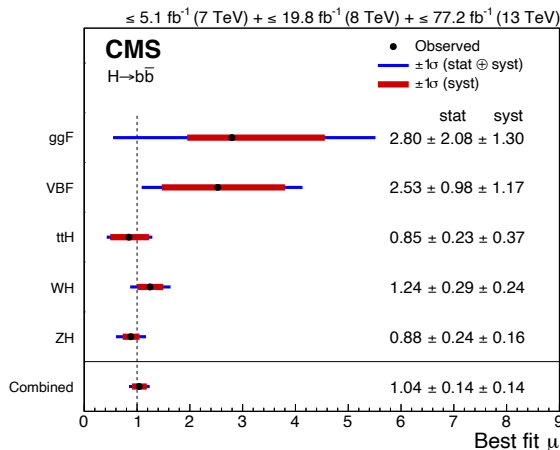
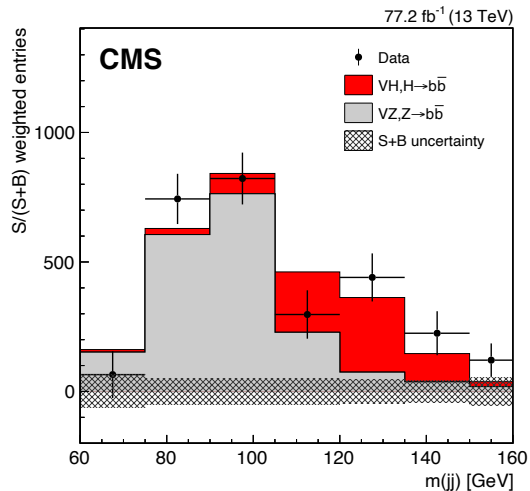
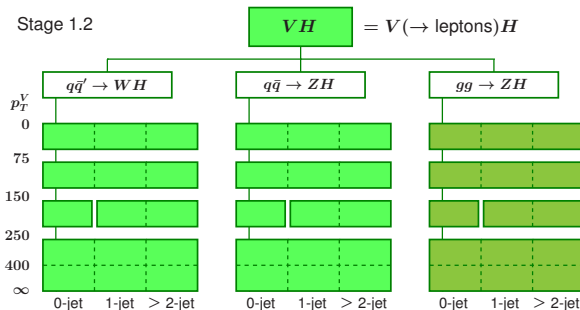
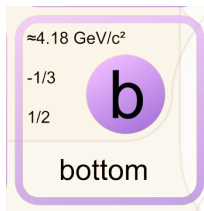
2017-08-04 21:48:42 CEST

Coupling to bottom quark: observation of VH/H → bb in 2018!

- Difficult channel despite large BR(58%) due to large backgrounds

✓ VH production most sensitive → **Observation in 2018!**

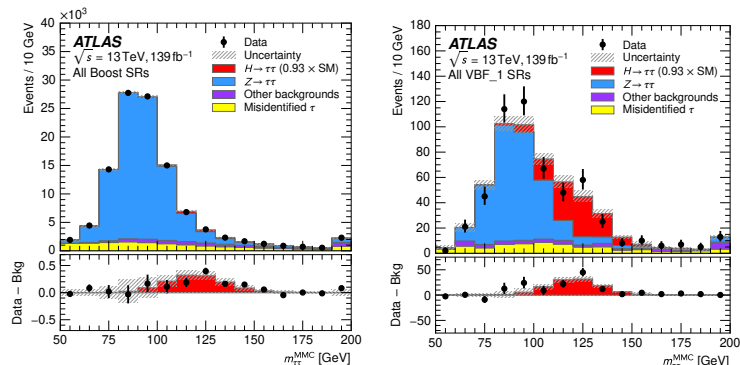
- ATLAS: 5.4σ (5.5σ) ATLAS: arXiv:1808.08238
- CMS: 5.5σ (5.6σ) CMS: arXiv:1808.08242



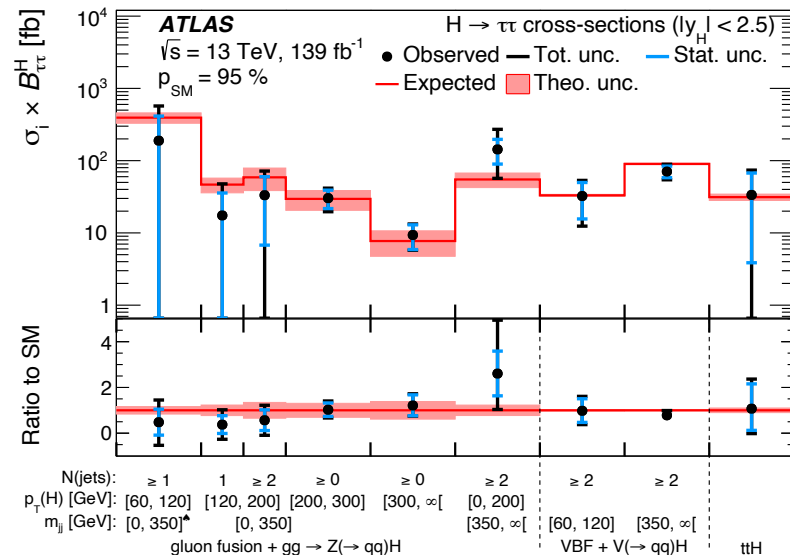
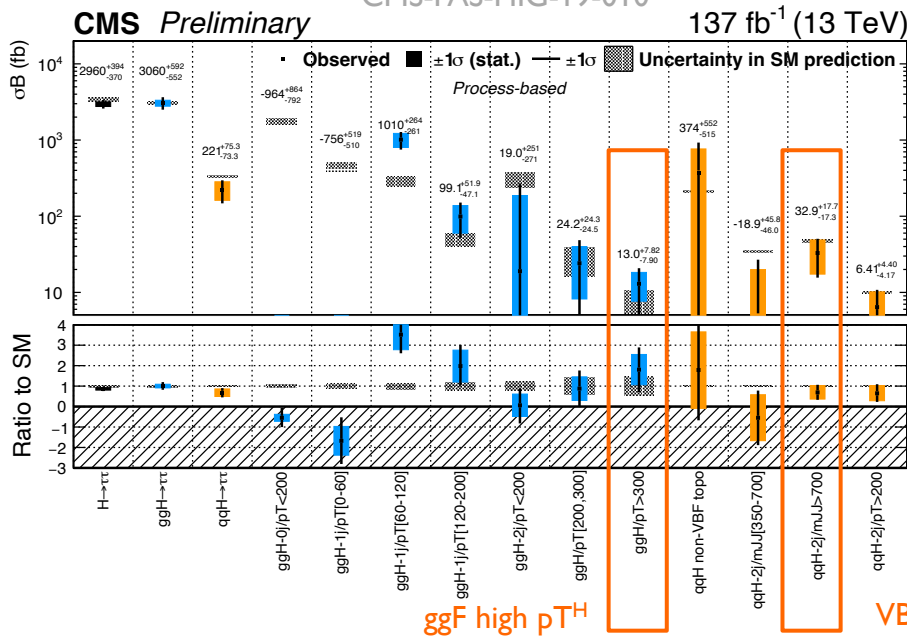
ATLAS-CONF-2021-051

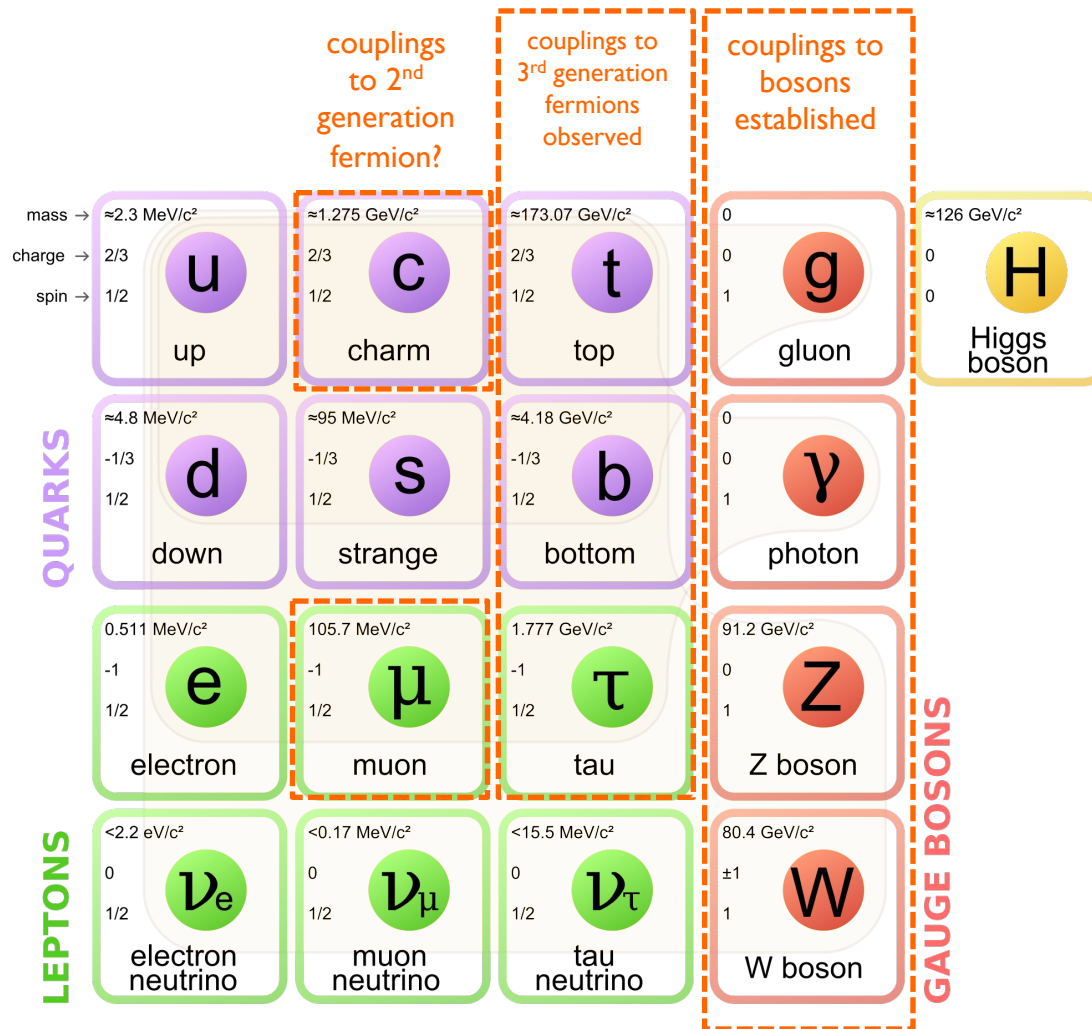
Coupling to 3rd generation leptons: $H \rightarrow \tau\tau$

- Observation already in Run I ATLAS+CMS coupling combination
- Now more precise measurements, bringing sensitivity to regions of the phase space less well measured by e.g. $H \rightarrow \gamma\gamma$ and $H \rightarrow 4l$
 - ✓ i.e. ggF high p_T^H and especially VBF



CMS-PAS-HIG-19-010



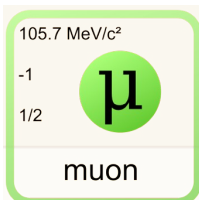


Is the Higgs boson responsible for the EW symmetry breaking also responsible for the masses of fermions?

Is the Higgs boson responsible for the masses of all fermions?

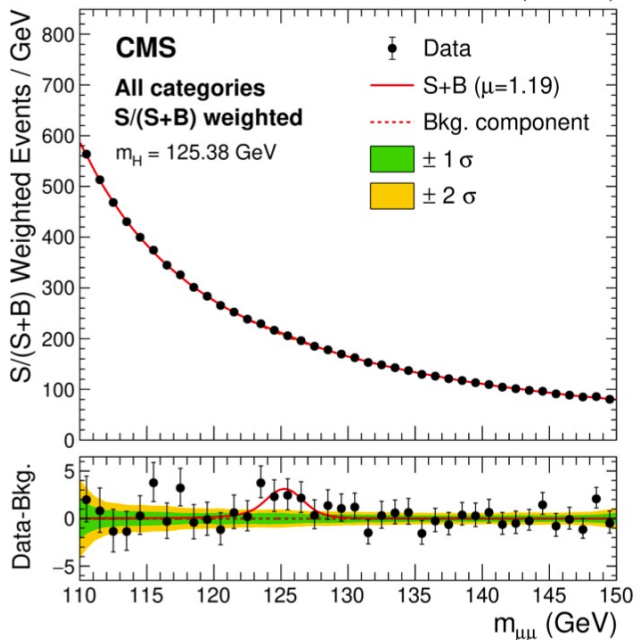
Coupling to 2nd generation leptons: $H \rightarrow \mu\mu$ evidence in 2020!

- $H \rightarrow \mu\mu$ very rare ($BR \sim 2 \cdot 10^{-4}$) with large resonant background from $DY \rightarrow \mu\mu$ ($S/B \sim 0.1\%$)



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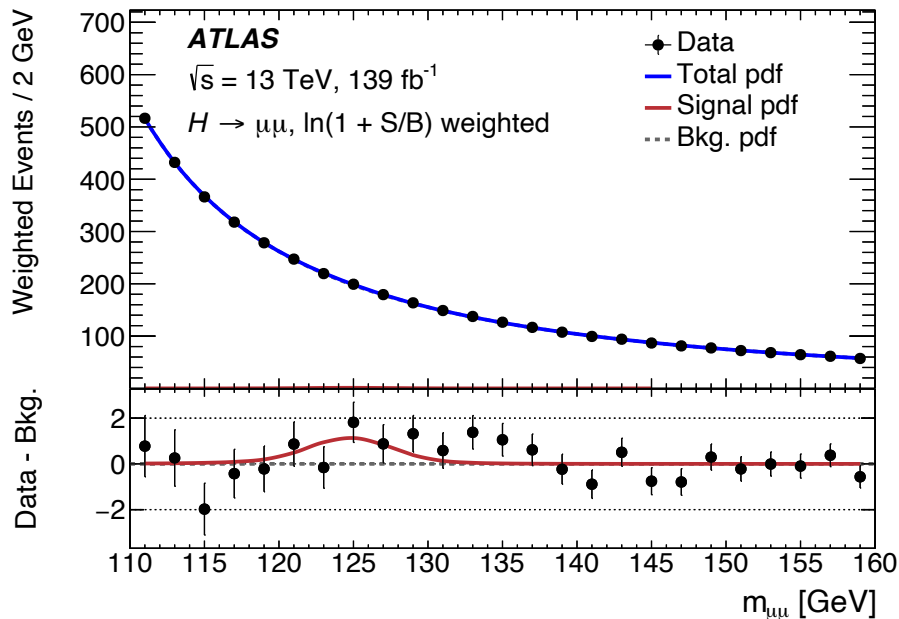
137 fb⁻¹ (13 TeV)



$\mu = 1.2 \pm 0.2$

Significance: **3.0 σ** (2.5 expected)

PLB 812(2021) 135980



$\mu = 1.2 \pm 0.6$

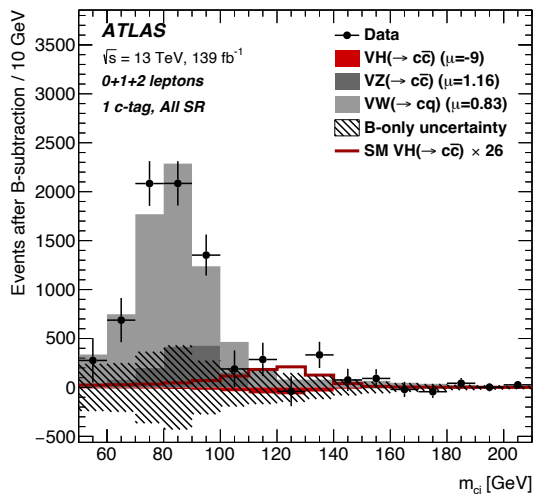
Significance: **2.0 σ** (1.7 expected)

The next frontier: coupling to 2nd generation quarks $H \rightarrow c\bar{c}$

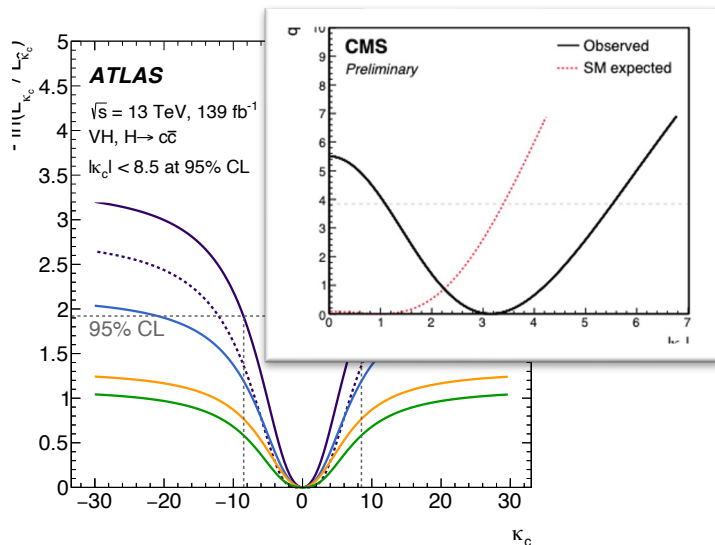
- Very challenging channel: large backgrounds from multi-jets
 - ✓ c-tagging needed to discriminate $H \rightarrow b\bar{b}$

$$\mu_{VH(c\bar{c})}(\kappa_c) = \frac{\kappa_c^2}{1 + B_{H \rightarrow c\bar{c}}^{\text{SM}}(\kappa_c^2 - 1)}$$

Eur. Phys. J. C 82 (2022) 717



VZ,VW: 3.8 σ for VW(\rightarrow cq)

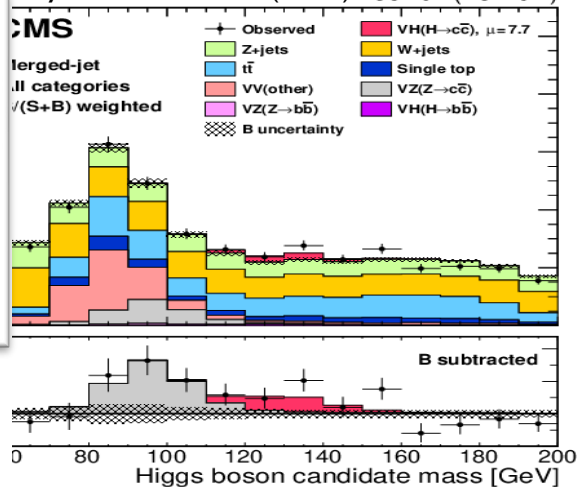


limits on κ_c coupling

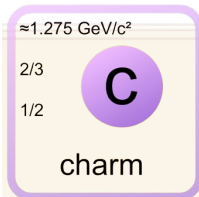
ATLAS: $\kappa_c = \sigma/\sigma_{\text{SM}} < 8.5$ (12.4 exp.)

CMS: $1.1 < \kappa_c < 5.5$ (< 3.4 exp.)

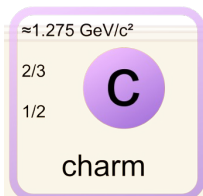
Phys. Rev. Lett. 131 (2023) 061801



$Z \rightarrow c\bar{c}$: 5.7 σ

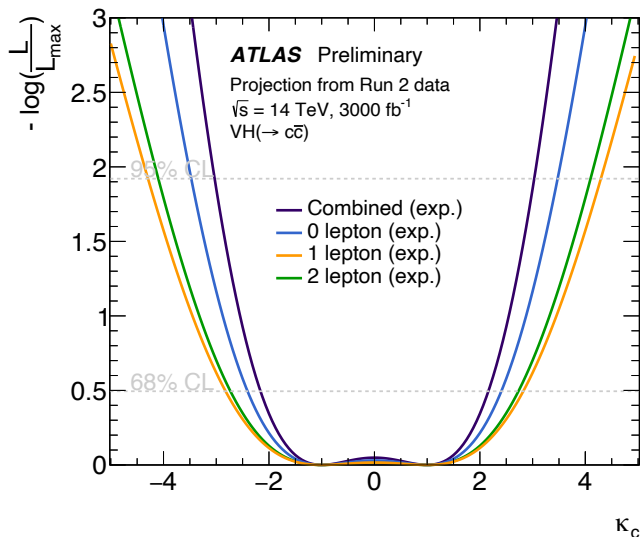


Will we observe $H \rightarrow cc$ in the future?

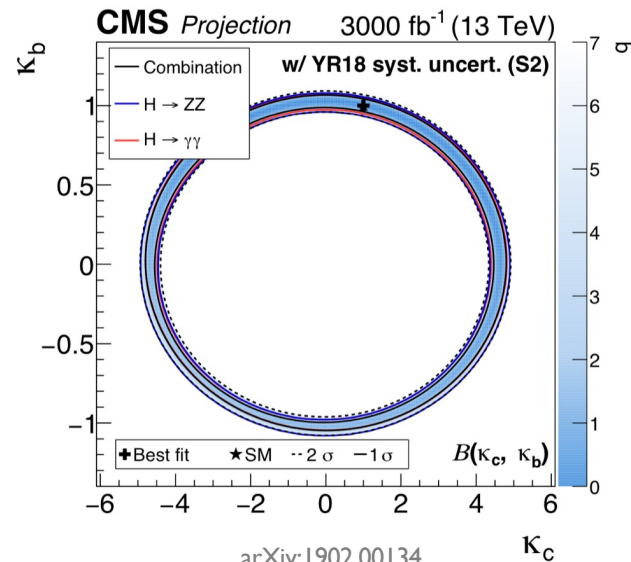
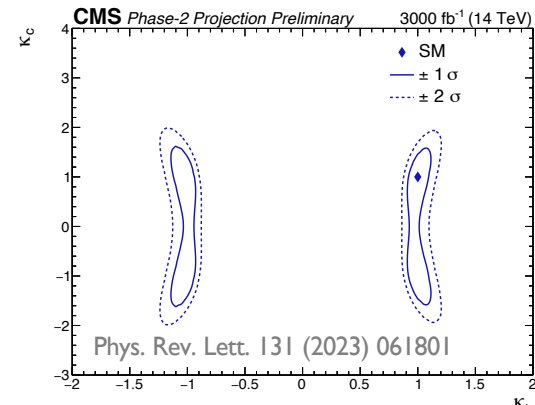
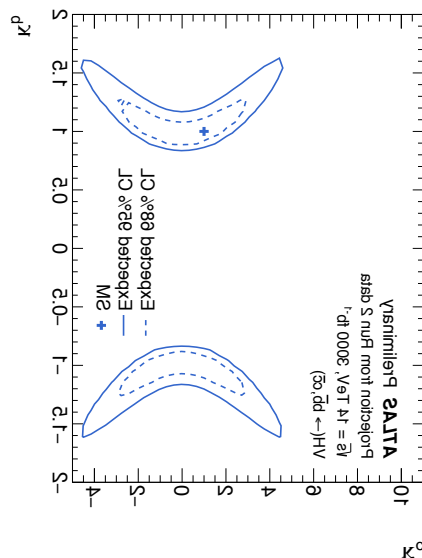


- Extrapolations to HL-LHC luminosity
 - ✓ Sensitivity to k_c (and k_b) from ATLAS $VH/H \rightarrow cc$ and $VH/H \rightarrow bb$ analyses
 - Sensitivity to k_c and k_b from CMS p_T^H differential measurements

ATL-PHYS-PUB-2021-039



$k_c = \sigma/\sigma_{SM} < 3 \text{ @ } 95\% \text{ CL}$



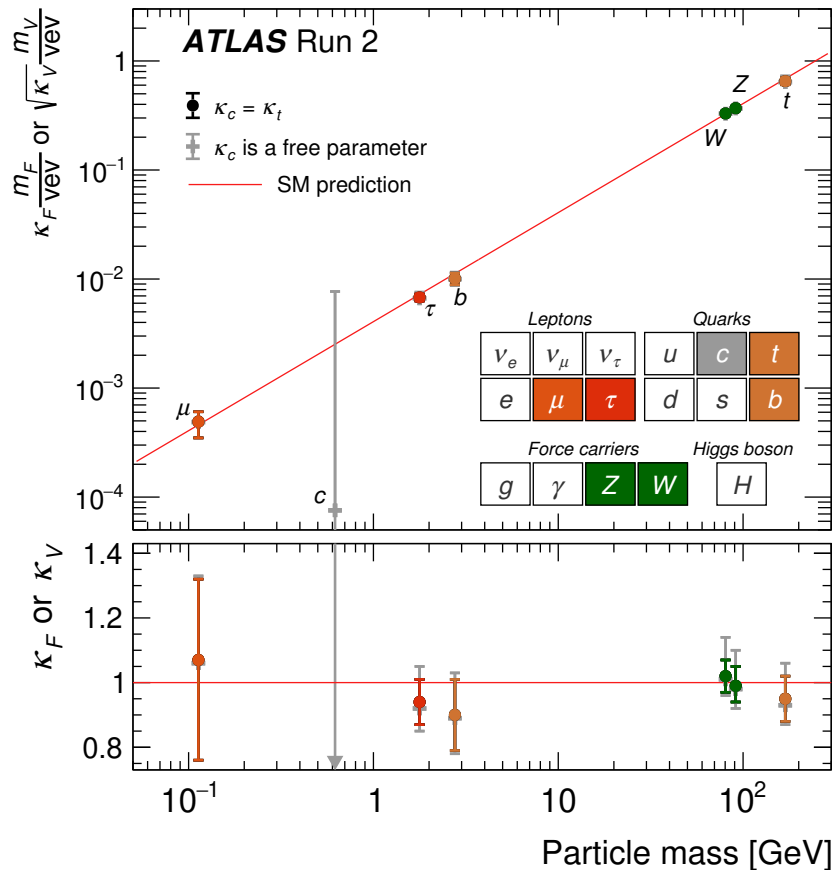
9.

One for all,
all for one!



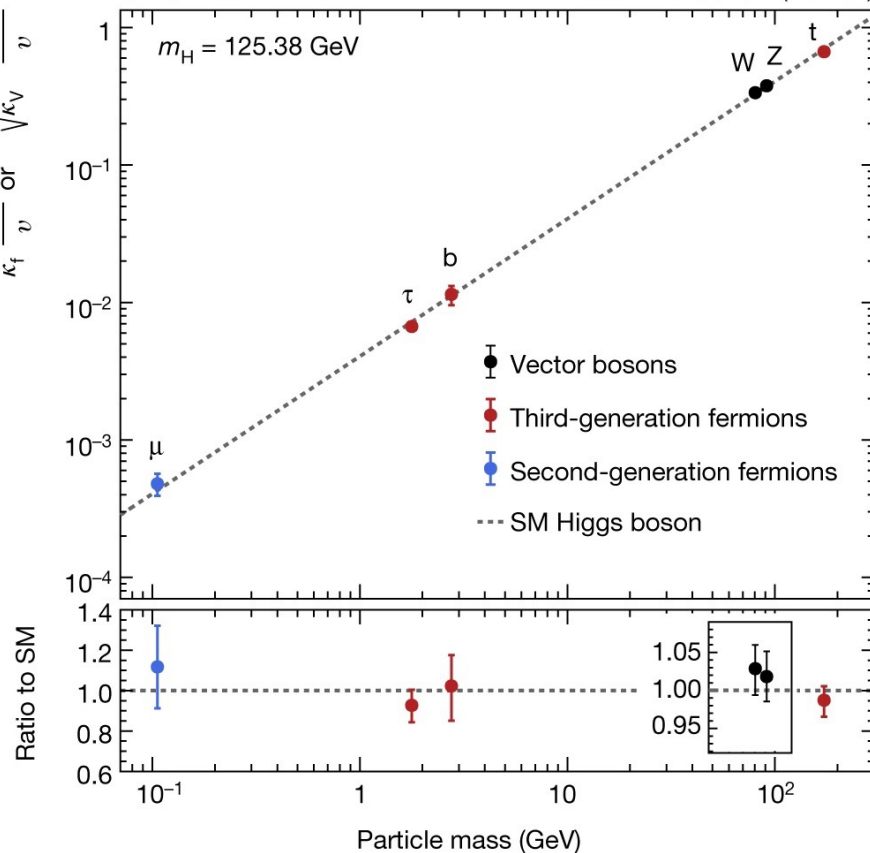
The Standard Model rules (over 3 order of magnitudes)!

Nature 607 (2022) 52

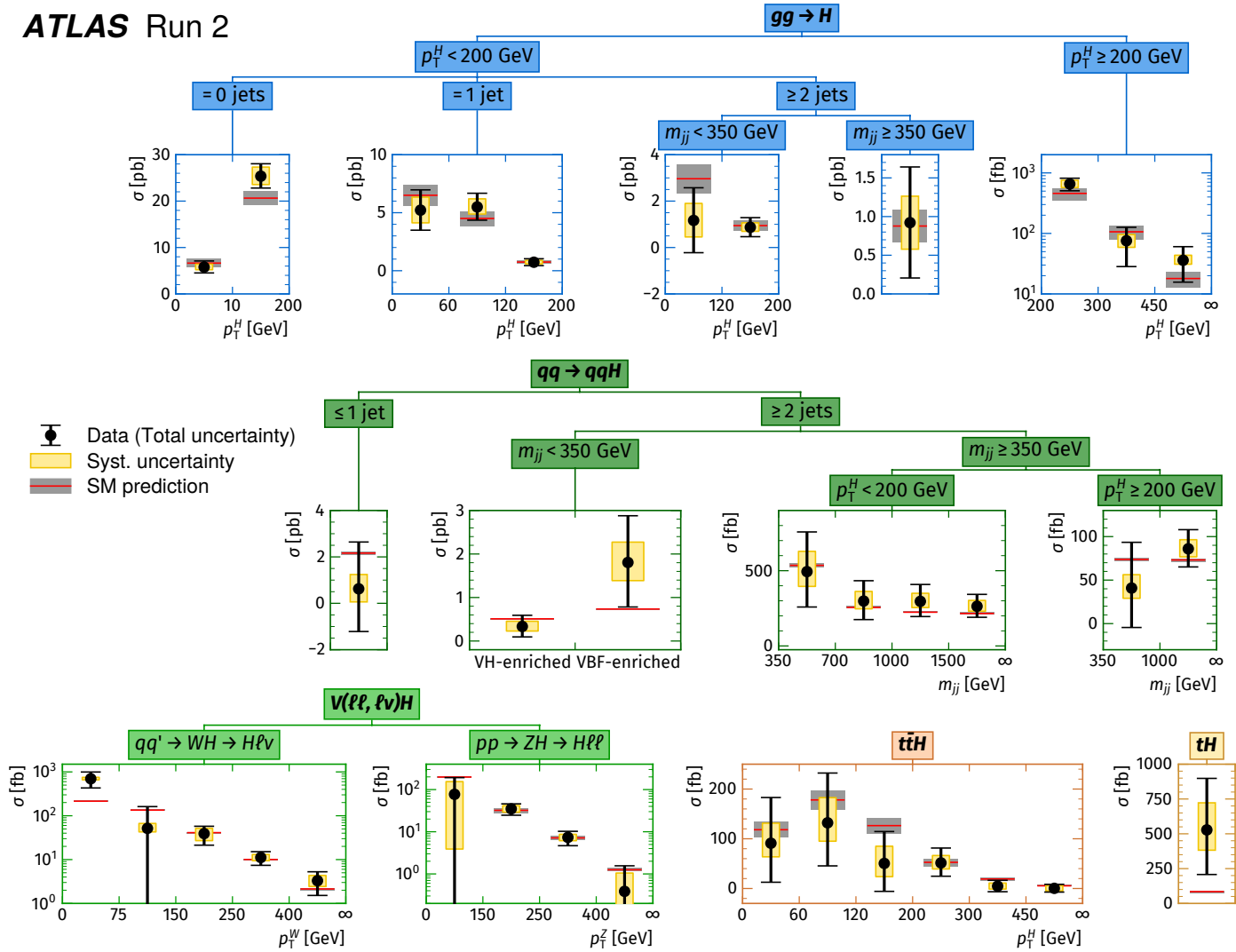


CMS

138 fb⁻¹ (13 TeV)



Nature 607 (2022) 60

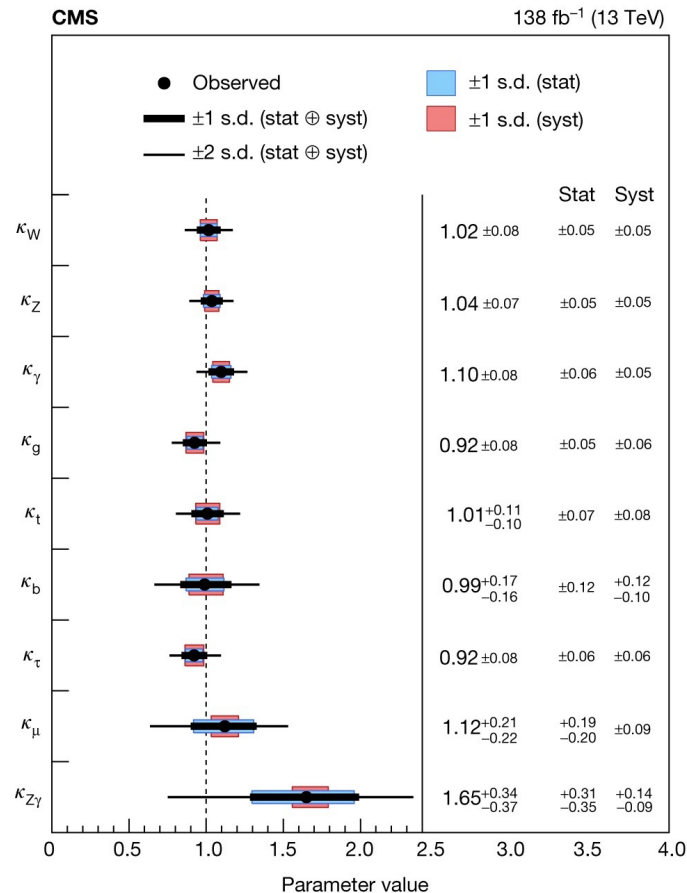
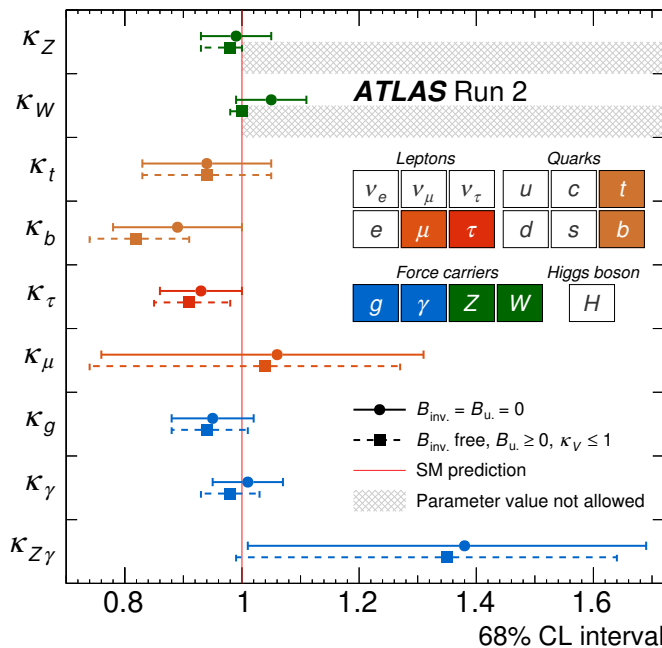


How well do we know the Higgs couplings?

- **Higgs couplings known to ~6-15%**
 - ✓ Some channel not included yet, or with partial statistics
 - ✓ Expect more improvements in next ~2 years, also from ATLAS+CMS combination

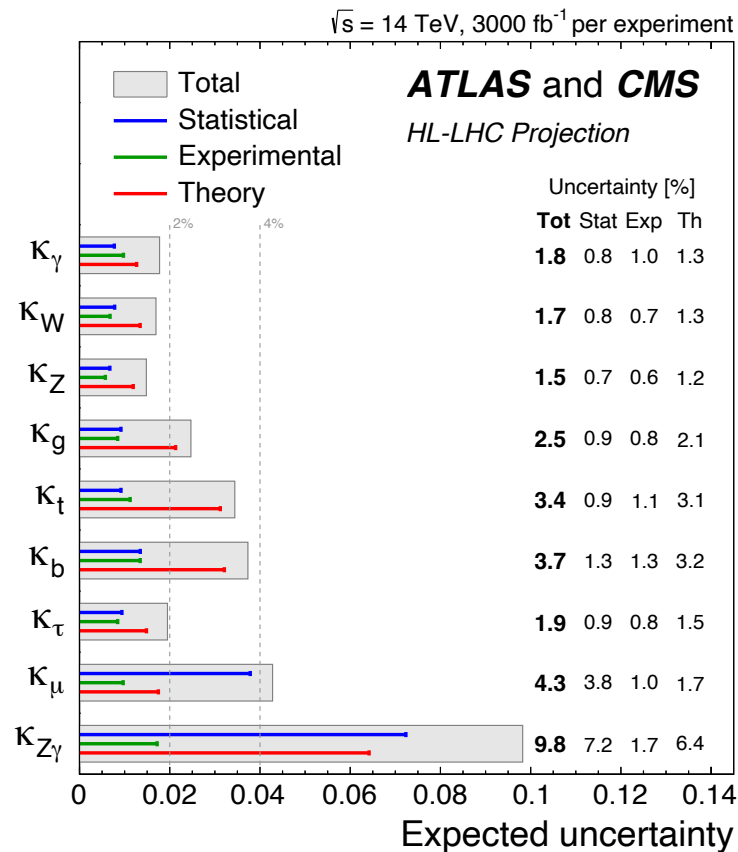
$$\kappa_j^2 = \sigma_j / \sigma_j^{\text{SM}}$$

$$\kappa_j^2 = \Gamma^j / \Gamma_{\text{SM}}^j$$



How well will we know the Higgs couplings?

- **2-4% precision** expected with $2 \times 3000 \text{ fb}^{-1}$



arXiv:1902.00134

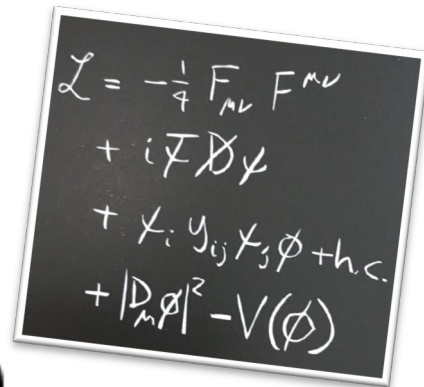
10.

Can the
Higgs boson
tell us something
about new phenomena?



Can Higgs properties constrain BSM phenomena today?

- Most recent approach: **Effective Field Theory (EFT)** interpretation



Handwritten Lagrangian on a chalkboard:

$$\mathcal{L} = -\frac{1}{4} F_{\mu\nu} F^{\mu\nu} + i\bar{\psi}\not{D}\psi + \bar{\psi}_i y_{ij} \psi_j \phi + h.c. + |D_\mu \phi|^2 - V(\phi)$$

$$\mathcal{L}_{\text{SMEFT}} = \mathcal{L}_{\text{SM}} + \sum_i^{N_{d6}} \frac{c_i}{\Lambda^2} \mathcal{O}_i^{(6)} + \sum_j^{N_{d8}} \frac{b_j}{\Lambda^4} \mathcal{O}_j^{(8)} + \dots,$$

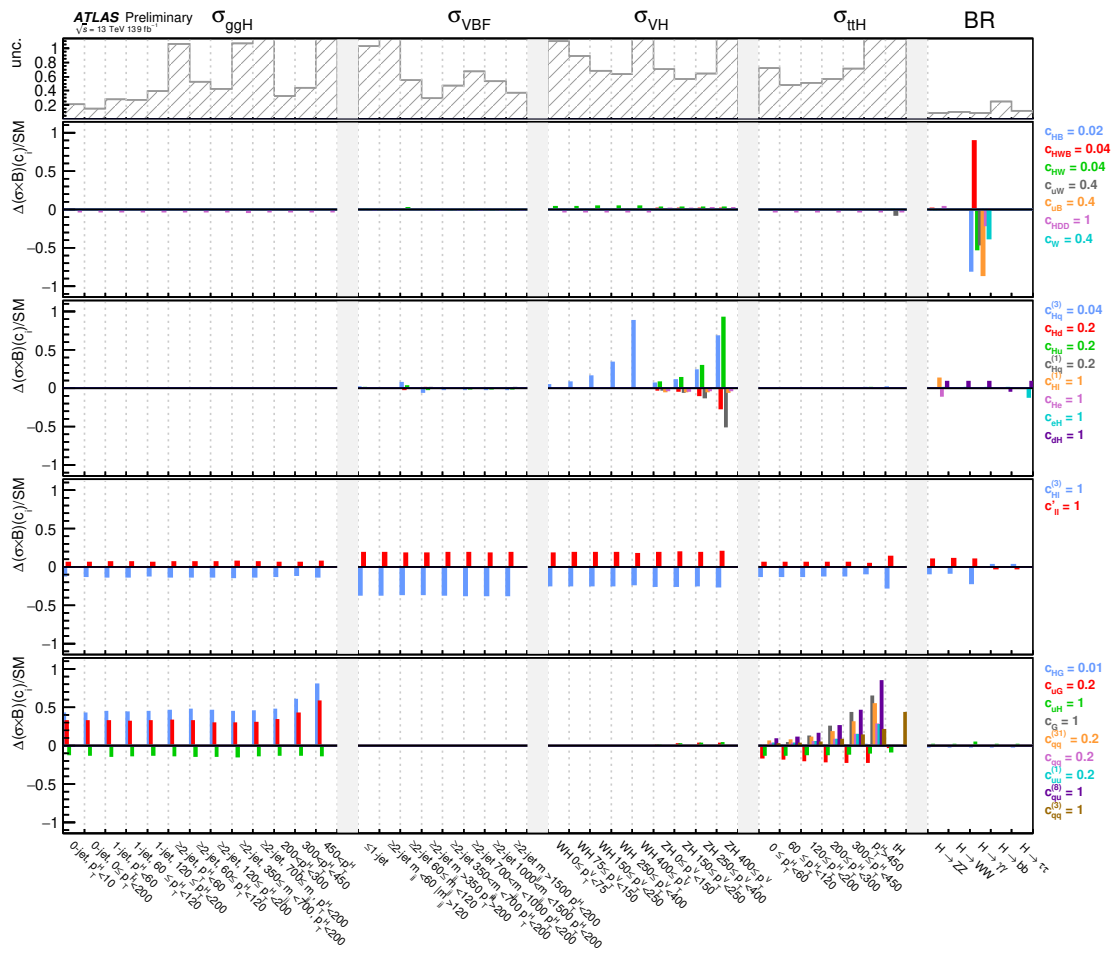
- $\mathcal{O}_i^{(n)}$ affect rates and kinematics
- STXS measurements and Higgs BR reparametrized in terms of dimension-six Wilson coefficients c_i in Warsaw basis
 - ✓ Correction for modified acceptance in $H \rightarrow 4l$ and $H \rightarrow l\nu l\nu$
 - ✓ U(3)5 flavor symmetry assumed, linearized in Wilson coefficients contribution (linear contribution $\propto \Lambda^{-2}$, expected to be leading)

Can Higgs properties constrain BSM phenomena today?

Wilson coefficient	Operator definition	Example diagram			
c_{HG}	$\Phi^\dagger \Phi G_{\mu\nu}^a G^{a\mu\nu}$		c_{Hl3}	$(i\Phi^\dagger \overleftrightarrow{D}_\mu^I \Phi)(\bar{\ell}\sigma^I \gamma^\mu \ell)$	
c_{HB}	$\Phi^\dagger \Phi B_{\mu\nu} B^{\mu\nu}$		c_{Hu}	$(i\Phi^\dagger \overleftrightarrow{D}_\mu^I \Phi)(\bar{u}\gamma^\mu u)$	
c_{HW}	$\Phi^\dagger \Phi W_{\mu\nu}^I W^{I\mu\nu}$		c_{Hd}	$(i\Phi^\dagger \overleftrightarrow{D}_\mu^I \Phi)(\bar{d}\gamma^\mu d)$	
c_{HWB}	$\Phi^\dagger \Phi W_{\mu\nu}^I B^{I\mu\nu}$		c_{He}	$(i\Phi^\dagger \overleftrightarrow{D}_\mu \Phi)(\bar{e}\gamma^\mu e)$	
c_{Hq1}	$(i\Phi^\dagger \overleftrightarrow{D}_\mu \Phi)(\bar{q}\gamma^\mu q)$		$ c_{uG} $	$(\bar{q}\sigma^{\mu\nu} T^a \Phi u) G_{\mu\nu}^a$	
c_{Hl1}	$(i\Phi^\dagger \overleftrightarrow{D}_\mu \Phi)(\bar{\ell}\gamma^\mu \ell)$		$ c_{eH} $	$(\Phi^\dagger \Phi)(\bar{\ell}e\Phi)$	
c_{Hq3}	$(i\Phi^\dagger \overleftrightarrow{D}_\mu^I \Phi)(\bar{q}\sigma^I \gamma^\mu q)$		$ c_{dH} $	$(\Phi^\dagger \Phi)(\bar{q}d\Phi)$	

EFT interpretation of Higgs couplings: c_i impact

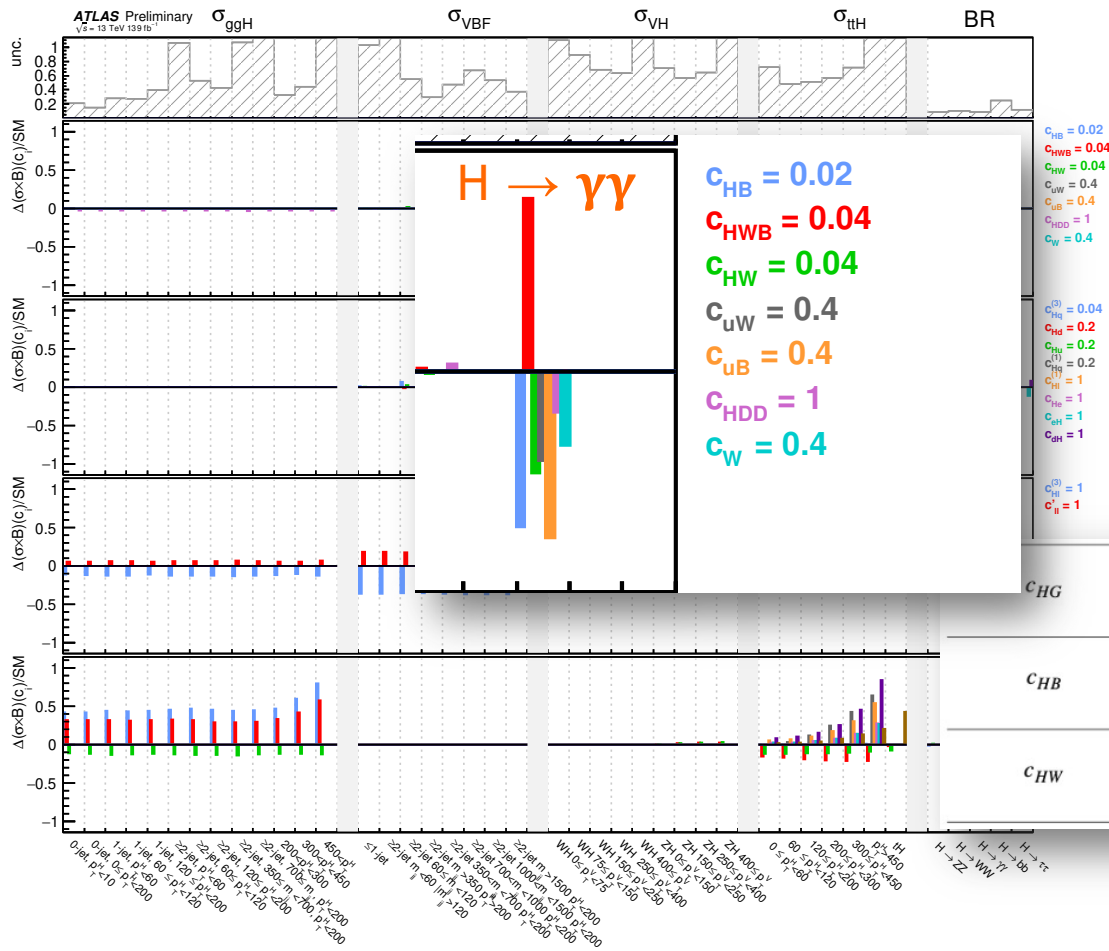
Measurement
uncertainty in STXS
bin / decay channel



Relative effect of
relevant operators

EFT interpretation of Higgs couplings: c_i impact

Measurement uncertainty in STXS bin / decay channel

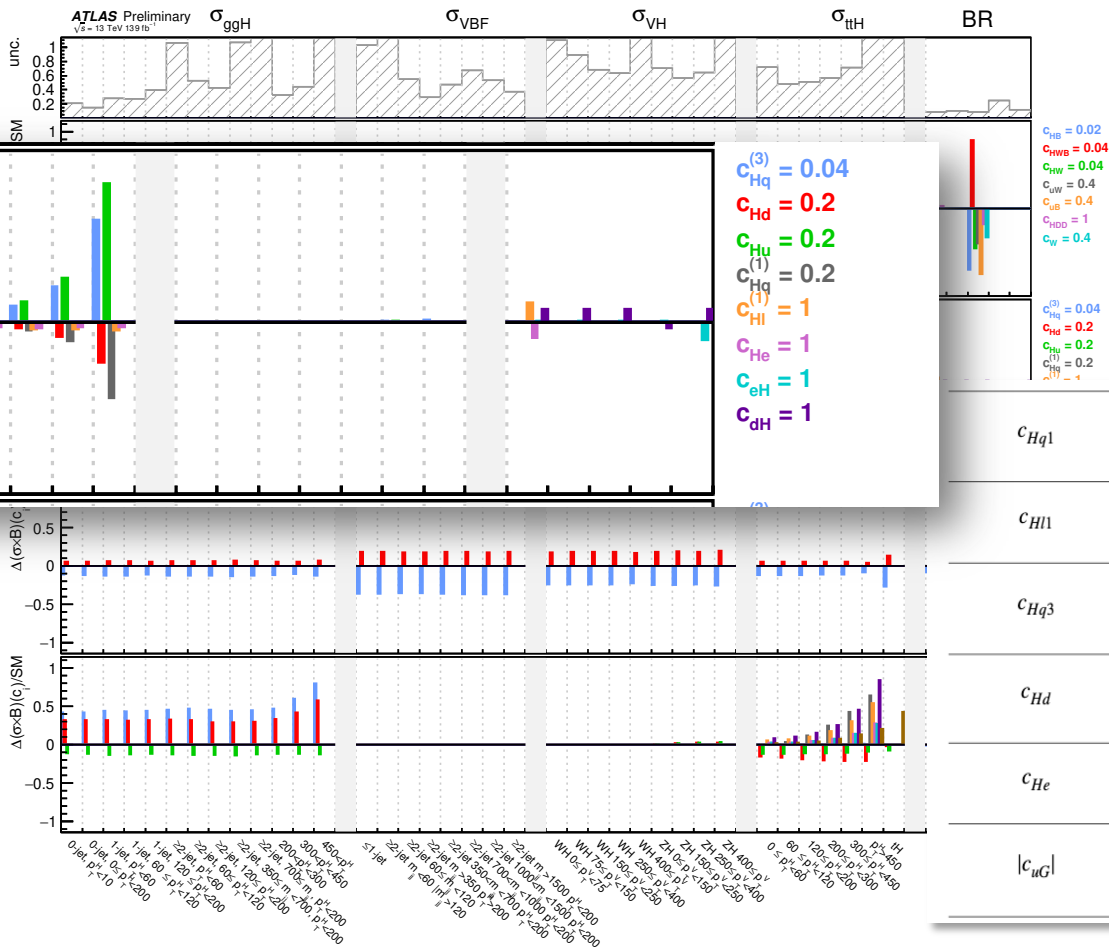


Relative effect of relevant operators

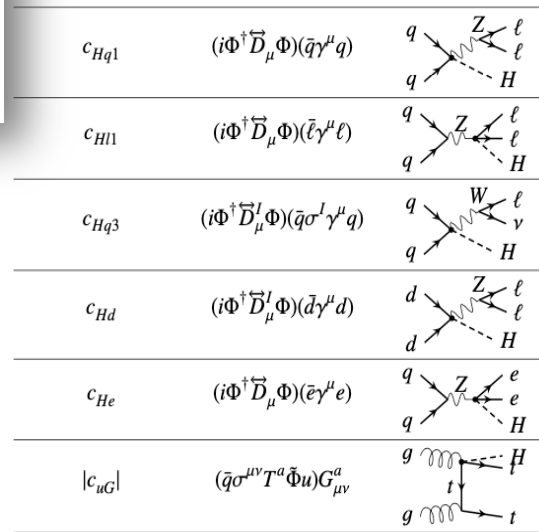
O(I) effects for small values of c_{HB} , c_{HW} , c_{HWB} (tree level $H \rightarrow \gamma\gamma$)

EFT interpretation of Higgs couplings: c_i impact

Measurement uncertainty in STXS bin / decay channel

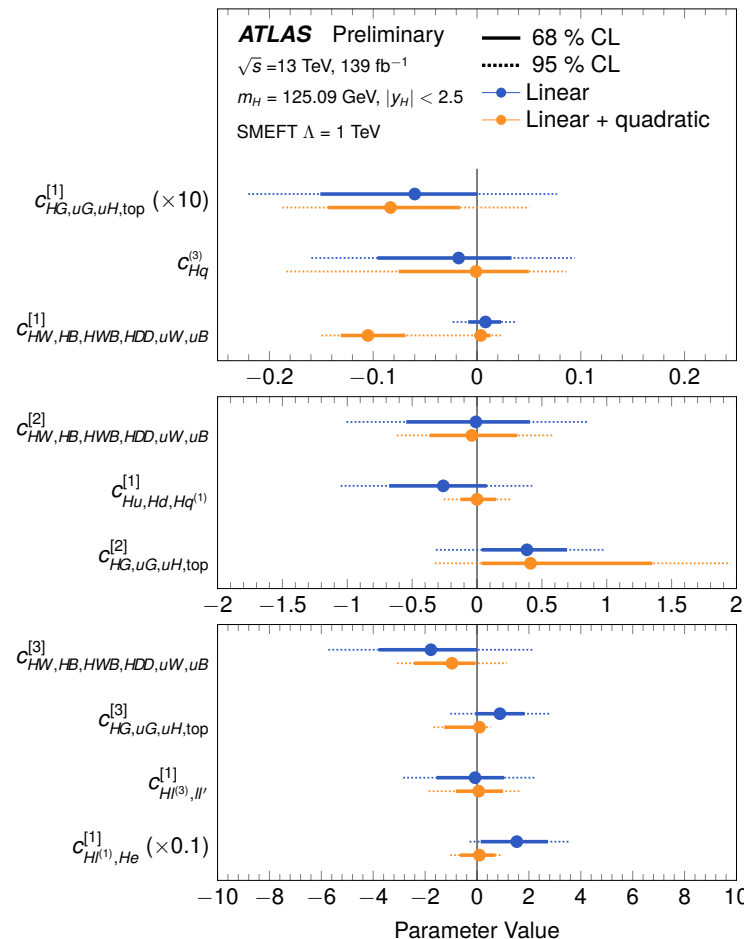


Strong energy growth of effect of $c^{(3)}_{Hg}$, $c^{(1)}_{Hg}$, c_{Hu} , c_{Hd}

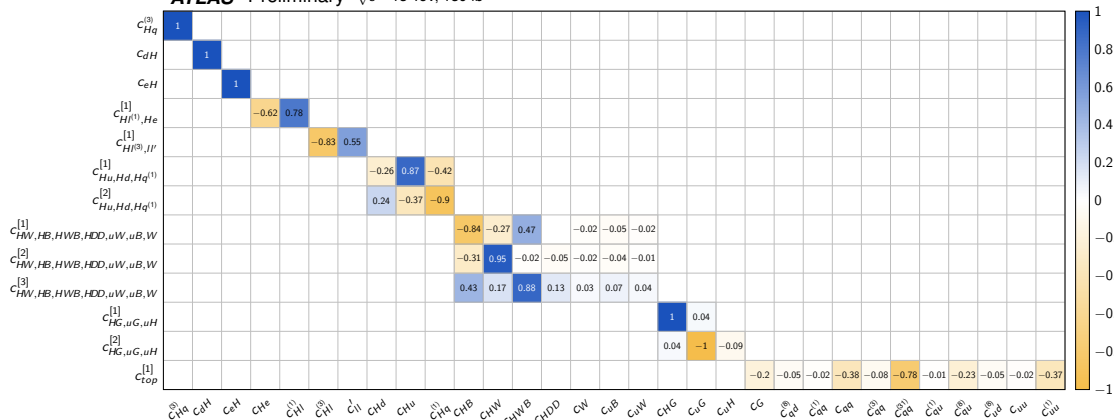


EFT interpretation of Higgs couplings

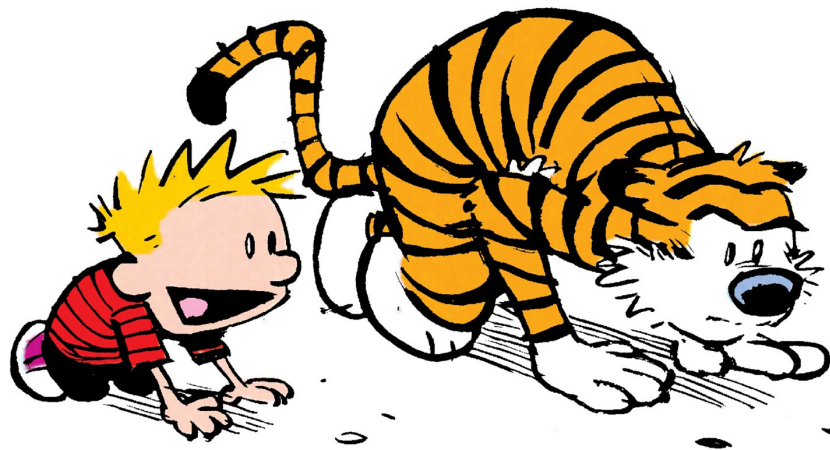
- Simultaneously constraining all coefficients impossible
- Fit performed in *sensitive directions*
 - ✓ linear combinations informed by principal component analysis
- Examples
 - ✓ $c^{[1]}_{HG,uG,uH}$: linear comb. with strong impact on $gg \rightarrow H$
 - ✓ $c^{[1]}_{HW,HB,HWB,HDD,uW,uB,W}$: strong impact on $H \rightarrow \gamma\gamma$
 - ✓ $c^{(3)}_{Hq}$: unique impact on VH



ATLAS Preliminary $\sqrt{s}=13$ TeV, 139 fb⁻¹

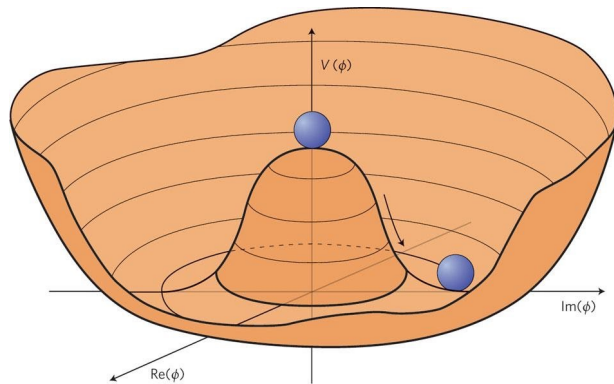


11



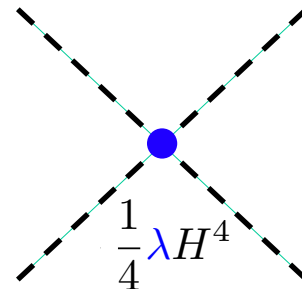
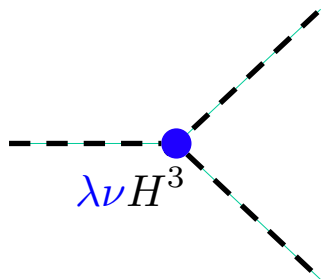
The SM
missing piece





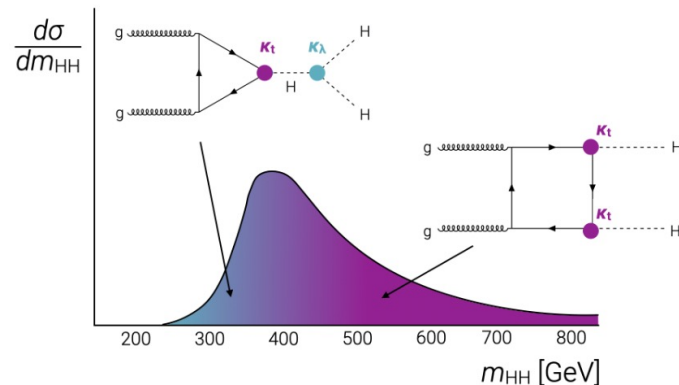
$$V(\Phi) = -\mu^2 \Phi^\dagger \Phi + \lambda (\Phi^\dagger \Phi)^2$$

$$V(\Phi) = V_0 + \frac{1}{2} m_H^2 H^2 + \lambda \nu H^3 + \frac{1}{4} \lambda H^4$$

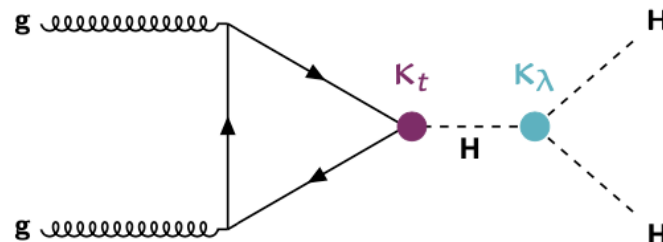
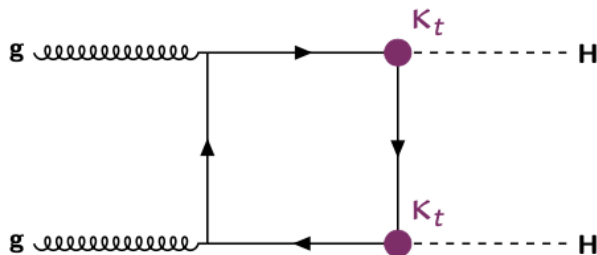


Can we access the Higgs self-coupling at the LHC?

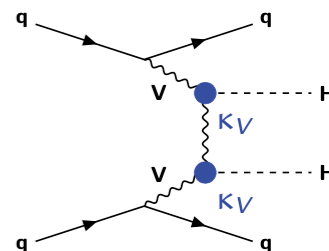
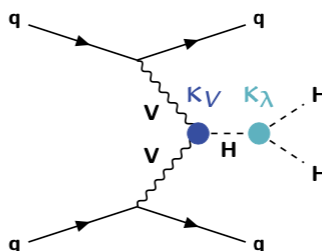
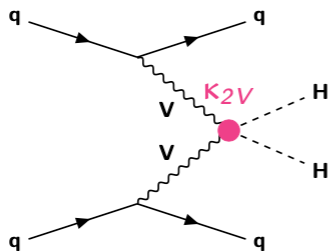
- Main avenues toward self-coupling measurement at LHC
 - ✓ Direct: HH production
 - ✓ Indirect: constraints from single Higgs measurements
- HH production
 - ✓ Very low cross-section ($\sim 1000\times$ smaller than single Higgs!)
 - ✓ Destructive interference between $gg \rightarrow HH$ processes



gluon-fusion
 $\sigma_{ggF}(pp \rightarrow HH) = 31.05 \text{ fb}$



VBF
 $\sigma_{VBF}(pp \rightarrow HH) = 1.726 \text{ fb}$

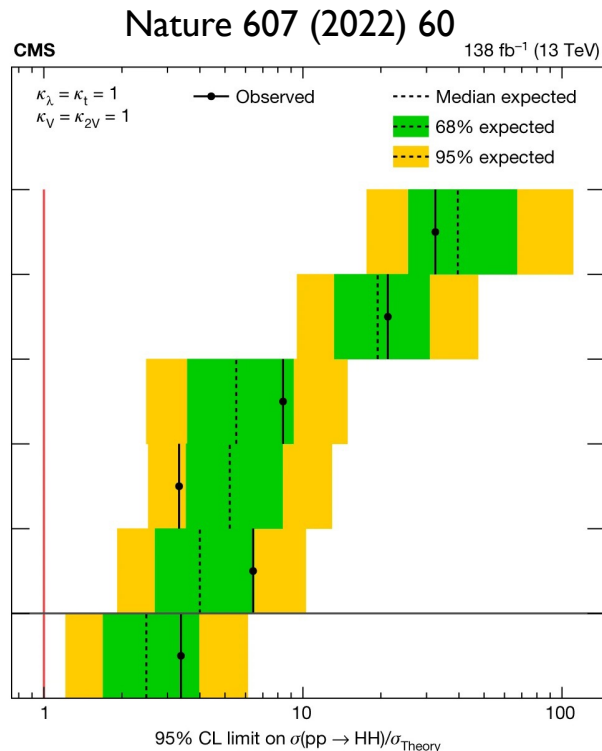
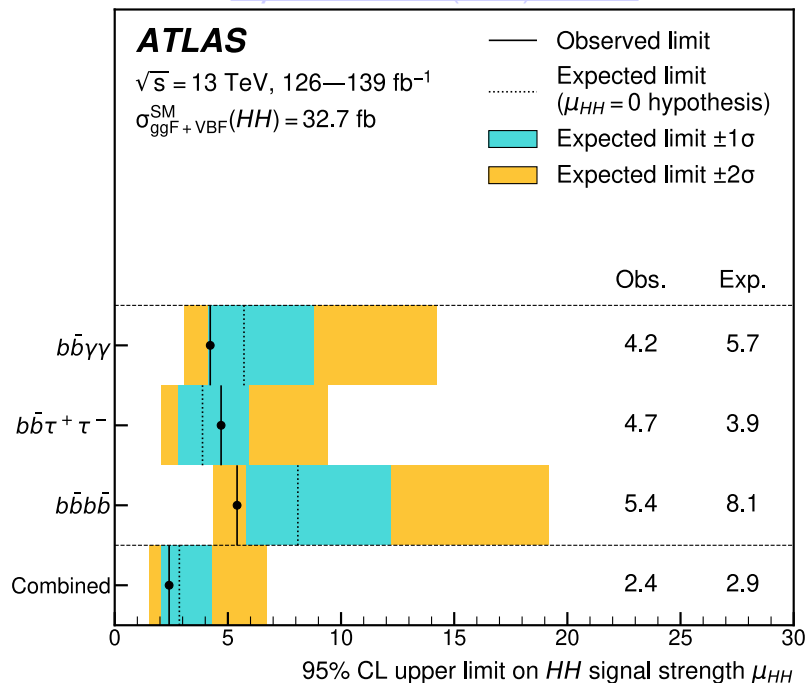


How to measure di-Higgs production?

	bb	WW	$\tau\tau$	ZZ	$\Upsilon\Upsilon$
bb	34%				
WW	25%	4.6%			
$\tau\tau$	7.3%	2.7%	0.39%		
ZZ	3.1%	1.1%	0.33%	0.069%	
$\Upsilon\Upsilon$	0.26%	0.10%	0.028%	0.012%	0.0005%

Closing in on di-Higgs production!

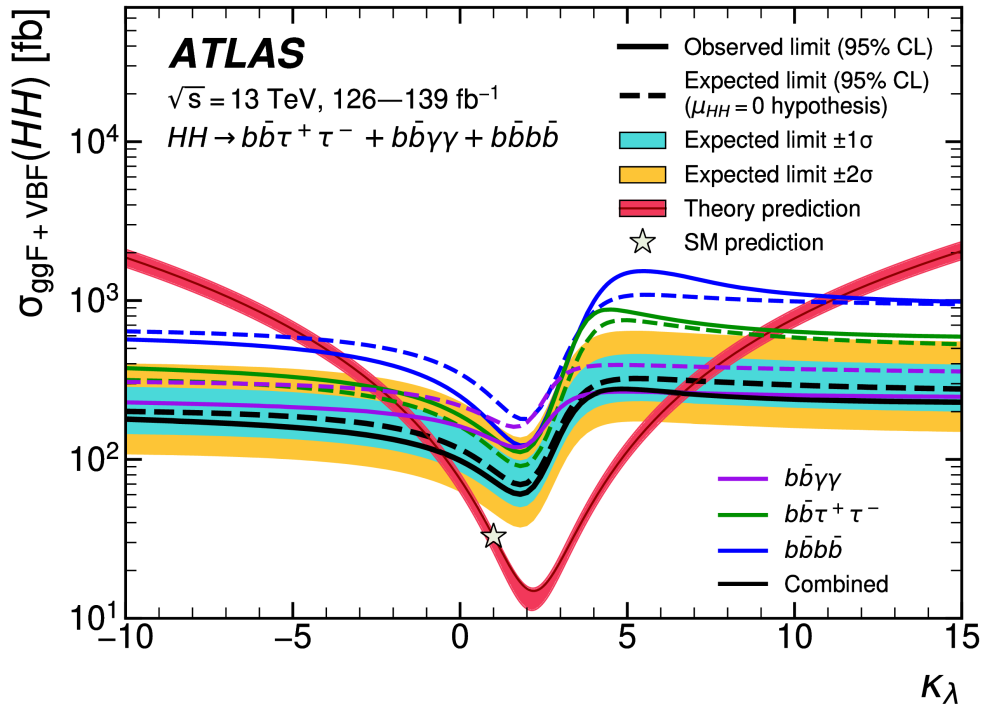
Phys. Lett. B 843 (2024) 137745



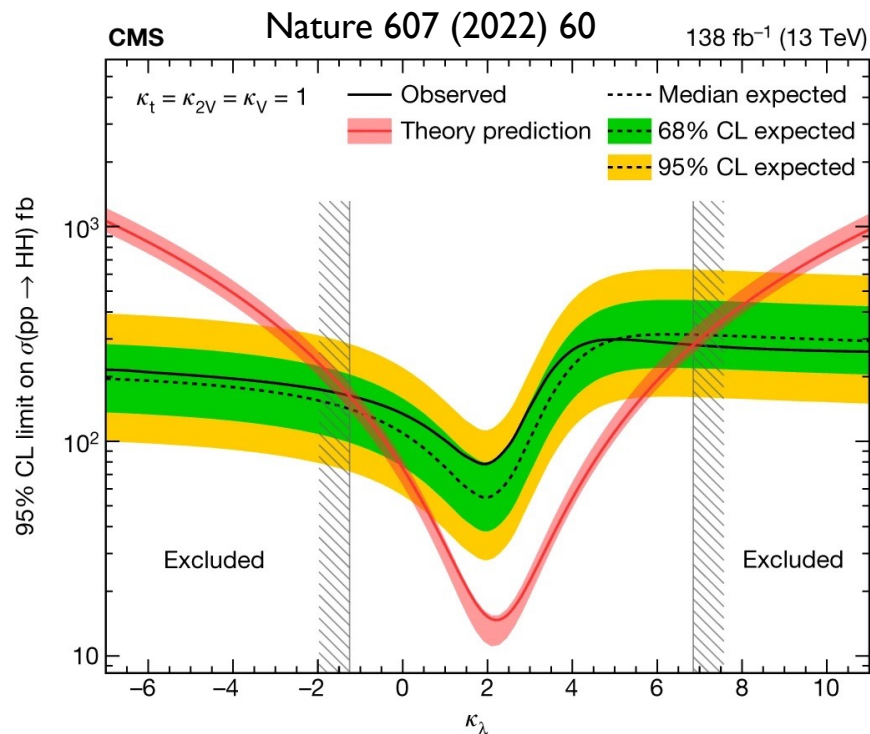
- Limits on κ_λ are improving dramatically, and there is more to be exploited (VBF production, combination of all channels and between ATLAS and CMS, ...)
- Di-Higgs studies will be a central aspect of the Run 3 (and of course HL-LHC)!

Closing in on di-Higgs production!

arXiv:2211.01216



$-0.4 < \kappa_\lambda < 6.3 @ 95\% \text{ CL}$



$-1.24 < \kappa_\lambda < 6.49 @ 95\% \text{ CL}$

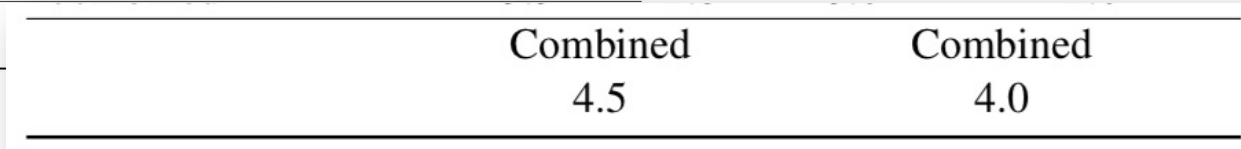
- Di-Higgs studies will be a central aspect of the Run 3 (and of course HL-LHC)!

When will we find the missing piece of the SM?

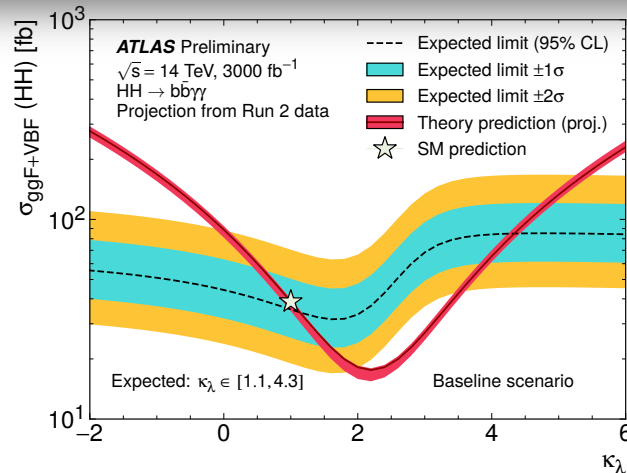
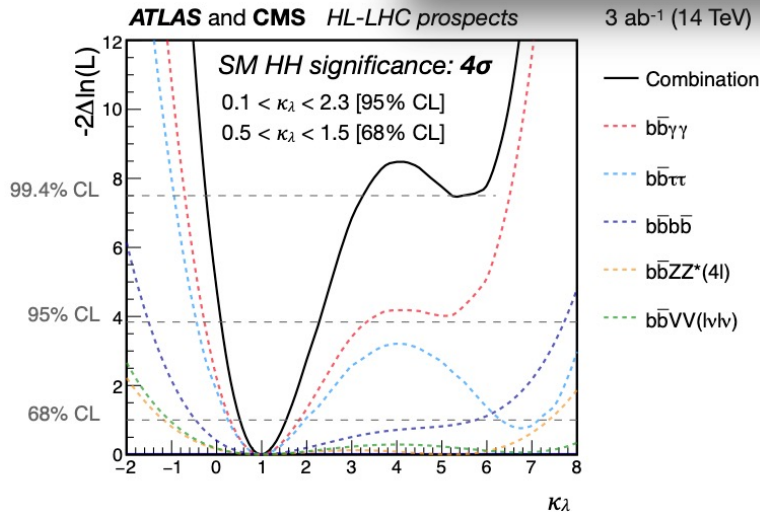
	Statistical-only		Statistical + Systematic	
	ATLAS	CMS	ATLAS	CMS
$HH \rightarrow b\bar{b}b\bar{b}$	1.4	1.2	0.61	0.95
$HH \rightarrow b\bar{b}\tau\tau$	2.5	1.6	2.1	1.4
$HH \rightarrow b\bar{b}\gamma\gamma$	2.1	1.8	2.0	1.8
$HH \rightarrow b\bar{b}VV(l\nu\nu)$	-	0.59	-	0.56
$HH \rightarrow b\bar{b}ZZ(4l)$	-	0.37	-	0.37
combined	3.5	2.8	3.0	2.6

A word of hope...

We often do better than our projections!



arXiv:1902.00134



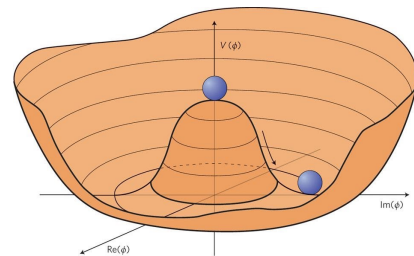
... and one of caution
 Sometimes projection are optimistic, and HL-LHC will present formidable challenges!

ATL-PHYS-PUB-2022-001

Twelve years with the Higgs boson



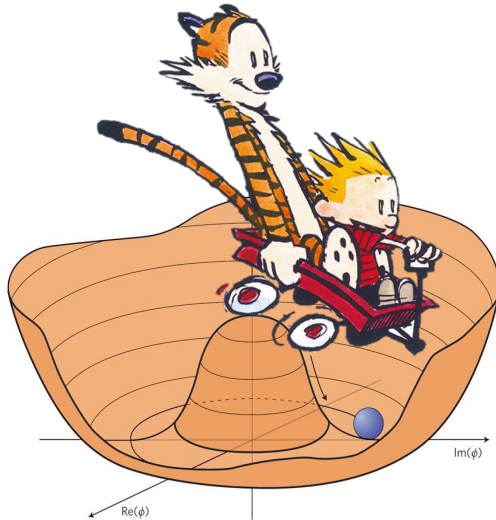
12.



The road ahead

We have a formidable tool to understand particle physics...

- **In 12 years we have measured with great precision many Higgs boson properties**
 - ✓ Mass, width, CP properties...
 - ✓ Coupling properties and differential distributions...
 - ✓ Closing in on coupling to 2nd generation, rare decays, self-interaction...
 - ✓ Higgs as a tool to constrain New Physics...
- **The LHC has restarted in 2022 for its Run 3 and will operate for a long time...**
 - ✓ A unique opportunity to continue characterizing the Higgs sector!

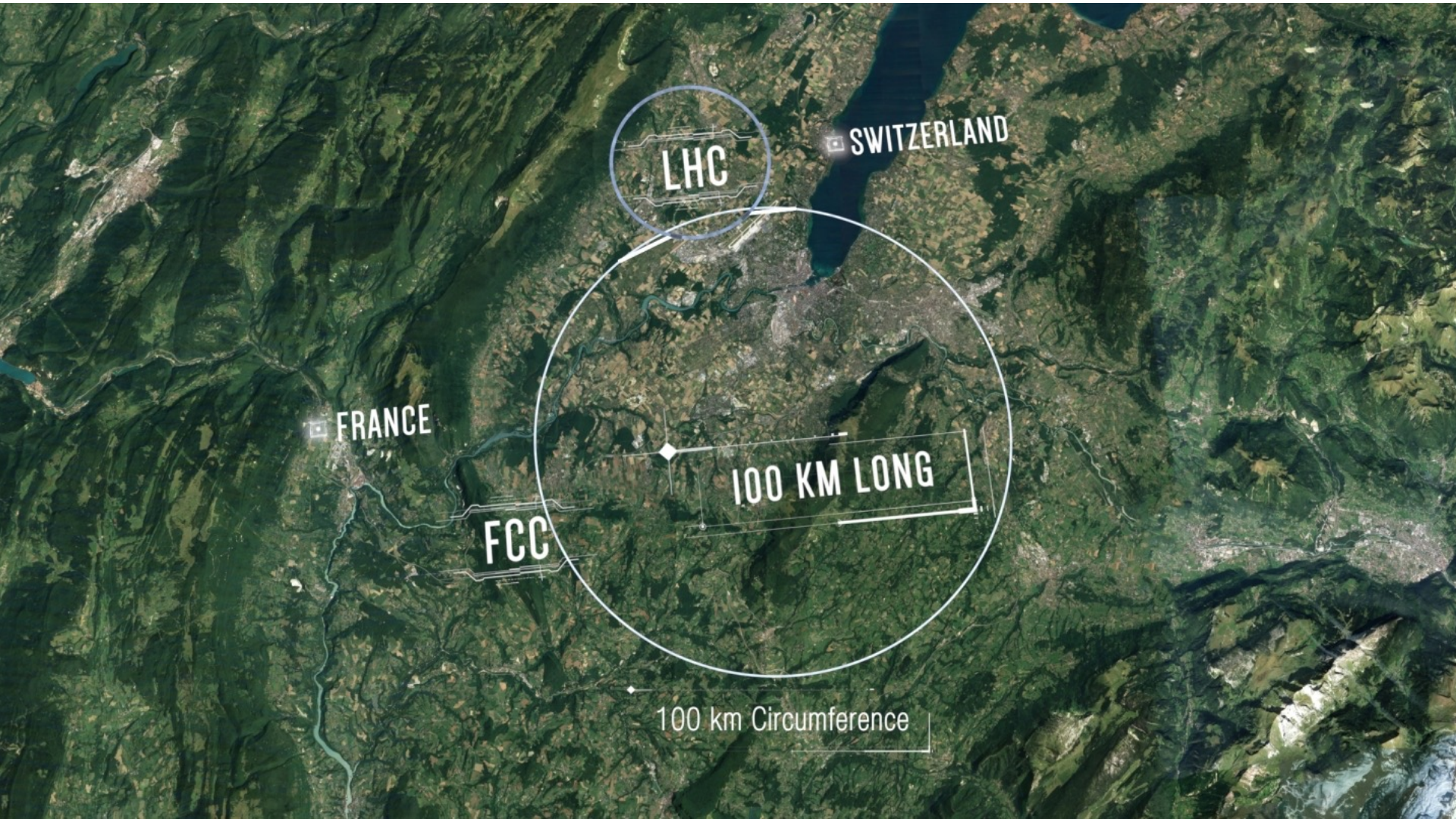


Every problem of the SM originates from Higgs interactions

$$\mathcal{L} = \lambda H \Psi \bar{\Psi} + \mu^2 |H|^2 - \lambda |H|^4 - V_0$$

↑ ↑ ↑ ↑
flavour naturalness stability C.C.

G. Giudice
via F. Gianotti



LHC

SWITZERLAND

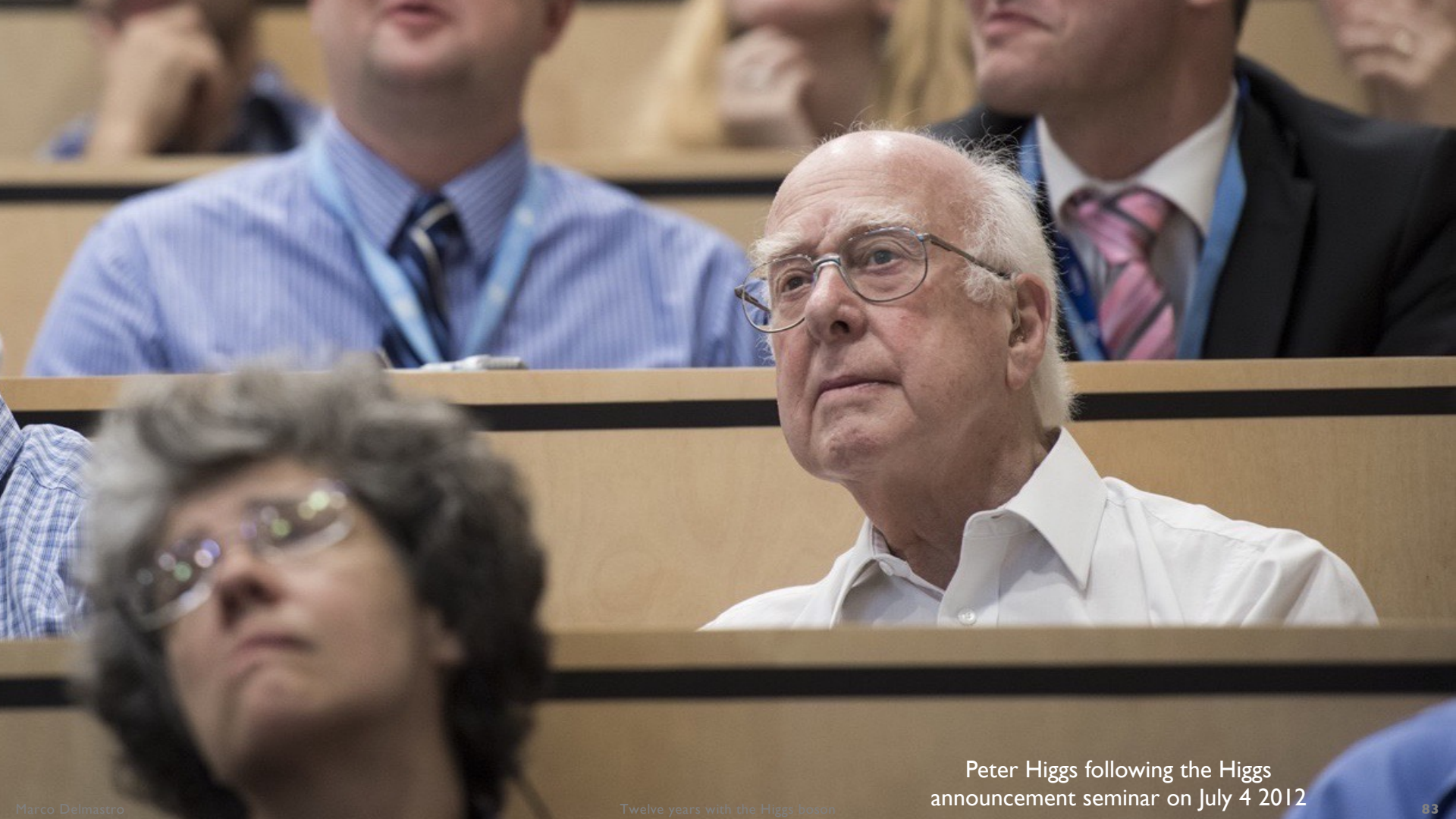
FRANCE

FCC

100 KM LONG

100 km Circumference



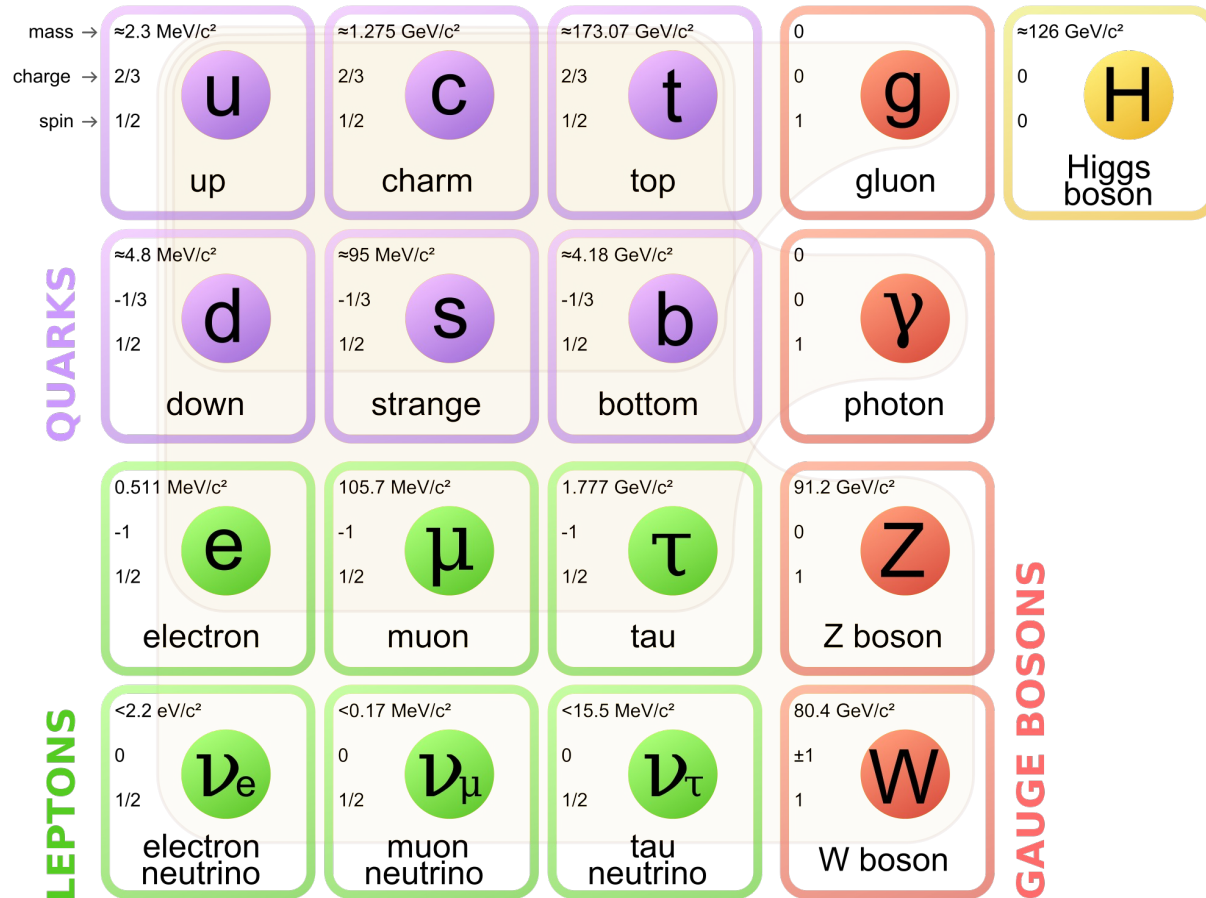


Peter Higgs following the Higgs
announcement seminar on July 4 2012

$e^{i\pi}$.

More material

The Standard Model



Higgs boson

Why is the electroweak interaction so much stronger than gravity?

- Are there new particles close to the mass of the Higgs boson?
- Is the Higgs boson elementary or made of other particles?
- Are there anomalies in the interactions of the Higgs boson with the W and Z bosons?

Why is there more matter than antimatter in the Universe?

- Are there charge-parity violating Higgs decays?
- Are there anomalies in the Higgs self-coupling that would imply a strong first-order early-Universe electroweak phase transition?
- Are there multiple Higgs sectors?

What is dark matter?

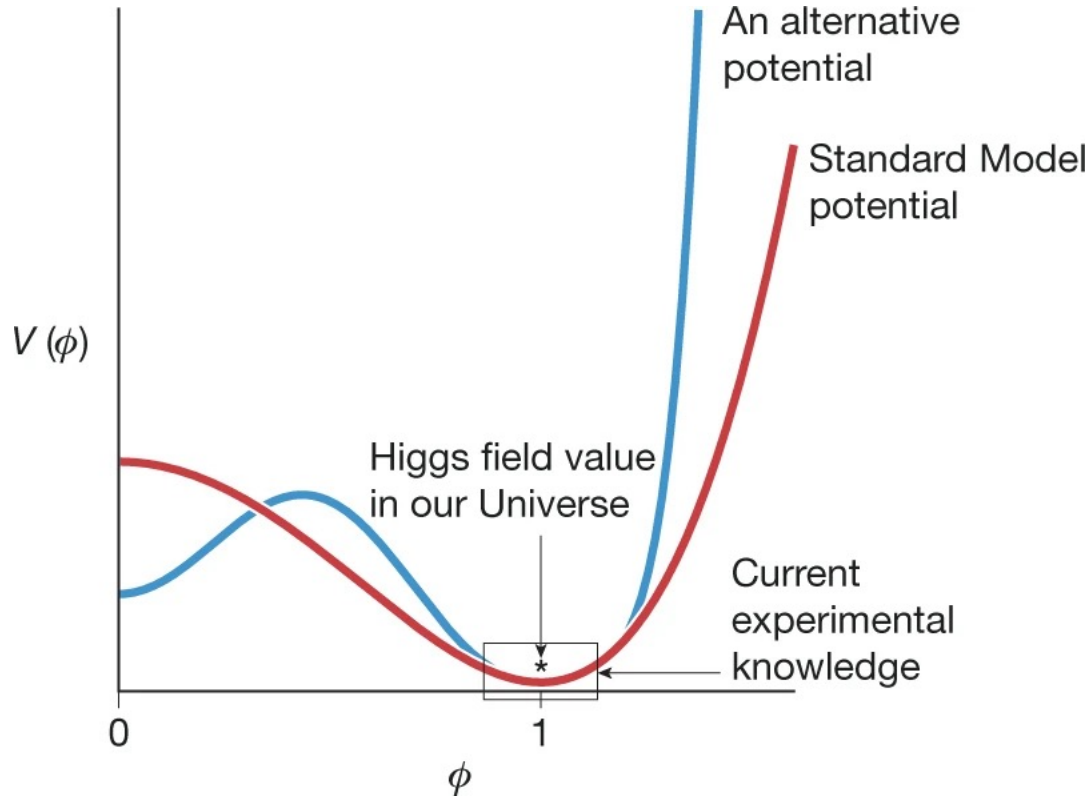
- Can the Higgs boson provide a portal to dark matter or a dark sector?
- Is the Higgs lifetime consistent with the Standard Model?
- Are there new decay modes of the Higgs boson?

What is the origin of the vast range of quark and lepton masses in the Standard Model?

- Are there modified interactions to the Higgs boson and known particles?
- Does the Higgs boson decay into pairs of quarks or leptons with distinct flavours (for example, $H \rightarrow \mu^+ \tau^-$)?

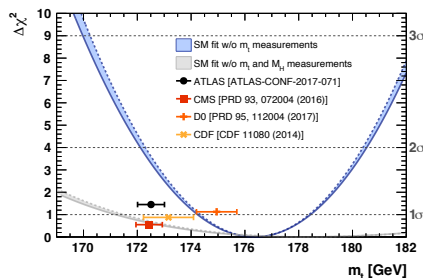
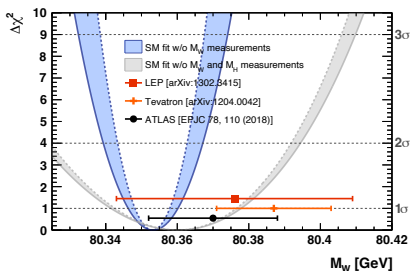
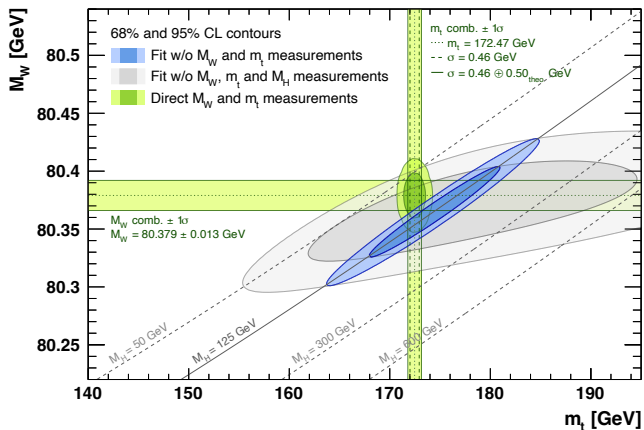
What is the origin of the early Universe inflation?

- Any imprint in cosmological observations?



Knowing the Higgs mass values...

... allows to make precision predictions



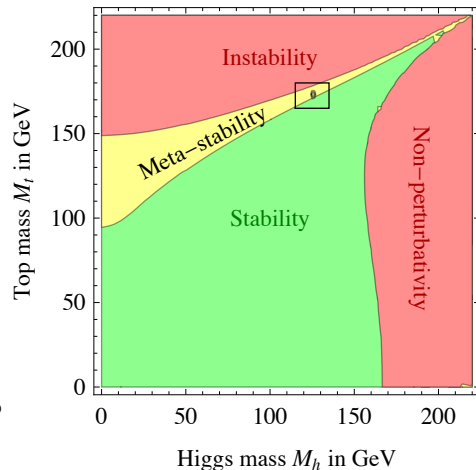
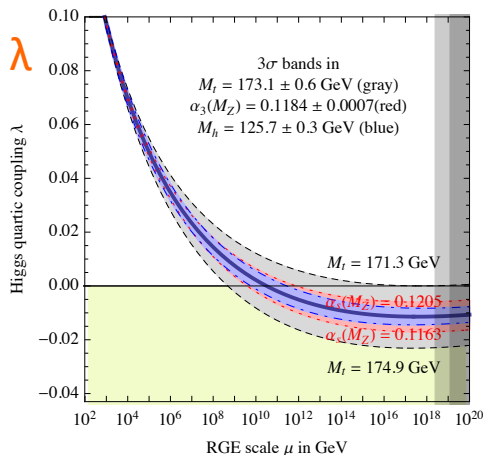
... but current m_H precision has little impact

... questions us on vacuum (meta) stability

$$V(\Phi) = -\mu^2 \Phi^\dagger \Phi + \lambda (\Phi^\dagger \Phi)^2$$

$$m_H = \sqrt{2\lambda} \nu$$

Running of the Higgs self coupling, assuming SM only at high scale

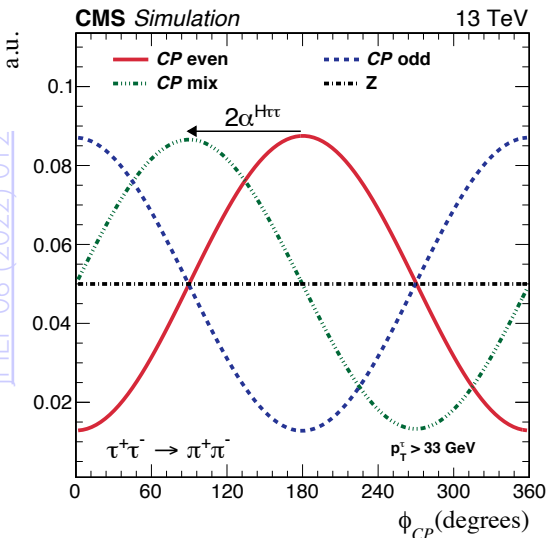
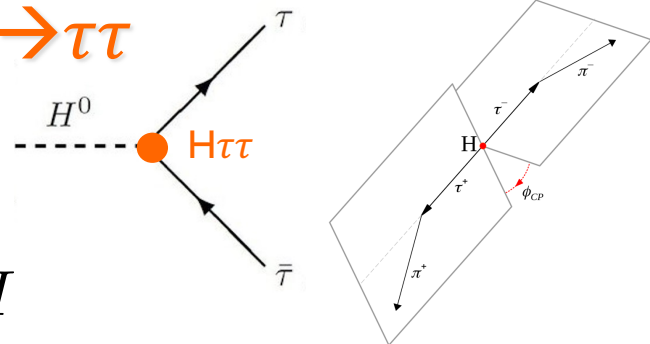


Vanishing quartic coupling at the Planck scale? Potential not bounded?
... but interpretation more sensitive to precision on top quark mass

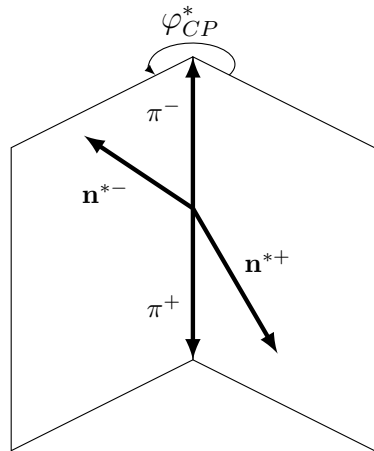
CP properties of Higgs- τ coupling with $H \rightarrow \tau\tau$

Effective Lagrangian for Yukawa coupling to tau leptons parameterized by **CP-Even** and **CP-odd** components

$$\mathcal{L}_{H\tau\tau} = -\frac{m_\tau}{\nu} \kappa_\tau \left(\cos \phi_\tau \bar{\tau}\tau + \sin \phi_\tau \bar{\tau}i\gamma_5\tau \right) H$$



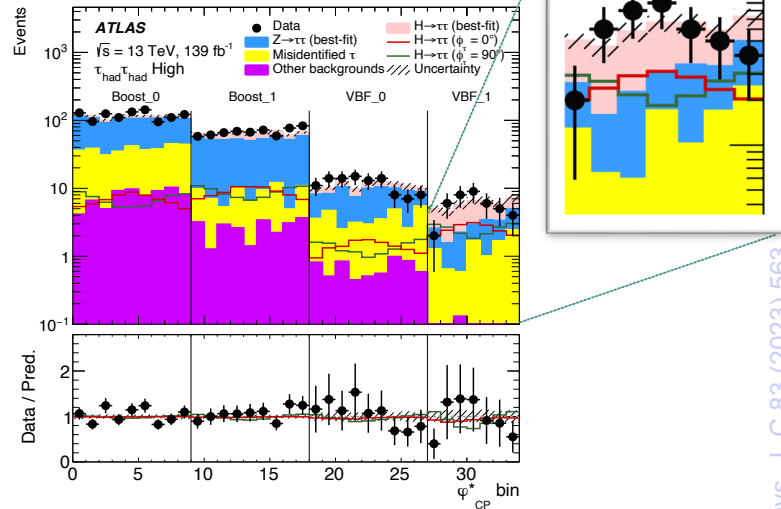
$$\alpha^{H\tau\tau}(\text{CMS}) = \phi_\tau(\text{ATLAS})$$



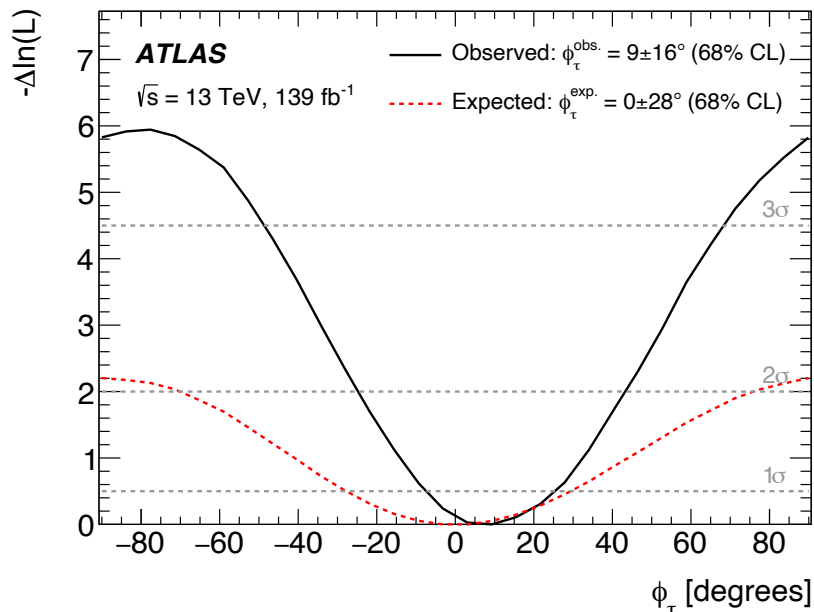
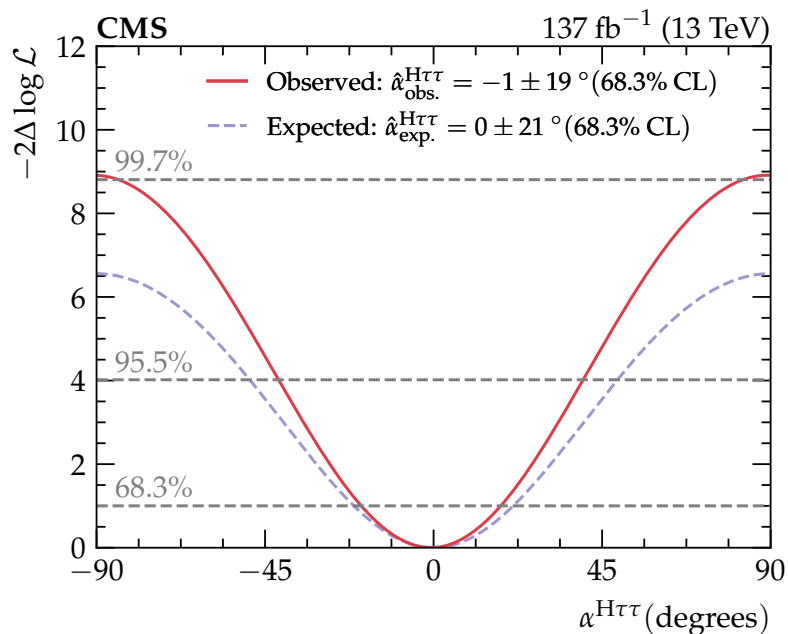
$$H \rightarrow \tau^+\tau^- \rightarrow \pi^+\pi^- + 2\nu$$

$$\phi^* = \arccos(\hat{n}_\perp^{*+} \cdot \hat{n}_\perp^{*-})$$

Uses either the impact parameter direction for single-prong taus (τ^\pm) or τ^\pm momentum for $\tau \rightarrow \tau^\pm \nu_\tau$ (also for 3-prong decays with $\rho \rightarrow \tau^\pm \nu_\tau$)



CP properties of Higgs- τ coupling with $H \rightarrow \tau\tau$



$$\phi_\tau = -1 \pm 19^\circ (21^\circ \text{ exp})$$

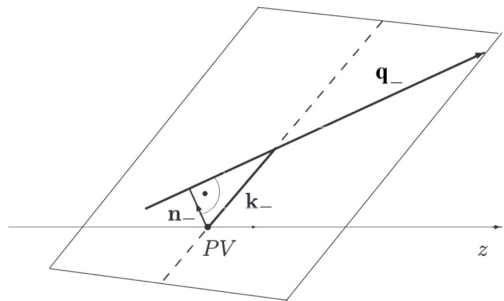
$$\alpha^{H\tau\tau}(\text{CMS}) = \phi_\tau(\text{ATLAS})$$

$$\phi_\tau = 9 \pm 16^\circ (28^\circ \text{ exp})$$

pure CP-odd coupling
excluded at 3σ

pure CP-odd coupling
excluded at 3.4σ

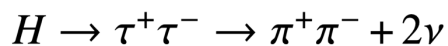
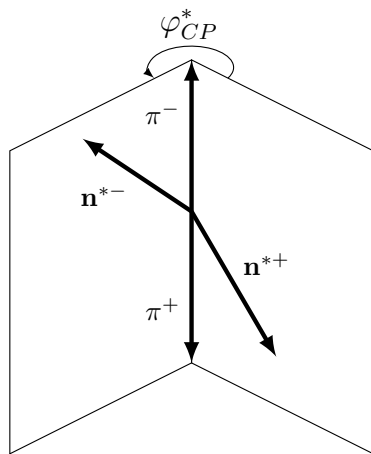
H → ττ decay CP



Impact parameter

directional distance of closest approach of charged particle's track to reconstructed PV of the event

4-vectors boosted to the rest frame of visible di-τ Zero Momentum Frame (e.g. two decay charged particles)

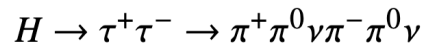
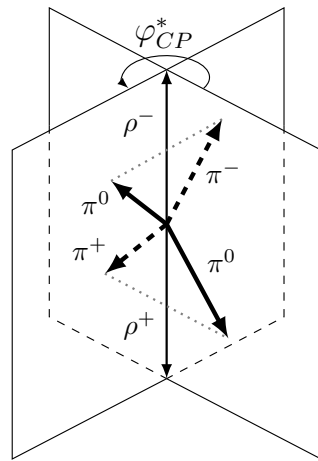


impact parameter

$$\varphi^* = \arccos(\hat{\mathbf{n}}_{\perp}^{*+} \cdot \hat{\mathbf{n}}_{\perp}^{*-})$$

$$O_{CP}^* = \hat{\mathbf{q}}^{*-} \cdot (\hat{\mathbf{n}}_{\perp}^{*+} \times \hat{\mathbf{n}}_{\perp}^{*-})$$

$$\varphi_{CP}^* = \begin{cases} \varphi^* & \text{if } O_{CP}^* \geq 0 \\ 360^\circ - \varphi^* & \text{if } O_{CP}^* < 0 \end{cases}$$

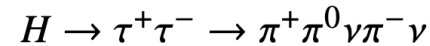
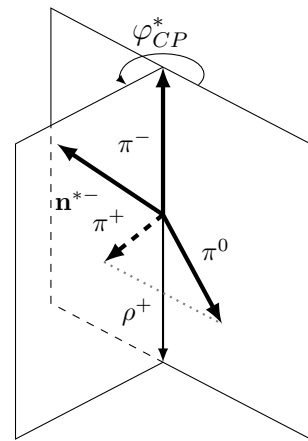


ρ decay plane

$$\varphi^* = \arccos(\hat{\mathbf{q}}_{\perp}^{*0+} \cdot \hat{\mathbf{q}}_{\perp}^{*0-})$$

$$O_{CP}^* = \hat{\mathbf{q}}^{*-} \cdot (\hat{\mathbf{q}}_{\perp}^{*0+} \times \hat{\mathbf{q}}_{\perp}^{*0-})$$

$$\varphi^{*'} = \begin{cases} \varphi^* & \text{if } O_{CP}^* \geq 0 \\ 360^\circ - \varphi^* & \text{if } O_{CP}^* < 0 \end{cases} \quad \varphi_{CP}^{*'} = \begin{cases} \varphi^{*'} & \text{if } y_+^{\rho}y_-^{\rho} \geq 0 \\ \varphi^{*'} + 180^\circ & \text{if } y_+^{\rho}y_-^{\rho} < 0 \end{cases} \quad y_{\pm}^{\rho} = \frac{E_{\pi^{\pm}} - E_{\pi^0}}{E_{\pi^{\pm}} + E_{\pi^0}}$$



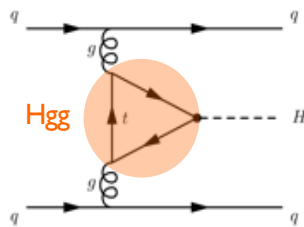
impact parameter +
ρ decay plane

CP properties of Higgs-gluon coupling with ggF+VBF

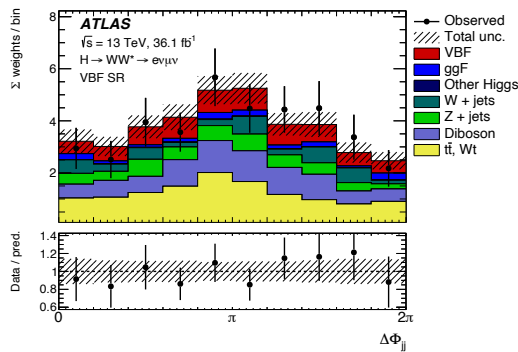
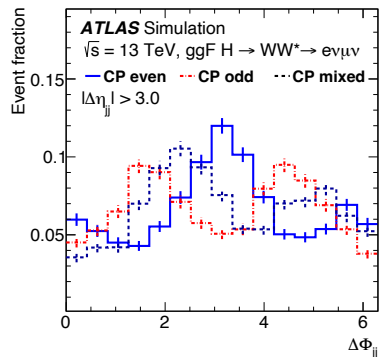
Study ggF+2 jets events: jet kinematics sensitive to Hgg coupling CP properties

ATLAS: $\Delta\Phi_{jj}$ ggF+VBF $H \rightarrow WW^*$

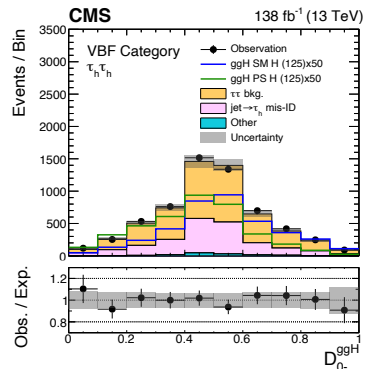
CMS: $\Delta\Phi_{jj}$ or MELA discriminant ggF+VBF $H \rightarrow \tau\tau$



arXiv:2109.13808



arXiv:2205.05120

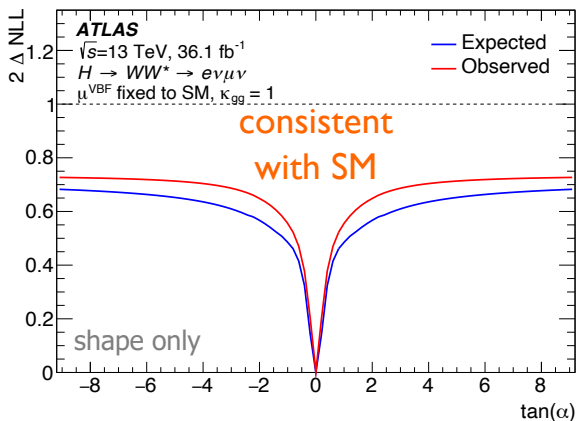


$$D_{2jet}^{VBF} = \frac{\mathcal{P}_{SM}^{ggH} + \mathcal{P}_{0-}^{ggH}}{\mathcal{P}_{SM}^{ggH} + \mathcal{P}_{0-}^{ggH} + \mathcal{P}_{SM}^{VBF}} \quad \text{ggF vs VBF}$$

$$D_{CP}^{ggH} = \frac{\mathcal{P}_{SM-0-}^{ggH}}{\mathcal{P}_{SM}^{ggH} + \mathcal{P}_{0-}^{ggH}} \quad \begin{array}{l} \text{ggF CP-even vs} \\ \text{ggF CP-odd} \end{array}$$

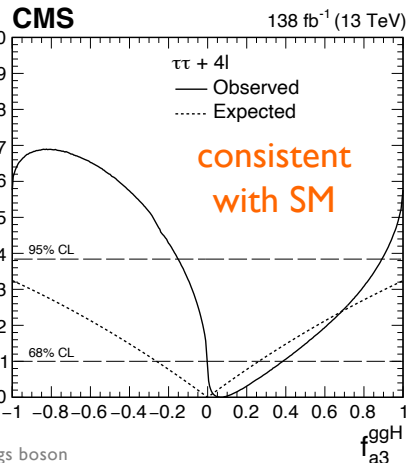
$$D_{0-}^{ggH} = \frac{\mathcal{P}_{SM}^{ggH}}{\mathcal{P}_{SM}^{ggH} + \mathcal{P}_{0-}^{ggH}} \quad \begin{array}{l} \text{CP-even CP-odd} \\ \text{interference} \end{array}$$

$H \rightarrow WW^*$



α parameterizes CP-even/CP-odd mixture on Hgg coupling

where α is the phase of the Higgs boson



$H \rightarrow \tau\tau$
(combined with $H \rightarrow 4l$)

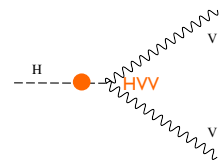
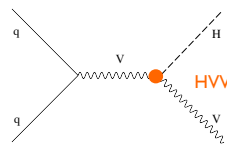
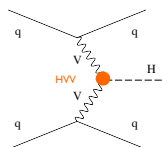
f_{a3}^{ggH} parameterizes fraction of CP-odd BSM point-like Hgg coupling

CP properties of HVV coupling (mult. prod. + $H \rightarrow 4l$)

HVV couplings parameterized by **tensor structures in scattering amplitude**

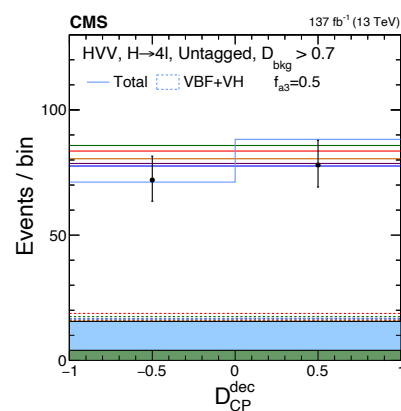
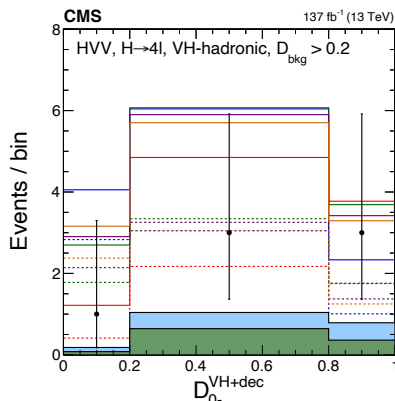
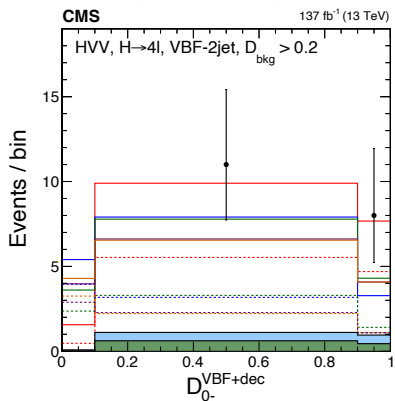
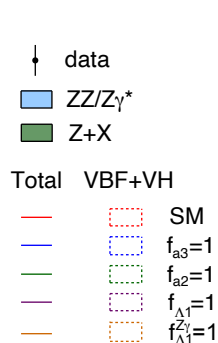
$$\mathcal{A}(HVV) \sim \left[a_1^{VV} + \frac{\kappa_1^{VV} q_1^2 + \kappa_2^{VV} q_2^2}{(\Lambda_1^{VV})^2} \right] m_{V1}^2 \epsilon_{V1}^* \epsilon_{V2}^* + a_2^{VV} f_{\mu\nu}^{*(1)} f^{*(2)\mu\nu} + a_3^{VV} f_{\mu\nu}^{*(1)} \tilde{f}^{*(2)\mu\nu}.$$

$$f_{a3} = \frac{|a_3|^2 \sigma_3}{|a_1|^2 \sigma_1 + |a_2|^2 \sigma_2 + |a_3|^2 \sigma_3 + |\kappa_1|^2 \sigma_{\Lambda 1} + |\kappa_1^{Z\gamma}|^2 \sigma_{\Lambda 1}^{Z\gamma}} \text{sgn} \left(\frac{a_3}{a_1} \right)$$

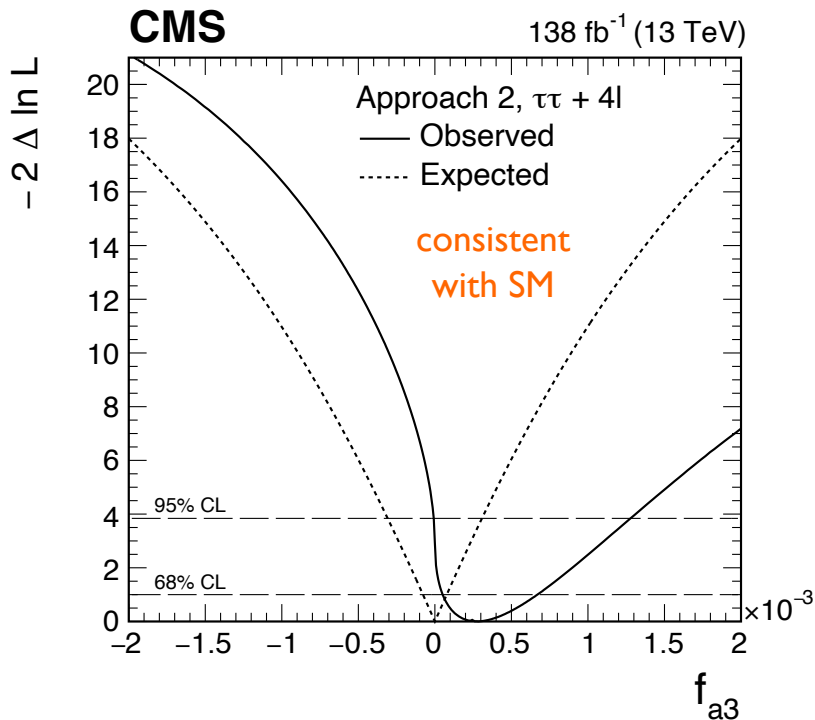


Multiple analyses constraining HVV couplings with ggH , VBF, VH, ttH , tH production and $H \rightarrow \tau\tau$, $H \rightarrow 4l$, $H \rightarrow \gamma\gamma$ decay

f_{a3} parameterizes fractional contribution of **CP-odd HZZ** coupling

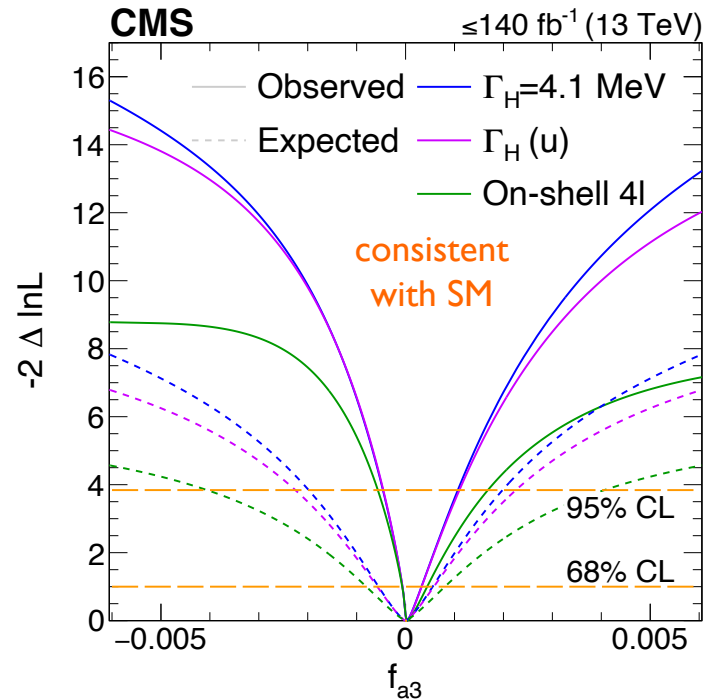


CP properties of HVV coupling (mult. prod. + $H \rightarrow 4l$; offshell)



Constraint from combined analyses
Other anomalous HVV fixed to 0

$$f_{a3} = 0.20^{+0.26}_{-0.16} \times 10^{-3}$$



Constraints from $H \rightarrow 4l$ onshell + off-shell
Other anomalous couplings floated in fits

$$f_{a3} = 0.024^{+0.32}_{-0.064} \times 10^{-3}$$

CP properties of HVV coupling with VBF

SM Lagrangian augmented with CP-odd dim-6 operators involving Higgs and EW gauge fields in EFT formalism

$$|\mathcal{M}|^2 = |\mathcal{M}_{SM}|^2 + \tilde{d} \cdot 2\text{Re}(\mathcal{M}_{SM}^* \mathcal{M}_{CP\text{-odd}}) + \tilde{d}^2 \cdot |\mathcal{M}_{CP\text{-odd}}|^2$$

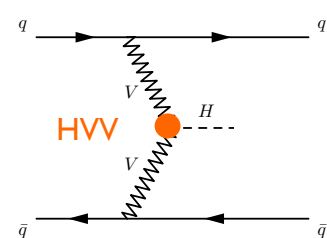
Approach pursued with various decay modes, here most recent ATLAS $H \rightarrow \gamma\gamma$, then combined with $36 \text{ fb}^{-1} H \rightarrow \tau\tau$

$$c_{H\tilde{W}} = c_{H\tilde{B}}$$

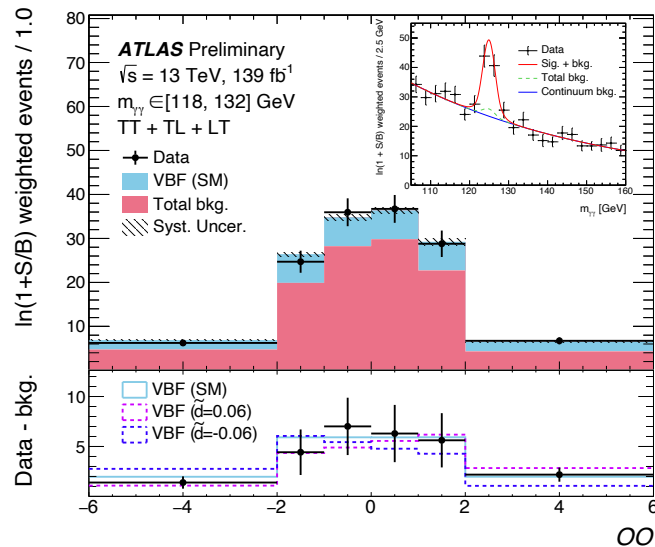
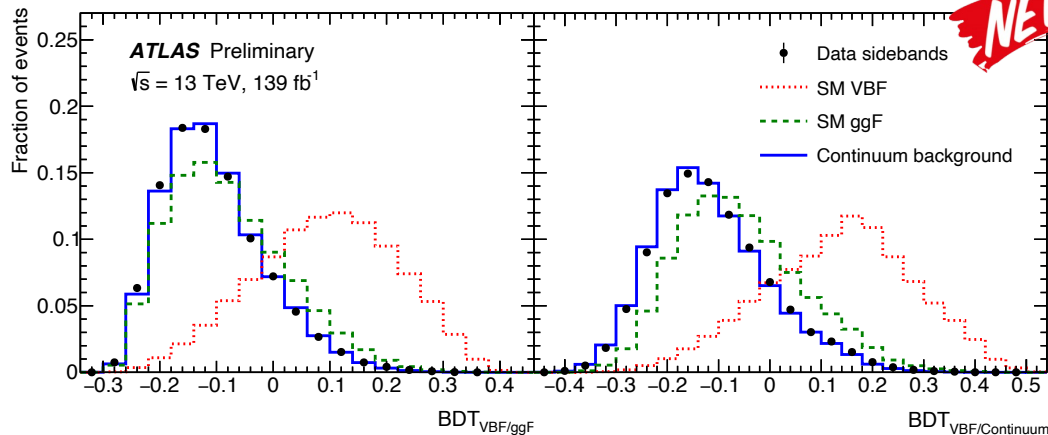
$$c_{HW\tilde{B}} = 0$$

$$\tilde{d} = \frac{\nu^2}{\Lambda^2} c_{H\tilde{W}}$$

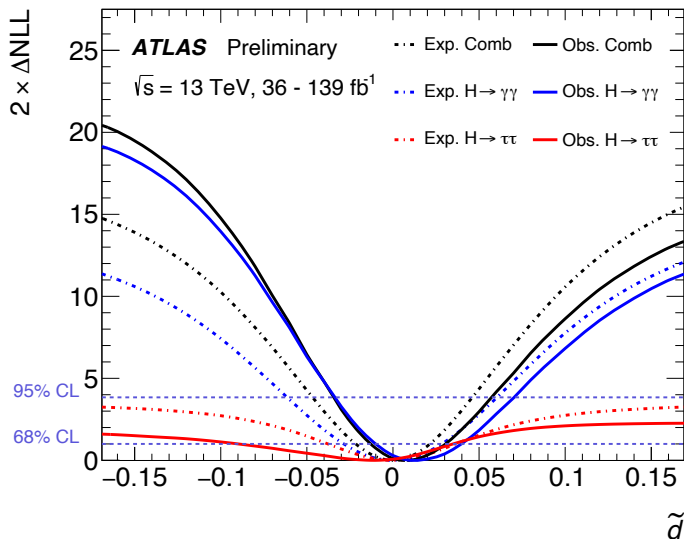
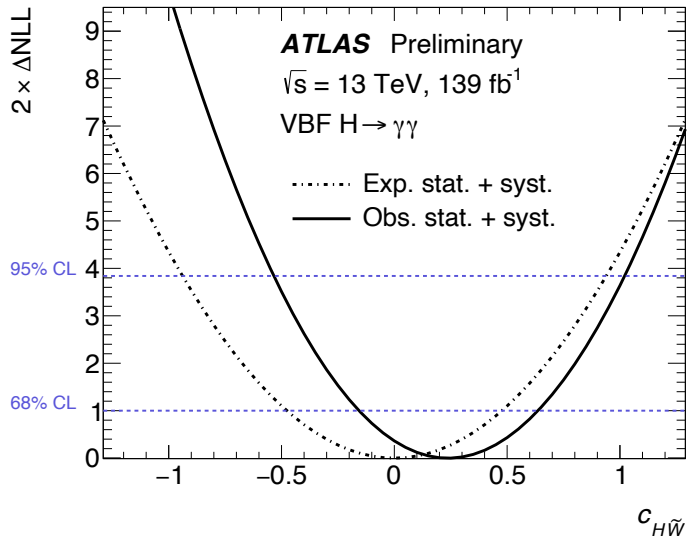
$$OO = \frac{2\text{Re}(\mathcal{M}_{SM}^* \mathcal{M}_{CP\text{-odd}})}{|\mathcal{M}_{SM}|^2}$$



2 dedicated BDT's (VBF vs ggF, VBF vs $\gamma\gamma$)



CP properties of HVV coupling with VBF



$$\tilde{d} = \frac{\nu^2}{\Lambda^2} c_{H\tilde{W}}$$

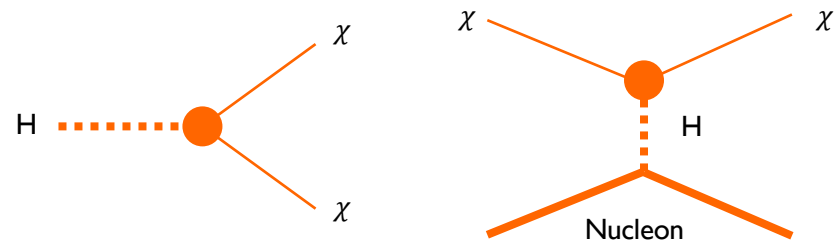
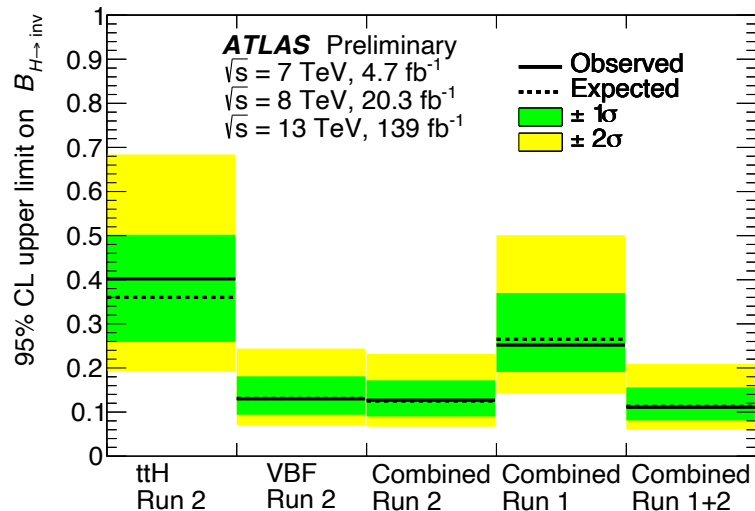
	68% (exp.)	95% (exp.)	68% (obs.)	95% (obs.)
\tilde{d} (inter. only)	[-0.027, 0.027]	[-0.055, 0.055]	[-0.011, 0.036]	[-0.032, 0.059]
\tilde{d} (inter.+quad.)	[-0.028, 0.028]	[-0.061, 0.060]	[-0.010, 0.040]	[-0.034, 0.071]
\tilde{d} from $H \rightarrow \tau\tau$	[-0.038, 0.036]	-	[-0.090, 0.035]	-
Combined \tilde{d}	[-0.022, 0.021]	[-0.046, 0.045]	[-0.012, 0.030]	[-0.034, 0.057]
$c_{H\tilde{W}}$ (inter. only)	[-0.48, 0.48]	[-0.94, 0.94]	[-0.16, 0.64]	[-0.53, 1.02]
$c_{H\tilde{W}}$ (inter.+quad.)	[-0.48, 0.48]	[-0.95, 0.95]	[-0.15, 0.67]	[-0.55, 1.07]

most stringent constraints on CP-properties of HVV coupling to date

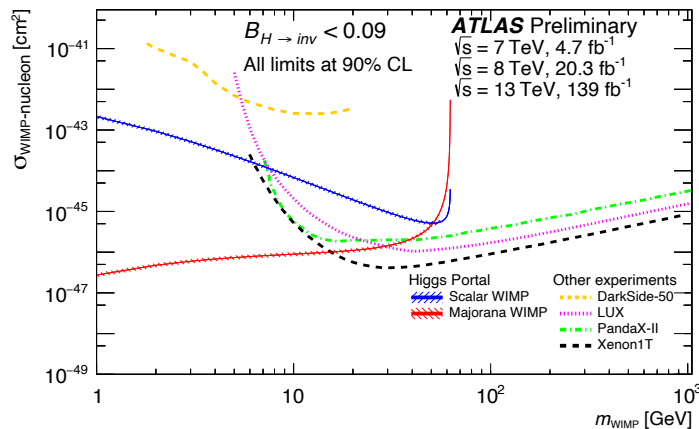
Other information about the Higgs width?

- Part of Higgs width could be due to decays to undetectable particles ($H \rightarrow \text{invisible}$)
 - ✓ Limits from direct searches in $H \rightarrow \text{invisible}$, explicitly requiring MET in the event

[ATLAS-CONF-2020-052](#)



Limits on $BR(H \rightarrow \text{inv})$ can be recast in term of limit on Dark Matter production



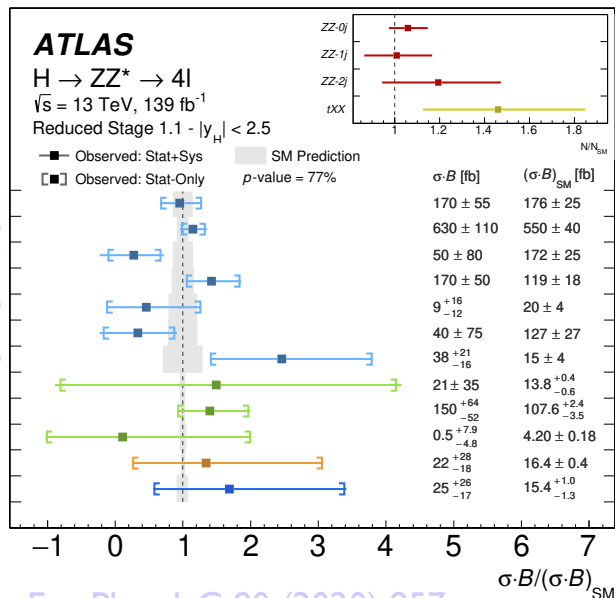
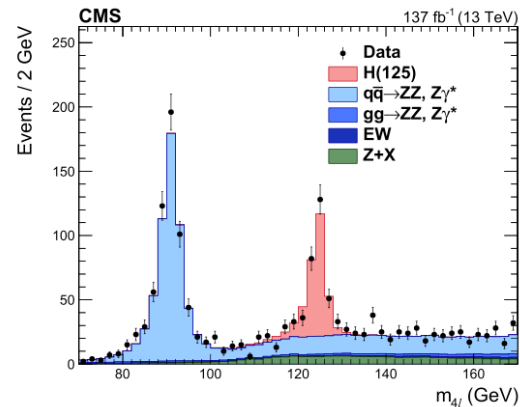
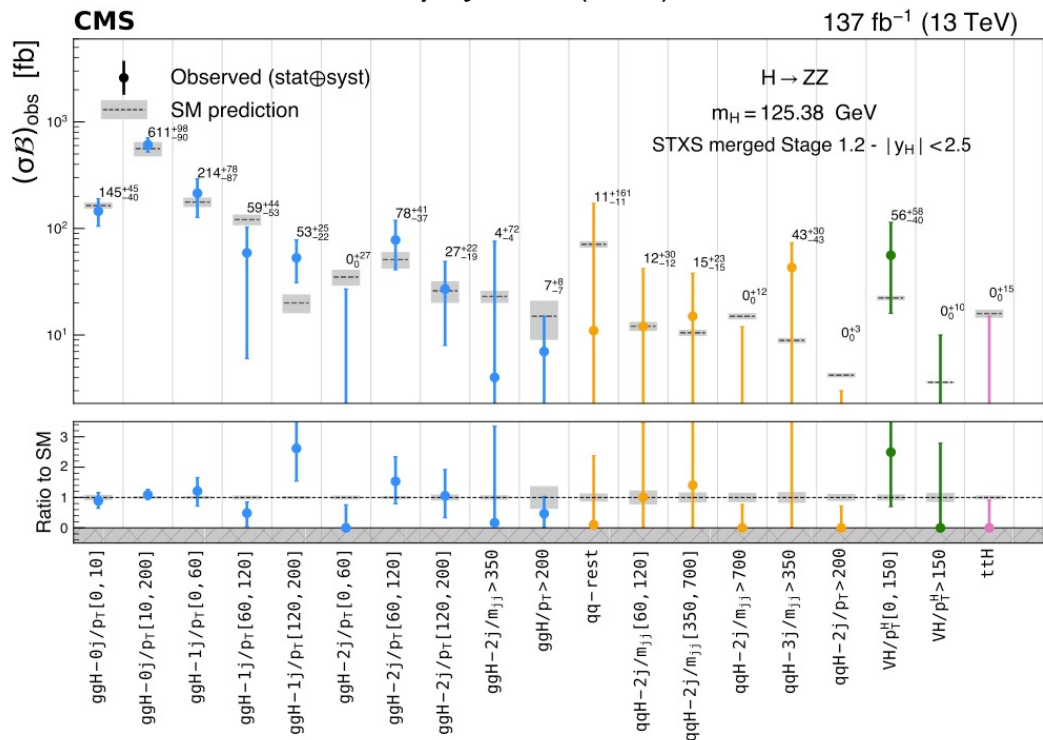
ATLAS: $BR(H \rightarrow \text{inv}) < 0.11$ (exp. 0.11) @ 95% CL

CMS: $BR(H \rightarrow \text{inv}) < 0.17$ (exp. 0.11) @ 95% CL

[CMS-PAS-HIG-20-003](#)

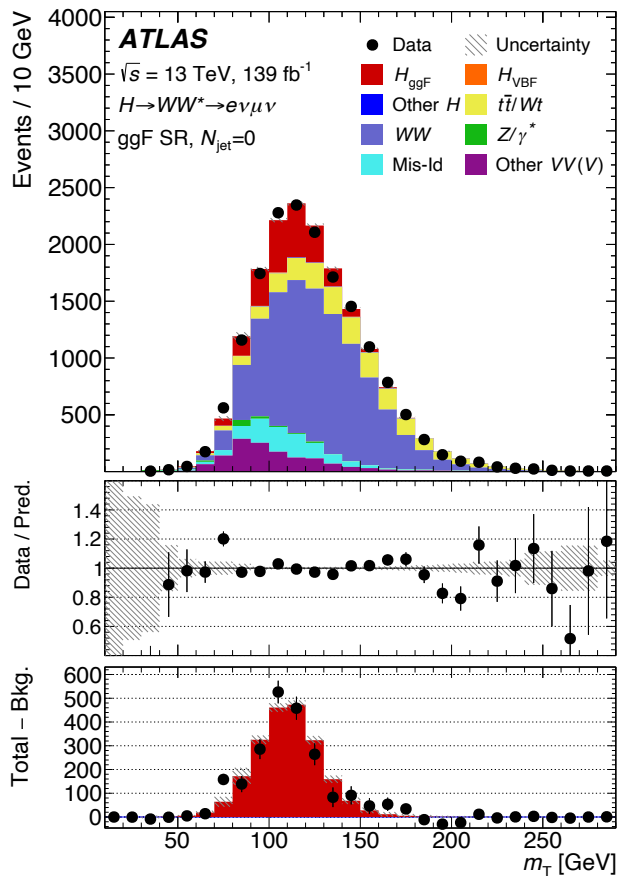
STXS at work: $H \rightarrow ZZ^*$

Eur. Phys. J. C 81 (2021) 488

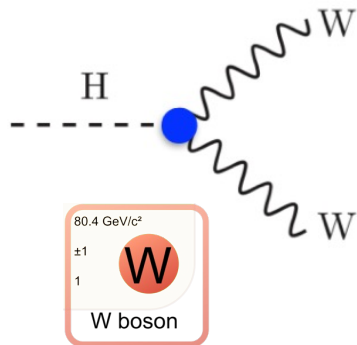
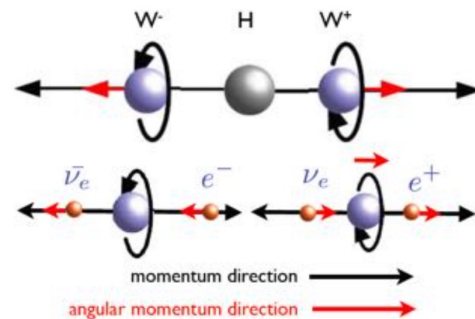


Eur. Phys. J. C 80 (2020) 957

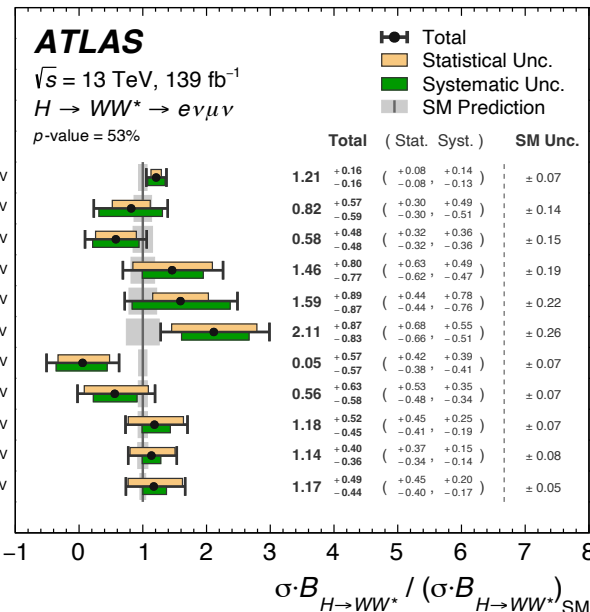
STXS at work: $H \rightarrow WW^*$



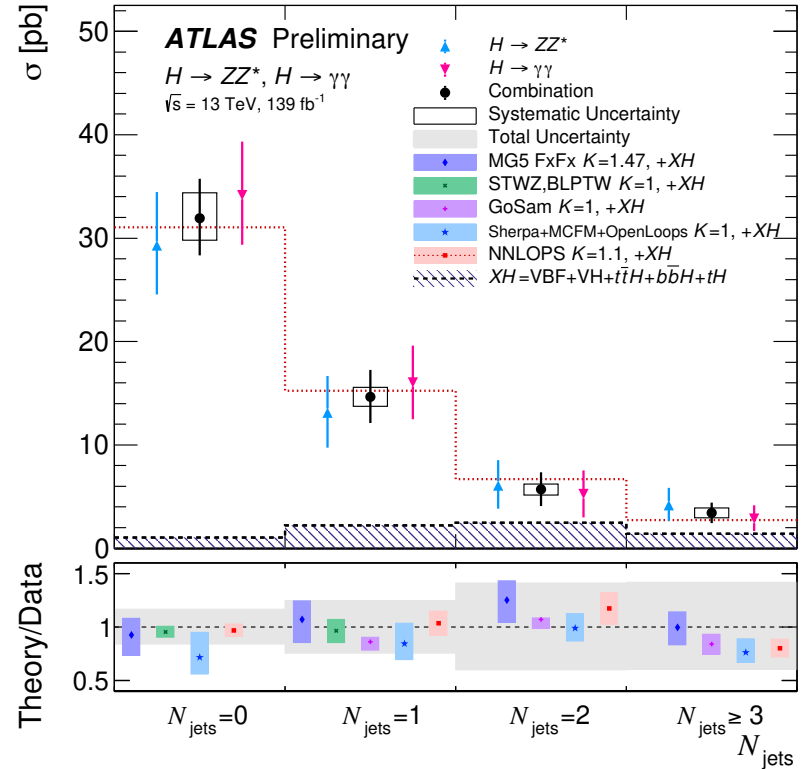
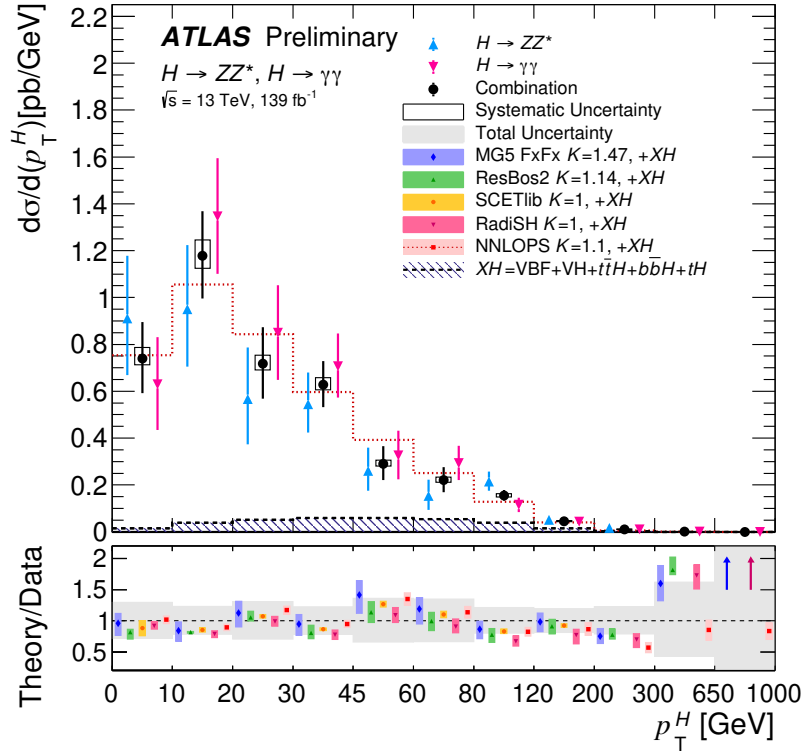
- Precision measurements have become a possibility also for those channels where reconstructing the final state is not possible, e.g. $H \rightarrow WW^*$ (neutrinos, missing transverse energy)!



arXiv:2207.0033

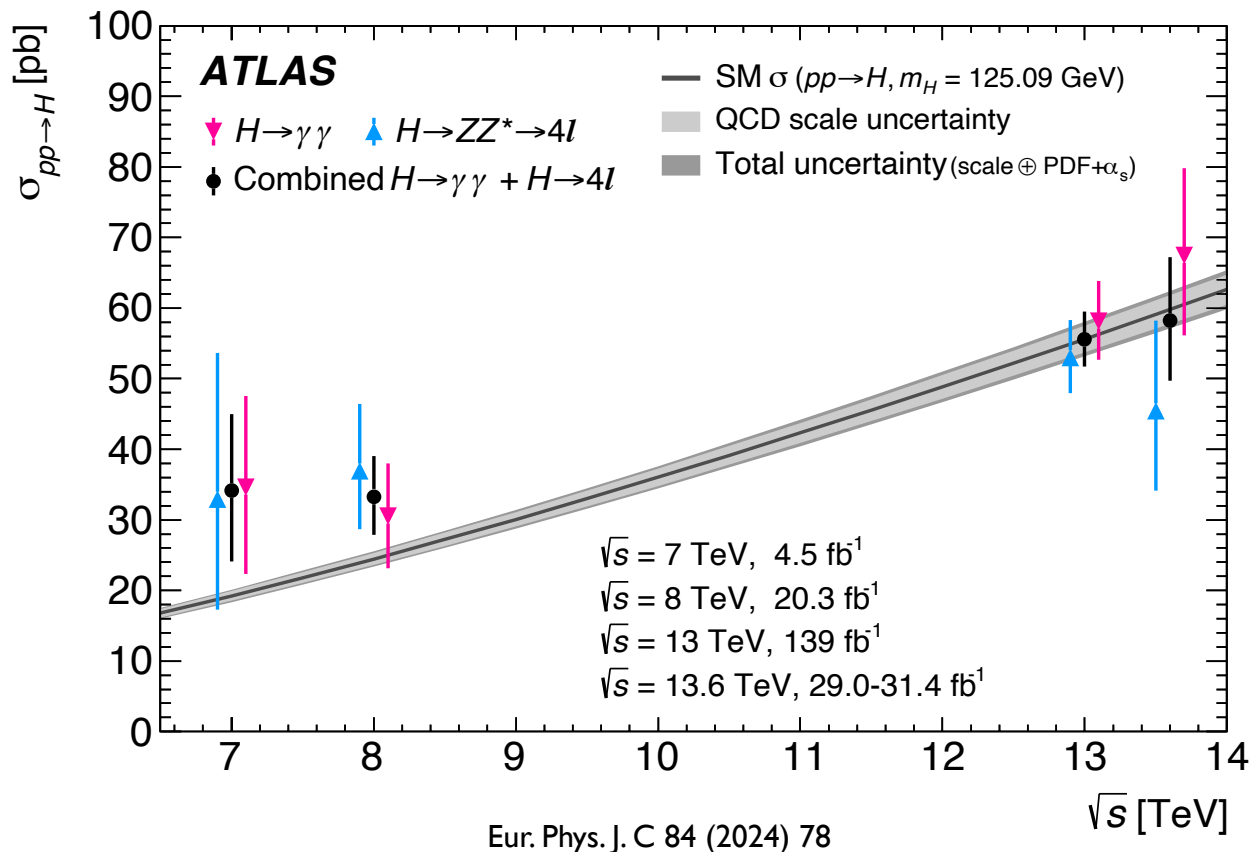


H \rightarrow $\gamma\gamma$ and H \rightarrow ZZ*: differential and fiducial cross-sections



ATLAS-CONF-2022-002

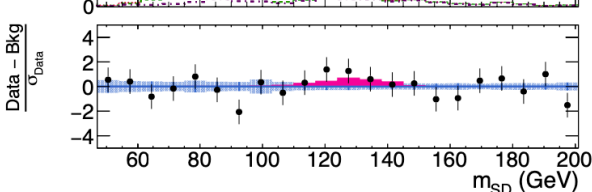
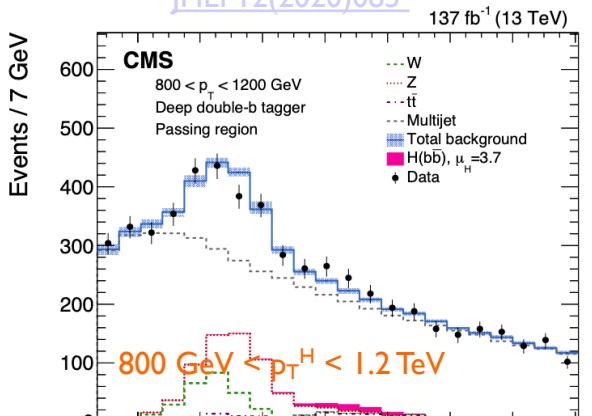
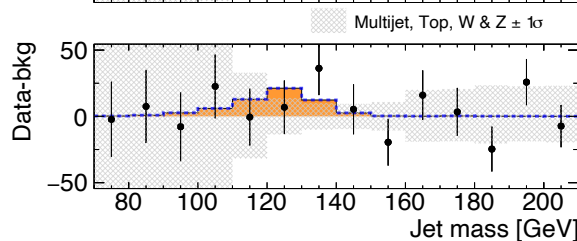
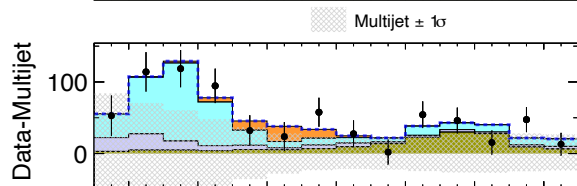
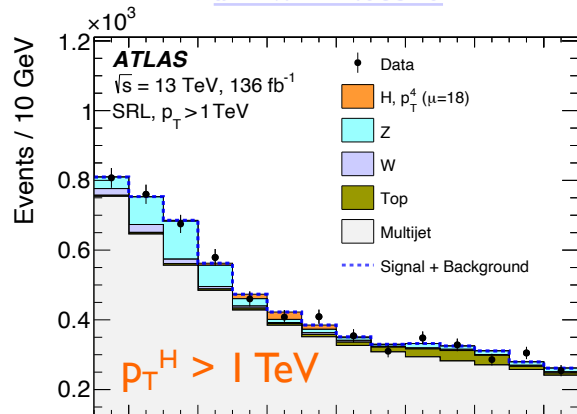
$H \rightarrow \gamma\gamma$ and $H \rightarrow ZZ^*$: inclusive cross-sections



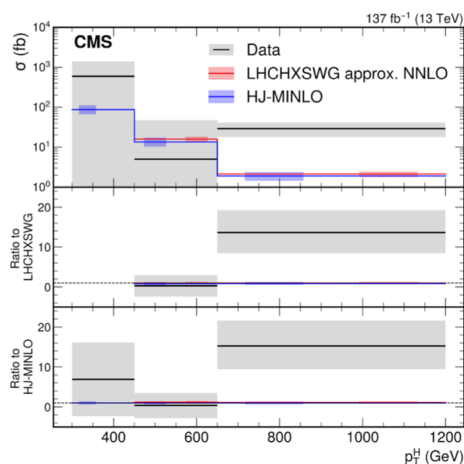
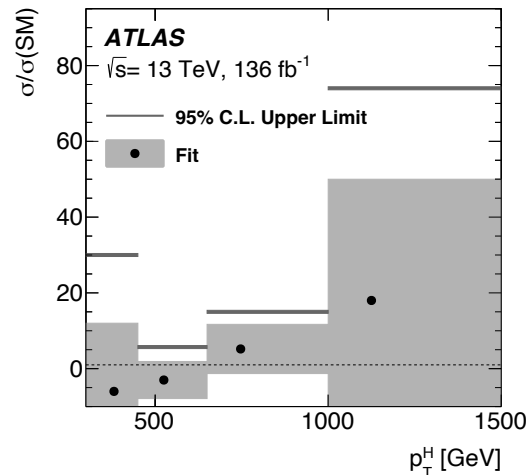
Reaching the highest p_T^H with $H \rightarrow bb$

arXiv:2111.08340

JHEP12(2020)085



- Jets substructure with 1 or 2 b-tags
- Fiducial cross-section (mostly ggF)
- Can be recasted as STXS
- Highest p_T^H most sensitive to BSM effects



Other rare Higgs decays within reach?

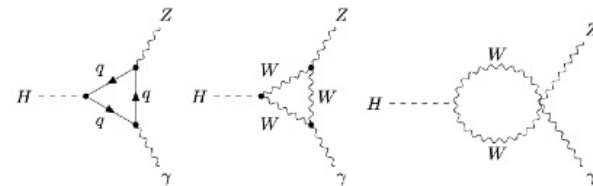
evd. = evidence

Analysis	Final states	13 TeV \mathcal{L} [fb $^{-1}$]	BR (SM)	$\sigma_H \times \text{BR(H decays)}$	ATLAS/CMS references
$c\bar{c}$		139, 35.9	2.9%	$26 \times \text{SM}, 70 \times \text{SM}$	ATLAS-CONF-2021-021, JHEP 03 (2020) 131
invisible		139, 35.9	$\sim 10^{-3}$		ATLAS-CONF-2020-052, PLB 793 (2019) 520
$Z\gamma$	$ee/\mu\mu+\gamma$	139, 35.9	$\sim 10^{-3}$	evd. $\mu = 1.5, 3.9 \times \text{SM}$	arXiv:2103.10322 (PLB 2021), JHEP 11 (2018) 152
$\mu^+\mu^-$		139, 137	$\sim 10^{-4}$	$2.2 \times \text{SM}, \text{evd. } \mu = 1.19$	PLB 812 (2021) 135980, JHEP 01 (2021) 148
BR(H decays)					
$\rho\gamma$	$\pi^+\pi^-\gamma$	35.6	$\sim 10^{-5}$	8.8×10^{-4}	JHEP 07 (2018) 127
$\phi\gamma$	$K^+K^-\gamma$		$\sim 10^{-6}$	4.8×10^{-4}	
$Z\rho$	$ee/\mu\mu + \pi^+\pi^-$	137	$\sim 10^{-5}$	$(1.04-1.31) \times 10^{-2}$	JHEP 11 (2020) 039
$Z\phi$	$ee/\mu\mu + K^+K^-$	137	$\sim 10^{-6}$	$(3-4) \times 10^{-3}$	
$Z\eta_c$	$ee/\mu\mu + \text{had}$	139	$\sim 10^{-5}$	$(\sigma \times \text{BR} = 110 \text{ pb})$	PRL 125 (2020) 221802
ZJ/ψ			$\sim 10^{-6}$	$(\sigma \times \text{BR} = 100 \text{ pb})$	
$J/\psi \gamma$	} $\mu^+\mu^-\gamma$	36.1, 35.9	$\sim 10^{-6}$	$3.5 \times 10^{-4}, 7.6 \times 10^{-4}$	PLB 786 (2018) 134, EPJ C 79 (2019) 94
$\psi(2S)\gamma$		36.1	$\sim 10^{-6}$	2.0×10^{-3}	PLB 786 (2018) 134
$\Upsilon(nS)\gamma$ (n=1,2,3)		36.1	$\sim 10^{-9}$	$(4.9, 5.9, 5.7) \times 10^{-4}$	
$\Upsilon\Upsilon$	} 4μ		$\sim 10^{-9}$	1.4×10^{-3}	
$J/\psi J/\psi$		37.5	$\sim 10^{-10}$	1.8×10^{-3}	PLB 797 (2019) 134811
e^+e^-		139	$\sim 10^{-9}-10^{-10}$	3.6×10^{-4}	PLB 801 (2020) 135148

Imma Riu @ LHCP 2021

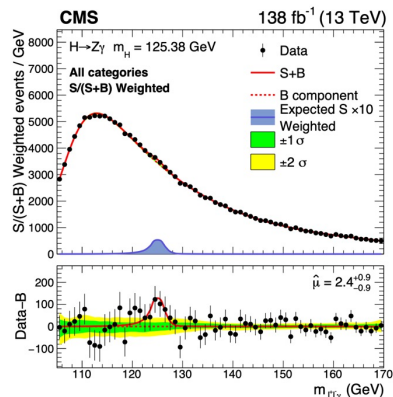
Other rare Higgs decays within reach?

- $SU(2)_L$ symmetry relates the HWW , HZZ , $H\gamma\gamma$ and $HZ\gamma$ interactions
 - ✓ If New Physics respects $SU(2)_L$, effect correlated in all four channels
 - ✓ Important to observed and measure also $H \rightarrow Z\gamma$

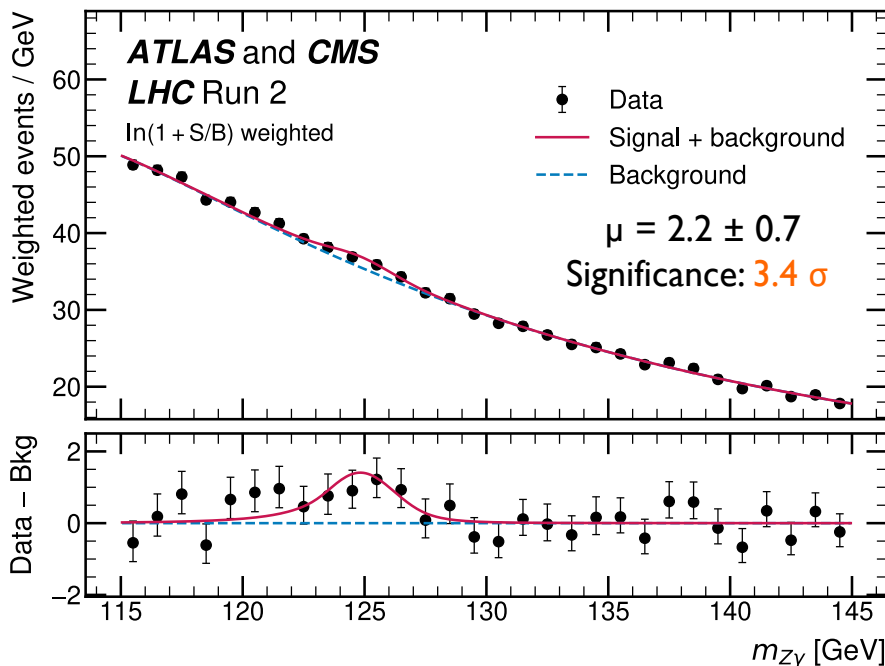


[Phys. Rev. Lett. 132 \(2024\) 021803](#)

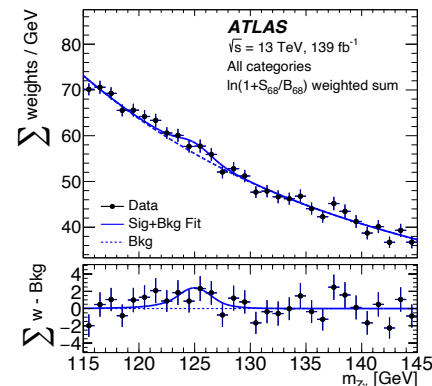
JHEP 05 (2023) 233



$\mu = 2.4 \pm 0.9$
Significance: 2.7σ
(1.2 expected)



[Phys. Lett. B 809 \(2020\) 135754](#)



$\mu = 2.0 \pm 0.9$
Significance: 2.2σ
(1.2 expected)

How well *should* we know the Higgs couplings?

SMALL CORRECTIONS EXPECTED IN MANY BSM MODELS

If new physics is at 1 TeV:

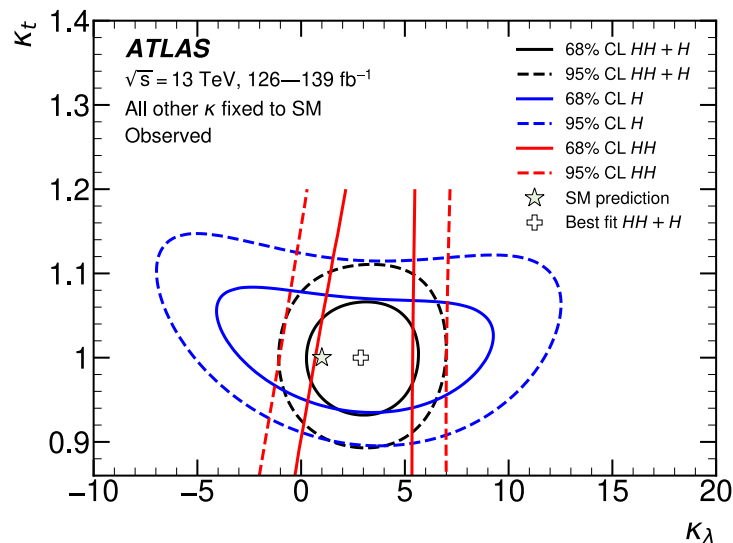
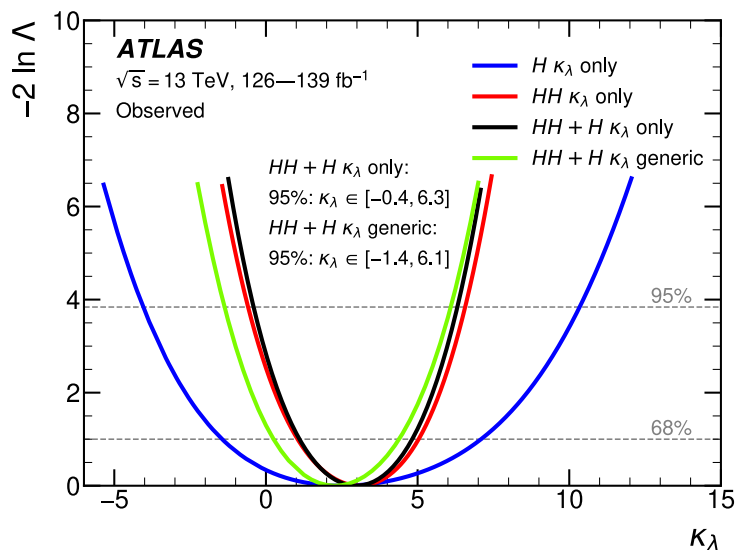
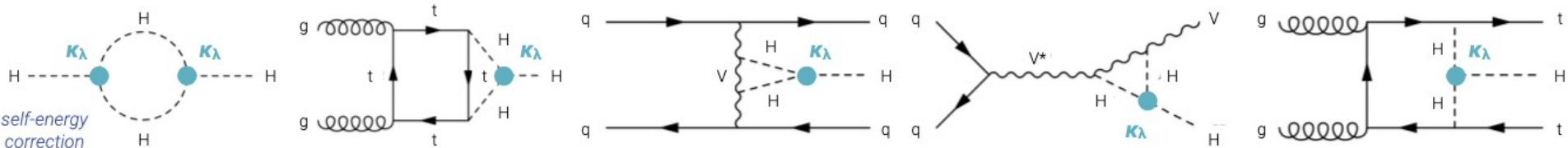
	$\delta\kappa_V$	$\delta\kappa_b$	$\delta\kappa_\gamma$
Singlet	<6%	<6%	<6%
2HDM (large $\tan\beta$)	~1%	~10%	~1%
MSSM	~.001%	~1.6%	~-.4%
Composite	~-3%	~-(3-9)%	~-9%
Top Partner	~-2%	~-2%	~1%

Patterns of deviations can pinpoint specific BSM physics

- Generically new physics effects on couplings $\sim \frac{v^2}{M^2} \sim \mathcal{O}(6\%)$ for $M=1$ TeV
- Only now are we approaching sensitivity where we expect deviations

More tools at hand! Constraining κ_λ with single-Higgs

- Single Higgs boson productions and decays as well as kinematics sensitive to the self-coupling through EW corrections

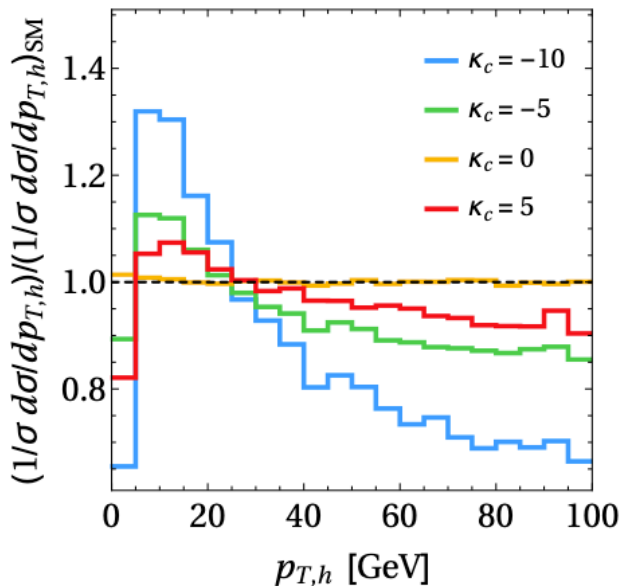


arXiv:2211.01216

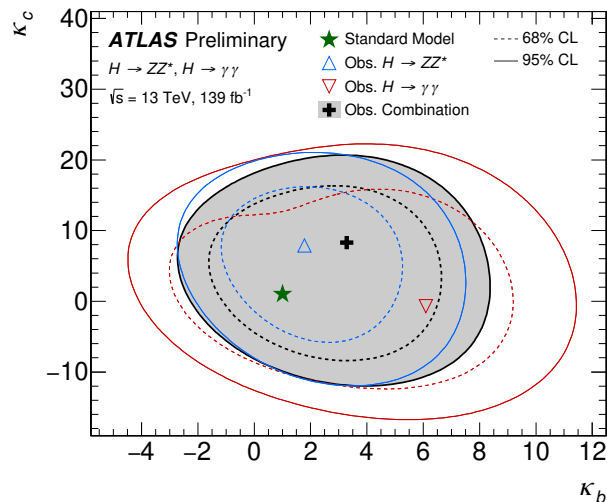
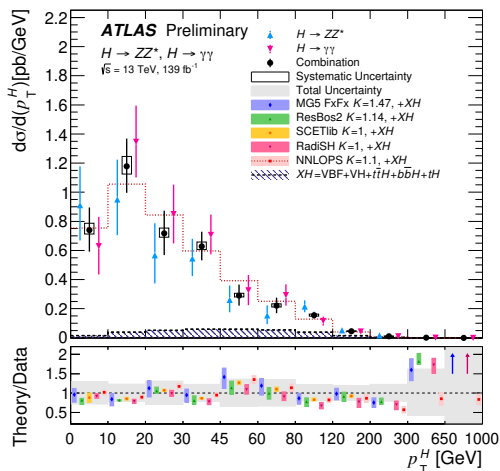
Direct $H \rightarrow cc$ search is not the only tool...

- Charm quark contributes to ggF loop, modification to Higgs-charm coupling can alter $p_{T,H}$
- Analogous effect on quark-initiated production of the Higgs boson
- Shape of differential $p_{T,H}$ cross-section can be interpreted in term of Higgs-charm coupling

[arXiv:1606.09253](https://arxiv.org/abs/1606.09253)



[ATLAS-CONF-2022-002](https://arxiv.org/abs/2202.00022)

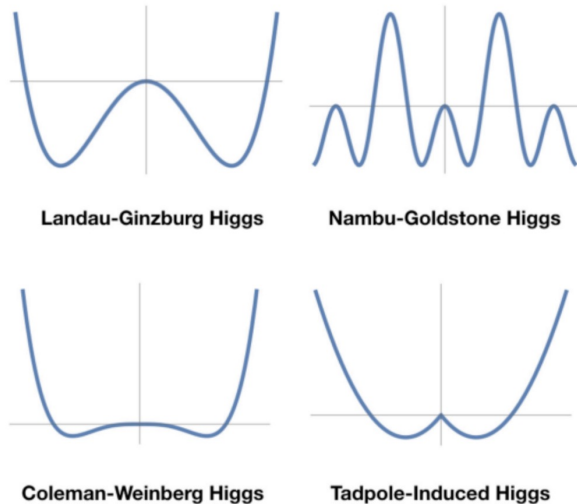


Channel	Parameter best-fit	Observed 95% confidence interval	Expected 95% confidence interval
Combined	$\kappa_b = 3.3^{+2.4}_{-4.1}$	$[-2.1, 7.4]$	$[-2.2, 7.4]$
	$\kappa_c = 8.3^{+5.5}_{-13.8}$	$[-10.1, 18.3]$	$[-10.3, 16.6]$

Is it the Standard Model Higgs potential?

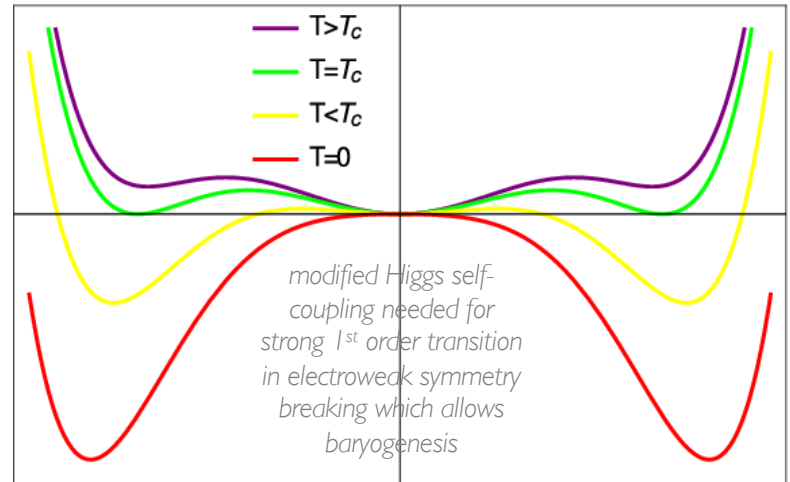
- How the Higgs potential is realized in Nature can provide important answers to many open points in our understanding of Nature...
 - ✓ e.g. BSM scenarios of Higgs sector, electroweak phase transition in the SM, baryogenesis...

Different models predicts different potential shape



[Phys. Rev. D 101, 075023 \(2020\)](#)

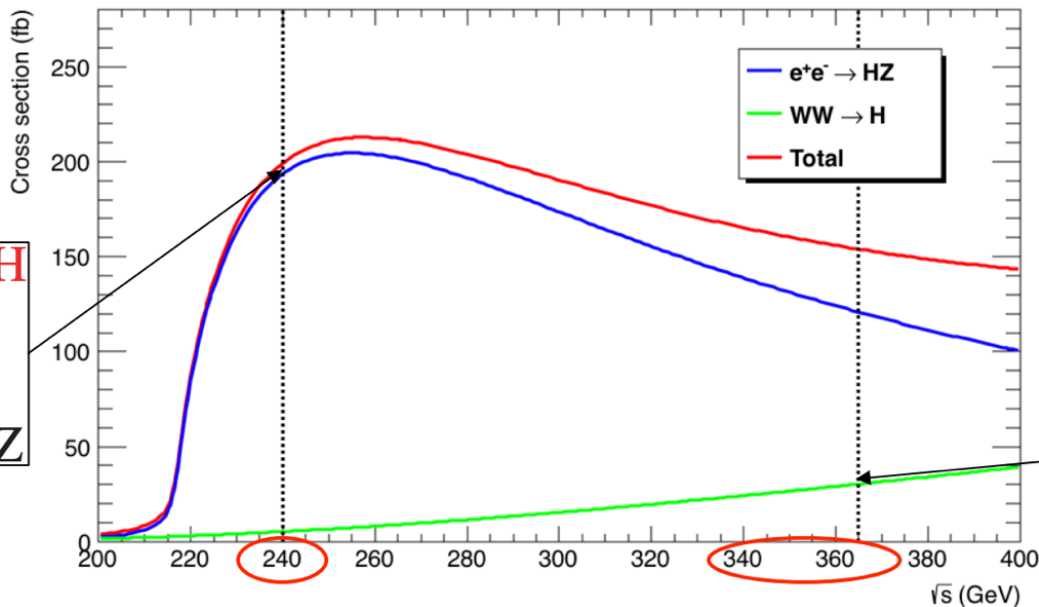
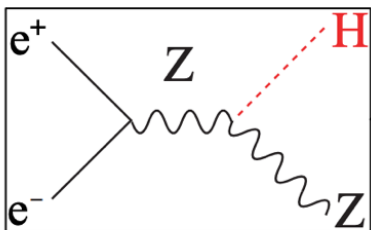
Evolution of Higgs potential could explain matter/antimatter asymmetry via electroweak baryogenesis



[arXiv:1511.03969](#)

Higgs at FC-ee

$L = 5 \text{ ab}^{-1}$
 $ZH = 10^6$
 $VBF = 2 \times 10^4$



$L = 1.5 \text{ ab}^{-1}$
 $ZH = 2.5 \times 10^5$
 $VBF = 5 \times 10^4$

